

# Recent View about Kähler Geometry and Spin Structure of "World of Classical Worlds"

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### Abstract

The construction of Kähler geometry of WCW (“world of classical worlds”) is fundamental to TGD program. I ended up with the idea about physics as WCW geometry around 1985 and made a breakthrough around 1990, when I realized that Kähler function for WCW could correspond to Kähler action for its preferred extremals defining the analogs of Bohr orbits so that classical theory with Bohr rules would become an exact part of quantum theory and path integral would be replaced with genuine integral over WCW. The motivating construction was that for loop spaces leading to a unique Kähler geometry. The geometry for the space of 3-D objects is even more complex than that for loops and the vision still is that the geometry of WCW is unique from the mere existence of Riemann connection.

This chapter represents the updated version of the construction providing a solution to the problems of the previous construction. The basic formulas remain as such but the expressions for WCW super-Hamiltonians defining WCW Hamiltonians (and matrix elements of WCW metric) as their anticommutator are replaced with those following from the dynamics of the Kähler-Dirac action.

## 1 Introduction

The construction of Kähler geometry of WCW (“world of classical worlds”) is fundamental to TGD program. I ended up with the idea about physics as WCW geometry around 1985 and made a breakthrough around 1990, when I realized that Kähler function for WCW could correspond to Kähler action for its preferred extremals defining the analogs of Bohr orbits so that classical theory with Bohr rules would become an exact part of quantum theory and path integral would be replaced with genuine integral over WCW. The motivating construction was that for loop spaces leading to a unique Kähler geometry [A8]. The geometry for the space of 3-D objects is even more complex than that for loops and the vision still is that the geometry of WCW is unique from the mere existence of Riemann connection.

The basic idea is that WCW is union of symmetric spaces  $G/H$  labelled by zero modes which do not contribute to the WCW metric. There have been many open questions and it seems the details of the earlier approach [?] must be modified at the level of detailed identifications and interpretations.

1. A longstanding question has been whether one could assign Equivalence Principle (EP) with the coset representation formed by the super-Virasoro representation assigned to  $G$  and  $H$  in such a way that the four-momenta associated with the representations and identified as inertial and gravitational four-momenta would be identical. This does not seem to be the case. The recent view will be that EP reduces to the view that the classical four-momentum associated with Kähler action is equivalent with that assignable to Kähler-Dirac action supersymmetrically related to Kähler action: quantum classical correspondence (QCC) would be in question. Also strong form of general coordinate invariance implying strong form of holography in turn implying

that the super-symplectic representations assignable to space-like and light-like 3-surfaces are equivalent could imply EP with gravitational and inertial four-momenta assigned to these two representations.

At classical level EP follows from the interpretation of GRT space-time as effective space-time obtained by replacing many-sheeted space-time with Minkowski space with effective metric determined as a sum of Minkowski metric and sum over the deviations of the induced metrics of space-time sheets from Minkowski metric. Poincare invariance suggests strongly classical EP for the GRT limit in long length scales at least.

2. The detailed identification of groups  $G$  and  $H$  and corresponding algebras has been a longstanding problem. Symplectic algebra associated with  $\delta M_{\pm}^4 \times CP^2$  ( $\delta M_{\pm}^4$  is light-cone boundary - or more precisely, with the boundary of causal diamond (CD) defined as Cartesian product of  $CP_2$  with intersection of future and past direct light cones of  $M^4$  has Kac-Moody type structure with light-like radial coordinate replacing complex coordinate  $z$ . Virasoro algebra would correspond to radial diffeomorphisms. I have also introduced Kac-Moody algebra assigned to the isometries and localized with respect to internal coordinates of the light-like 3-surfaces at which the signature of the induced metric changes from Minkowskian to Euclidian and which serve as natural correlates for elementary particles (in very general sense!). This kind of localization by force could be however argued to be rather ad hoc as opposed to the inherent localization of the symplectic algebra containing the symplectic algebra of isometries as sub-algebra. It turns out that one obtains direct sum of representations of symplectic algebra and Kac-Moody algebra of isometries naturally as required by the success of p-adic mass calculations.
3. The dynamics of Kähler action is not visible in the earlier construction. The construction also expressed WCW Hamiltonians as 2-D integrals over partonic 2-surfaces. Although strong form of general coordinate invariance (GCI) implies strong form of holography meaning that partonic 2-surfaces and their 4-D tangent space data should code for quantum physics, this kind of outcome seems too strong. The progress in the understanding of the solutions of Kähler-Dirac equation led however to the conclusion that spinor modes other than right-handed neutrino are localized at string world sheets with strings connecting different partonic 2-surfaces. This leads to a modification of earlier construction in which WCW super-Hamiltonians are essentially integrals with integrand identified as a Noether super current for the deformations in  $G$ . Each spinor mode gives rise to super current and the modes of right-handed neutrino and other fermions differ in an essential ways. Right-handed neutrino would correspond to symplectic algebra and other modes to the Kac-Moody algebra and one obtains the crucial 5 tensor factors of Super Virasoro required by p-adic mass calculations.

The matrix elements of WCW metric between Killing vectors are expressible as anti-commutators of super-Hamiltonians identifiable as contractions of WCW gamma matrices with these vectors and give Poisson brackets of corresponding Hamiltonians. The anti-commutation relates of induced spinor fields are

dictated by this condition. Everything is 3-dimensional although one expects that symplectic transformations localized within interior of  $X^3$  act as gauge symmetries so that in this sense effective 2-dimensionality is achieved. The components of WCW metric are labelled by standard model quantum numbers so that the connection with physics is extremely intimate.

4. An open question in the earlier visions was whether finite measurement resolution is realized as discretization at the level of fundamental dynamics. This would mean that only certain string world sheets from the slicing by string world sheets and partonic 2-surfaces are possible. The requirement that anti-commutations are consistent suggests that string world sheets correspond to surfaces for which Kähler magnetic field is constant along string in well defined sense ( $J_{\mu\nu}\epsilon^{\mu\nu}g^{1/2}$  remains constant along string). It however turns that by a suitable choice of coordinates of 3-surface one can guarantee that this quantity is constant so that no additional constraint results.
5. Quantum criticality is one of the basic notions of quantum TGD and its relationship to coset construction has remained unclear. In this chapter the concrete realization of criticality in terms of symmetry breaking hierarchy of Super Virasoro algebra acting on symplectic and Kac-Moody algebras. Also a connection with finite measurement resolution - second key notion of TGD - emerges naturally.

The appendix of the book gives a summary about basic concepts of TGD with illustrations. Pdf representation of same files serving as a kind of glossary can be found at <http://tgdtheory.fi/tgdglossary.pdf> [?].

## 2 WCW As A Union Of Homogenous Or Symmetric Spaces

The physical interpretation and detailed mathematical understanding of superconformal symmetries has developed rather slowly and has involved several side tracks. In the following I try to summarize the basic picture with minimal amount of formulas with the understanding that the statement “Noether charge associated with geometrically realized Kac-Moody symmetry” is enough for the reader to write down the needed formula explicitly. Formula oriented reader might deny the value of the approach giving weight to principles. My personal experience is that piles of formulas too often hide the lack of real understanding.

In the following the vision about WCW as union of coset spaces is discussed in more detail.

### 2.1 Basic Vision

The basic view about coset space construction for WCW has not changed.

1. The idea about WCW as a union of coset spaces  $G/H$  labelled by zero modes is extremely attractive. The structure of homogenous space [A1] (<http://>

`tinyurl.com/y7u2t8jo` ) means at Lie algebra level the decomposition  $g = h \oplus t$  to sub-Lie-algebra  $h$  and its complement  $t$  such that  $[h, t] \subset t$  holds true. Homogeneous spaces have  $G$  as its isometries. For symmetric space the additional condition  $[t, t] \subset h$  holds true and implies the existence of involution changing at the Lie algebra level the sign of elements of  $t$  and leaving the elements of  $h$  invariant. The assumption about the structure of symmetric space [A4] (<http://tinyurl.com/ycouv7uh> ) implying covariantly constant curvature tensor is attractive in infinite-dimensional case since it gives hopes about calculability.

An important source of intuition is the analogy with the construction of  $CP_2$ , which is symmetric space. A particular choice of  $h$  corresponds to Lie-algebra elements realized as Killing vector fields which vanish at particular point of WCW and thus leave 3-surface invariant. A preferred choice for this point is as maximum or minimum of Kähler function. For this 3-surface the Hamiltonians of  $h$  should be stationary. If symmetric space property holds true then commutators of  $[t, t]$  also vanish at the minimum/maximum. Note that Euclidian signature for the metric of WCW requires that Kähler function can have only maximum or minimum for given zero modes.

2. The basic objection against TGD is that one cannot use the powerful canonical quantization using the phase space associated with configuration space - now WCW . The reason is the extreme non-linearity of the Kähler action and its huge vacuum degeneracy, which do not allow the construction of Hamiltonian formalism. Symplectic and Kähler structure must be realized at the level of WCW . In particular, Hamiltonians must be represented in completely new manner. The key idea is to construct WCW Hamiltonians as anti-commutators of super-Hamiltonians defining the contractions of WCW gamma matrices with corresponding Killing vector fields and therefore defining the matrix elements of WCW metric in the tangent vector basis defined by Killing vector fields. Super-symmetry therefore gives hopes about constructing quantum theory in which only induced spinor fields are second quantized and embedding space coordinates are treated purely classically.
3. It is important to understand the difference between symmetries and isometries assigned to the Kähler function. Symmetries of Kähler function do not affect it. The symmetries of Kähler action are also symmetries of Kähler action because Kähler function is Kähler action for a preferred extremal (here there have been a lot of confusion). Isometries leave invariant only the quadratic form defined by Kähler metric  $g_{M\bar{N}} = \partial_M \partial_{\bar{L}} K$  but not Kähler function in general. For  $G/H$  decomposition  $G$  represents isometries and  $H$  both isometries and symmetries of Kähler function.

$CP_2$  is familiar example:  $SU(3)$  represents isometries and  $U(2)$  leaves also Kähler function invariant since it depends on the  $U(2)$  invariant radial coordinate  $r$  of  $CP_2$ . The origin  $r = 0$  is left invariant by  $U(2)$  but for  $r > 0$   $U(2)$  performs a rotation at  $r = \text{constant}$  3-sphere. This simple picture helps to understand what happens at the level of WCW .

How to then distinguish between symmetries and isometries? A natural guess

is that one obtains also for the isometries Noether charges but the vanishing of boundary terms at spatial infinity crucial in the argument leading to Noether theorem as  $\Delta S = \Delta Q = 0$  does not hold true anymore and one obtains charges which are not conserved anymore. The symmetry breaking contributions would now come from effective boundaries defined by wormhole throats at which the induce metric changes its signature from Minkowskian to Euclidian. A more delicate situation is in which first order contribution to  $\Delta S$  vanishes and therefore also  $\Delta Q$  and the contribution to  $\Delta S$  comes from second variation allowing also to define Noether charge which is not conserved.

4. The simple picture about  $CP_2$  as symmetric space helps to understand the general vision if one assumes that WCW has the structure of symmetric space. The decomposition  $g = h + t$  corresponds to decomposition of symplectic deformations to those which vanish at 3-surface ( $h$ ) and those which do not ( $t$ ).

For the symmetric space option, the Poisson brackets for super generators associated with  $t$  give Hamiltonians of  $h$  identifiable as the matrix elements of WCW metric. They would not vanish although they are stationary at 3-surface meaning that Riemann connection vanishes at 3-surface. The Hamiltonians which vanish at 3-surface  $X^3$  would correspond to  $t$  and the Hamiltonians for which Killing vectors vanish and which therefore are stationary at  $X^3$  would correspond to  $h$ . Outside  $X^3$  the situation would of course be different. The metric would be obtained by parallel translating the metric from the preferred point of WCW to elsewhere and symplectic transformations would make this parallel translation.

For the homogenous space option the Poisson brackets for super generators of  $t$  would still give Hamiltonians identifiable as matrix elements of WCW metric but now they would be necessary those of  $h$ . In particular, the Hamiltonians of  $t$  do not in general vanish at  $X^3$ .

## 2.2 Equivalence Principle And WCW

### 2.3 Equivalence Principle At Quantum And Classical Level

Quite recently I returned to an old question concerning the meaning of Equivalence Principle (EP) in TGD framework.

Heretic would of course ask whether the question about whether EP is true or not is a pseudo problem due to uncritical assumption there really are two different four-momenta which must be identified. If even the identification of these two different momenta is difficult, the pondering of this kind of problem might be waste of time.

At operational level EP means that the scattering amplitudes mediated by graviton exchange are proportional to the product of four-momenta of particles and that the proportionality constant does not depend on any other parameters characterizing particle (except spin). The are excellent reasons to expect that the stringy picture for interactions predicts this.

1. The old idea is that EP reduces to the coset construction for Super Virasoro algebra using the algebras associated with  $G$  and  $H$ . The four-momenta



assignable to these algebras would be identical from the condition that the differences of the generators annihilate physical states and identifiable as inertial and gravitational momenta. The objection is that for the preferred 3-surface  $H$  by definition acts trivially so that time-like translations leading out from the boundary of CD cannot be contained by  $H$  unlike  $G$ . Hence four-momentum is not associated with the Super-Virasoro representations assignable to  $H$  and the idea about assigning EP to coset representations does not look promising.

2. Another possibility is that EP corresponds to quantum classical correspondence (QCC) stating that the classical momentum assignable to Kähler action is identical with gravitational momentum assignable to Super Virasoro representations. This forced to reconsider the questions about the precise identification of the Kac-Moody algebra and about how to obtain the magic five tensor factors required by p-adic mass calculations [K25].

A more precise formulation for EP as QCC comes from the observation that one indeed obtains two four-momenta in TGD approach. The classical four-momentum assignable to the Kähler action and that assignable to the Kähler-Dirac action. This four-momentum is an operator and QCC would state that given eigenvalue of this operator must be equal to the value of classical four-momentum for the space-time surfaces assignable to the zero energy state in question. In this form EP would be highly non-trivial. It would be justified by the Abelian character of four-momentum so that all momentum components are well-defined also quantum mechanically. One can also consider the splitting of four-momentum to longitudinal and transversal parts as done in the parton model for hadrons: this kind of splitting would be very natural at the boundary of CD. The objection is that this correspondence is nothing more than QCC.

3. A further possibility is that duality of light-like 3-surfaces and space-like 3-surfaces holds true. This is the case if the action of symplectic algebra can be defined at light-like 3-surfaces or even for the entire space-time surfaces. This could be achieved by parallel translation of light-cone boundary providing slicing of CD. The four-momenta associated with the two representations of super-symplectic algebra would be naturally identical and the interpretation would be in terms of EP.

One should also understand how General Relativity and EP emerge at classical level. The understanding comes from the realization that GRT is only an effective theory obtained by endowing  $M^4$  with effective metric.

1. The replacement of superposition of fields with superposition of their effects means replacing superposition of fields with the set-theoretic union of space-time surfaces. Particle experiences sum of the effects caused by the classical fields at the space-time sheets.
2. This is true also for the classical gravitational field defined by the deviation from flat Minkowski metric in standard  $M^4$  coordinates for the space-time sheets. One can define effective metric as sum of  $M^4$  metric and deviations. This effective metric would correspond to that of General Relativity. This resolves long standing issues relating to the interpretation of TGD.

3. Einstein's equations could hold true for the effective metric. They are motivated by the underlying Poincare invariance which cannot be realized as global conservation laws for the effective metric. The conjecture vanishing of divergence of Kähler energy momentum tensor can be seen as the microscopic justification for the claim that Einstein's equations hold true for the effective space-time.
4. The breaking of Poincare invariance could have interpretation as effective breaking in zero energy ontology (ZEO), in which various conserved charges are length dependent and defined separately for each causal diamond (CD).

One can of course consider the possibility that Einstein's equations generalize for preferred extremals of Kähler action. This would actually represent at space-time level the notion of QCC rather than realise QCC interpreted as EP. The condition that the energy momentum tensor for Kähler action has vanishing covariant divergence would be satisfied in GRT if Einstein's equations with cosmological term hold true. This is the case also now but one can consider also more general solutions in which one has two cosmological constants which are not genuine constants anymore: this idea is however not promising.

An interesting question is whether inertial-gravitational duality generalizes to the case of color gauge charges so that color gauge fluxes would correspond to "gravitational" color charges and the charges defined by the conserved currents associated with color isometries would define "inertial" color charges. Since the induced color fields are proportional to color Hamiltonians multiplied by Kähler form they vanish identically for vacuum extremals in accordance with "gravitational" color confinement.

## 2.4 Criticism Of The Earlier Construction

The earlier detailed realization of super-Hamiltonians and Hamiltonians can be criticized.

1. Even after these more than twenty years it looks strange that the Hamiltonians should reduce to flux integrals over partonic 2-surfaces. The interpretation has been in terms of effective 2-dimensionality suggested strongly by strong form of general coordinate invariance stating that the descriptions based on light-like orbits of partonic 2-surfaces and space-like three surfaces at the ends of causal diamonds are dual so that only partonic 2-surfaces and 4-D tangent space data at them would matter. Strong form of holography implies effective 2-dimensionality but this should correspond gauge character for the action of symplectic generators in the interior the space-like 3-surfaces at the ends of CDs, which is something much milder.

One expects that the strings connecting partonic 2-surfaces could bring something new to the earlier simplistic picture. The guess is that embedding space Hamiltonian assignable to a point of partonic 2-surface should be replaced with that defined as integral over string attached to the point. Therefore the earlier picture would suffer no modification at the level of general formulas.

2. The fact that the dynamics of Kähler action and Kähler-Dirac action are not directly involved with the earlier construction raises suspicions. I have proposed that Kähler function could allow identification as Dirac determinant [K29] but one would expect more intimate connection. Here the natural question is whether super-Hamiltonians for the Kähler-Dirac action could correspond to Kähler charges constructible using Noether's theorem for corresponding deformations of the space-time surface and would also be identifiable as WCW gamma matrices.

## 2.5 Is WCW Homogenous Or Symmetric Space?

A key question is whether WCW can be symmetric space [A4] (<http://tinyurl.com/y8ojg1kb>) or whether only homogenous structure is needed. The lack of covariant constancy of curvature tensor might produce problems in infinite-dimensional context.

The algebraic conditions for symmetric space are  $g = h + t$ ,  $[h, t] \subset t$ ,  $[t, t] \subset h$ . The latter condition is the difficult one.

1.  $\delta CD$  Hamiltonians should induce diffeomorphisms of  $X^3$  indeed leaving it invariant. The symplectic vector fields would be parallel to  $X^3$ . A stronger condition is that they induce symplectic transformations for which all points of  $X^3$  remain invariant. Now symplectic vector fields vanish at preferred 3-surface (note that the symplectic transformations are  $r_M$  local symplectic transformations of  $S^2 \times CP_2$ ).
2. For Kac-Moody algebra inclusion  $H \subset G$  for the finite-dimensional Lie-algebra induces the structure of symmetric space. If entire algebra is involved this does not look physically very attractive idea unless one believes on symmetry breaking for both  $SU(3)$ ,  $U(2)_{ew}$ , and  $SO(3)$  and  $E_2$  (here complex conjugation corresponds to the involution). If one assumes only Kac-Moody algebra as critical symmetries, the number of tensor factors is 4 instead of five, and it is not clear whether one can obtain consistency with p-adic mass calculations.

Examples of 3-surfaces remaining invariant under  $U(2)$  are 3-spheres of  $CP_2$ . They could correspond to intersections of deformations of  $CP_2$  type vacuum extremals with the boundary of CD. Also geodesic spheres  $S^2$  of  $CP_2$  are invariant under  $U(2)$  subgroup and would relate naturally to cosmic strings. The corresponding 3-surface would be  $L \times S^2$ , where  $L$  is a piece of light-like radial geodesic.

3. In the case of symplectic algebra one can construct the embedding space Hamiltonians inducing WCW Hamiltonians as products of elements of the isometry algebra of  $S^2 \times CP_2$  for with parity under involution is well-defined. This would give a decomposition of the symplectic algebra satisfying the symmetric space property at the level embedding space. This decomposition does not however look natural at the level of WCW since the only single point of  $CP_2$  and light-like geodesic of  $\delta M_+^4$  can be fixed by  $SO(2) \times U(2)$  so that the 3-surfaces would reduce to pieces of light rays.

4. A more promising involution is the inversion  $r_M \rightarrow 1/r_M$  of the radial coordinate mapping the radial conformal weights to their negatives. This corresponds to the inversion in Super Virasoro algebra.  $t$  would correspond to functions which are odd functions of  $u \equiv \log(r_M/r_0)$  and  $h$  to even function of  $u$ . Stationary 3-surfaces would correspond to  $u = 1$  surfaces for which  $\log(u) = 0$  holds true. This would assign criticality with Virasoro algebra as one expects on general grounds.

$r_M = \text{constant}$  surface would most naturally correspond to a maximum of Kähler function which could indeed be highly symmetric. The elements with even  $u$ -parity should define Hamiltonians, which are stationary at the maximum of Kähler function. For other 3-surfaces obtained by  $/r_M$ -local) symplectic transformations the situation is different: now  $H$  is replaced with its symplectic conjugate  $hHg^{-1}$  of  $H$  is acceptable and if the conjecture is true one would obtained 3-surfaces assignable to perturbation theory around given maximum as symplectic conjugates of the maximum. The condition that  $H$  leaves  $X^3$  invariant in poin-twise manner is certainly too strong and imply that the 3-surface has single point as  $CP_2$  projection.

5. One can also consider the possibility that critical deformations correspond to  $h$  and non-critical ones to  $t$  for the preferred 3-surface. Criticality for given  $h$  would hold only for a preferred 3-surface so that this picture would be very similar that above. Symplectic conjugates of  $h$  would define criticality for other 3-surfaces. WCW would decompose to a union corresponding to different criticalities perhaps assignable to the hierarchy of sub-algebras of conformal algebra labelled by integer whose multiples give the allowed conformal weights. Hierarchy of breakings of conformal symmetries would characterize this hierarchy of sectors of WCW .

For sub-algebras of the conformal algebras (Kac-Moody and symplectic algebra) the condition  $[t, t] \subset h$  cannot hold true so that one would obtain only the structure of homogenous space.

## 2.6 Symplectic And Kac-Moody Algebras As Basic Building Bricks

### 3 Updated View About Kähler Geometry Of WCW

During last years the understanding of the mathematical aspects of TGD and of its connection with the experimental world has developed rapidly.

TGD differs in several respects from quantum field theories and string models. The basic mathematical difference is that the mathematically poorly defined notion of path integral is replaced with the mathematically well-defined notion of functional integral defined by the Kähler function defining Kähler metric for WCW (“world of classical worlds”). Apart from quantum jump, quantum TGD is essentially theory of classical WCW spinor fields with WCW spinors represented as fermionic Fock states. One can say that Einstein’s geometrization of physics program is generalized to the level of quantum theory.

It has been clear from the beginning that the gigantic super-conformal symmetries generalizing ordinary super-conformal symmetries are crucial for the existence of WCW Kähler metric. The detailed identification of Kähler function and WCW Kähler metric has however turned out to be a difficult problem. It is now clear that WCW geometry can be understood in terms of the analog of AdS/CFT duality between fermionic and space-time degrees of freedom (or between Minkowskian and Euclidian space-time regions) allowing to express Kähler metric either in terms of Kähler function or in terms of anti-commutators of WCW gamma matrices identifiable as super-conformal Noether super-charges for the symplectic algebra assignable to  $\delta M_{\pm}^4 \times CP_2$ . The string model type description of gravitation emerges and also the TGD based view about dark matter becomes more precise. String tension is however dynamical rather than pregiven and the hierarchy of Planck constants is necessary in order to understand the formation of gravitationally bound states. Also the proposal that sparticles correspond to dark matter becomes much stronger: sparticles actually are dark variants of particles.

A crucial element of the construction is the assumption that super-symplectic and other super-conformal symmetries having the same structure as 2-D super-conformal groups can be seen as broken gauge symmetries such that sub-algebra with conformal weights coming as  $n$ -ples of those for full algebra act as gauge symmetries. In particular, the Noether charges of this algebra vanish for preferred extremals- this would realize the strong form of holography implied by strong form of General Coordinate Invariance. This gives rise to an infinite number of hierarchies of conformal gauge symmetry breakings with levels labelled by integers  $n(i)$  such that  $n(i)$  divides  $n(i + 1)$  interpreted as hierarchies of dark matter with levels labelled by the value of Planck constant  $h_{eff} = n \times h$ . These hierarchies define also hierarchies of quantum criticalities, and are proposed to give rise to inclusion hierarchies of hyperfinite factors of  $II_1$  having interpretation in terms of finite cognitive resolution with inclusions being characterized by the integers  $n(+1)/n(i)$ .

These hierarchies are fundamental for the understanding of living matter. Living matter is fighting in order to stay at criticality and uses metabolic energy and homeostasis to achieve this. In the biological death of the system (self) a phase transition increasing  $h_{eff}$  finally takes place. The sub-selves of self experienced by self as mental images however die and are reborn at opposite boundary of the corresponding causal diamond (CD) and they genuinely evolve so that self can be said to become wiser even without dying! The purpose of this fighting against criticality would thus allow a possibility for sub-selves to evolve via subsequent reincarnations. One interesting prediction is the possibility of time reversed mental images. The challenge is to understand what they do mean at the level of conscious experience.

### **3.1 Kähler Function, Kähler Action, And Connection With String Models**

The definition of Kähler function in terms of Kähler action is possible because space-time regions can have also Euclidian signature of induced metric. Euclidian regions with 4-D  $CP_2$  projection - wormhole contacts - are identified as lines of generalized Feynman diagrams - space-time correlates for basic building bricks of

elementary particles. Kähler action from Minkowskian regions is imaginary and gives to the functional integrand a phase factor crucial for quantum field theoretic interpretation. The basic challenges are the precise specification of Kähler function of “world of classical worlds” ( WCW ) and Kähler metric.

There are two approaches concerning the definition of Kähler metric: the conjecture analogous to AdS/CFT duality is that these approaches are mathematically equivalent.

1. The Kähler function defining Kähler metric can be identified as Kähler action for space-time regions with Euclidian signature for a preferred extremal containing 3-surface as the ends of the space-time surfaces inside causal diamond (CD). Minkowskian space-time regions give to Kähler action an imaginary contribution interpreted as the counterpart of quantum field theoretic action. The exponent of Kähler function gives rise to a mathematically well-defined functional integral in WCW . WCW metric is dictated by the Euclidian regions of space-time with 4-D  $CP_2$  projection.

The basic question concerns the attribute ”preferred”. Physically the preferred extremal is analogous to Bohr orbit. What is the mathematical meaning of preferred extremal of Kähler action? The latest step of progress is the realization that the vanishing of generalized conformal charges for the ends of the space-time surface fixes the preferred extremals to high extent and is nothing but classical counterpart for generalized Virasoro and Kac-Moody conditions.

2. Fermions are also needed. The well-definedness of electromagnetic charge led to the hypothesis that spinors are restricted at string world sheets. One could also consider associativity as basic constraint to fermionic dynamics combined with the requirement that octonionic representation for gamma matrices is equivalent with the ordinary one. The conjecture is that this leads to the same outcome. This point is highly non-trivial and will be discussed below separately.
3. Second manner to define Kähler metric is as anticommutators of WCW gamma matrices identified as super-symplectic Noether charges for the Dirac action for induced spinors with string tension proportional to the inverse of Newton’s constant. These charges are associated with the 1-D space-like ends of string world sheets connecting the wormhole throats. WCW metric contains contributions from the spinor modes associated with various string world sheets connecting the partonic 2-surfaces associated with the 3-surface.

It is clear that the information carried by WCW metric about 3-surface is rather limited and that the larger the number of string world sheets, the larger the information. This conforms with strong form of holography and the notion of measurement resolution as a property of quantum state. Duality clearly means that Kähler function is determined either by space-time dynamics inside Euclidian wormhole contacts or by the dynamics of fermionic strings in Minkowskian regions outside wormhole contacts. This duality brings strongly in mind AdS/CFT duality. One could also speak about fermionic emergence since Kähler function is dictated by the Kähler metric part from a real part of

gradient of holomorphic function: a possible identification of the exponent of Kähler function is as Dirac determinant.

## 3.2 Symmetries of WCW

Towards the end of year 2023 a dramatic progress in the understanding of WCW geometry took place and the following piece of text summarizes the findings. It turned that the original intuitive picture was surprisingly near to what now looks the correct view.

### 3.2.1 The situation before 2023

WCW geometry exists only if it has maximal isometries. I have proposed that WCW could be regarded as a union of generalized symmetric spaces labelled by zero modes which do not contribute to the metric. The induced Kähler field is invariant under symplectic transformations of  $CP_2$  and would therefore define zero mode degrees of freedom if one assumes that WCW metric has symplectic transformations as isometries. In particular, Kähler magnetic fluxes would define zero modes and are quantized closed 2-surfaces. The induced metric appearing in Kähler action is however not zero mode degree of freedom. If the action contains volume term, the assumption about union of symmetric spaces is not well-motivated.

Symplectic transformations are not the only candidates for the isometries of WCW. The basic picture about what these maximal isometries could be, is partially inspired by string models.

1. A weaker proposal is that the symplectomorphisms of  $H$  define only symplectomorphisms of WCW. Extended conformal symmetries define also a candidate for isometry group. Remarkably, light-like boundary has an infinite-dimensional group of isometries which are in 1-1 correspondence with conformal symmetries of  $S^2 \subset S^2 \times R_+ = \delta M_+^4$ .
2. Extended Kac Moody symmetries induced by isometries of  $\delta M_+^4$  are also natural candidates for isometries. The motivation for the proposal comes from physical intuition deriving from string models. Note they do not include Poincare symmetries, which act naturally as isometries in the moduli space of causal diamonds (CDs) forming the "spine" of WCW.
3. The light-like orbits of partonic 2-surfaces might allow separate symmetry algebras. One must however notice that there is exchange of charges between interior degrees of freedom and partonic 2-surfaces. The essential point is that one can assign to these surface conserved charges when the dual light-like coordinate defines time coordinate. This picture also assumes a slicing of space-time surface by the partonic orbits for which partonic orbits associated with wormhole throats and boundaries of the space-time surface would be special. This slicing would correspond to Hamilton-Jacobi structure.
4. Fractal hierarchy of symmetry algebras with conformal weights, which are non-negative integer multiples of fundamental conformal weights, is essential

and distinguishes TGD from string models. Gauge conditions are true only the isomorphic subalgebra and its commutator with the entire algebra and the maximal gauge symmetry to a dynamical symmetry with generators having conformal weights below maximal value. This view also conforms with p-adic mass calculations.

5. The realization of the symmetries for 3-surfaces at the boundaries of CD and for light-like orbits of partonic 2-surfaces is known. The problem is how to extend the symmetries to the interior of the space-time surface. It is natural to expect that the symmetries at partonic orbits and light-cone boundary extend to the same symmetries.

### 3.2.2 Realization Of Super-Conformal Symmetries

The detailed realization of various super-conformal symmetries has been also a long standing problem.

1. Super-conformal symmetry requires that Dirac action for string world sheets is accompanied by string world sheet area as part of bosonic action. String world sheets are implied and can be present only in Minkowskian regions if one demands that octonionic and ordinary representations of induced spinor structure are equivalent (this requires vanishing of induced spinor curvature to achieve associativity in turn implying that  $CP_2$  projection is 1-D). Note that 1-dimensionality of  $CP_2$  projection is symplectically invariant property. Kähler action is not invariant under symplectic transformations. This is necessary for having non-trivial Kähler metric. Whether WCW really possesses super-symplectic isometries remains an open problem.
2. Super-conformal symmetry also demands that Kähler action is accompanied by what I call Kähler-Dirac action with gamma matrices defined by the contractions of the canonical momentum currents with embedding space-gamma matrices. Both the well-definedness of em charge and equivalence of octonionic spinor dynamics with ordinary one require the restriction of spinor modes to string world sheets with light-like boundaries at wormhole throats. K-D action with the localization of induced spinors at string world sheets is certainly the minimal option to consider.
3. Strong form of holography suggested by strong form of general coordinate invariance strongly suggests that super-conformal symmetry is broken gauge invariance in the sense that the classical super-conformal charges for a sub-algebra of the symplectic algebra with conformal weights vanishing modulo some integer  $n$  vanish. The proposal is that  $n$  corresponds to the effective Planck constant as  $h_{eff}/h = n$ . The standard conformal symmetries for spinors modes at string world sheets is always unbroken gauge symmetry.

### 3.2.3 The conserved charges associated with holomorphies

Generalized holomorphy not only solves explicitly the equations of motion [?] but, as found quite recently, also gives corresponding conserved Noether currents and charges.



1. Generalized holomorphy algebra generalizes the Super-Virasoro algebra and the Super-Kac-Moody algebra related to the conformal invariance of the string model. The corresponding Noether charges are conserved. Modified Dirac action allows to construct the supercharges having interpretation as WCW gamma matrices. This suggests an answer to a longstanding question related to the isometries of the "world of the classical worlds" (WCW).
2. Either the generalized holomorphies or the symplectic symmetries of  $H = M^4 \times CP_2$  or both together define WCW isometries and corresponding super algebra. It would seem that symplectic symmetries induced from  $H$  are not necessarily needed and might correspond to symplectic symmetries of WCW. One would obtain a close similarity with the string model, except that one has half-algebra for which conformal weights are proportional to non-negative integers and gauge conditions only apply to an isomorphic subalgebra. These are labeled by positive integers and one obtains a hierarchy.
3. By their light-likeness, the light cone boundary and orbits of partonic 2-surfaces allow an infinite-dimensional isometry group. This is possible only in dimension four. Its transformations are generalized conformal transformations of 2-sphere (partonic 2-surface) depending on light-like radial coordinate such that the radial scaling compensates for the usual conformal scaling of the metric. The WCW isometries would thus correspond to the isometries of the parton orbit and of the boundary of the light cone! These two representations could provide alternative representations for the charges if the strong form of holography holds true and would realize a strong form of holography. Perhaps these realizations deserve to be called inertial and gravitational charges.

Can these transformations leave the action invariant? For the light-cone boundary, this looks obvious if the light-cone is sliced by a surface parallel to the light-cone boundary. Note however that the tip of this surface might produce problems. A slicing defined by the Hamilton-Jacobi structure would be naturally associated with partonic orbits.

4. What about Poincare symmetries? They would act on the center of mass coordinates of causal diamonds (CDs) as found already earlier [?]. CDs form the "spine" of WCW, which can be regarded as fiber space with fiber for a given CD containing as a fiber the space-time surfaces inside it.

The super-symmetric counterparts of holomorphic charges for the modified Dirac action and bilinear in fermionic oscillator operators associated with the second quantization of free spinor fields in  $H$ , define gamma matrices of WCW. Their anticommutators define the Kähler metric of WCW. There is no need to calculate either the action defining the classical Kähler action defining the Kähler function or its derivatives with respect to WCW complex coordinates and their conjugates. What is important is that this makes it possible to speak about WCW metric also for number theoretical discretization of WCW with space-time surfaces replaced with their number theoretic discretizations.

### 3.2.4 Could generalized holomorphy allow to sharpen the existing views?

This picture is rather speculative, allows several variants, and is not proven. There is now however a rather convincing ansatz for the general form of preferred extremals. This proposal relies on the realization of holography as generalized 4-D holomorphy. Could it help to make the picture more precise?

1. Explicit solution of field equations in terms of the generalized holomorphy is now known. The solution ansatz is independent of action as long it is general coordinate invariance depending only on the induced geometric structures. Space-time surfaces would be minimal surfaces apart from lower-dimensional singular surfaces at which the field equations involve the entire action. Only the singularities, classical charges and positions of topological interaction vertices depend on the choice of the action [?]. Kähler action plus volume term is the choice of action forced by twistor lift making the choice of  $H$  unique.
2. The universality has a very intriguing implication. One can assign to any action of this kind conserved Noether currents and their fermionic counterparts (also super counterparts). One would have a huge algebra of conserved currents characterizing the space-time geometry. The corresponding charges can be made conserved by suitably modifying the form of holomorphic functions of the ansatz and therefore the time derivatives  $\partial_t h^k$  at the 3-D end of space-time surface at the boundary CD. This need not be the case for all deformations of partonic orbits. In any case, the 3-D holographic data seem to be dual as the strong form of holography suggests. The discussion of the symplectic symmetries leads to the conclusion that they give rise to conserved charges at the partonic 3-surfaces obeying Chern-Simons-Kähler dynamics, which is non-deterministic.
3. Hamilton-Jacobi structures emerge naturally as generalized conformal structures of space-time surfaces and  $M^4$  [?]. This inspires a proposal for a generalization of modular invariance and of moduli spaces as subspaces of Teichmüller spaces.
4. One can assign to holomorphy conserved Noether charges. The conservation reduces to the algebraic conditions satisfied for the same reason as field equations, i.e. the conservation conditions involving contractions of complex tensors of type (1,1) with tensors of type (2,0) and (0,2). The charges have the same form as Noether charges but it is not completely clear whether the action remains invariant under these transformations. This point is non-trivial since Noether theorem says that invariance of the action implies the existence of conserved charges but not vice versa. Could TGD represent a situation in which the equivalence between symmetries of action and conservation laws fails?

Also string models have conformal symmetries but in this case 2-D area form suffers conformal scaling. Also the fact that holomorphic ansatz is satisfied for such a large class of actions apart from singularities suggests that the action is not invariant.

5. The action should define Kähler function for WCW identified as the space of Bohr orbits. WCW Kähler metric is defined in terms of the second derivatives of the Kähler action of type (1,1) with respect to complex coordinates of WCW. Does the invariance of the action under holomorphies imply a trivial Kähler metric and constant Kähler function?

Here one must be very cautious since by holography the variations of the space-time surface are induced by those of 3-surface defining holographic data so that the entire space-time surface is modified and the action can change. The presence of singularities, analogous to poles and cuts of an analytic function and representing particles, suggests that the action represents the interactions of particles and must change. Therefore the action might not be invariant under holomorphies. The parameters characterizing the singularities should affect the value of the action just as the positions of these singularities in 2-D electrostatics affect the Coulomb energy.

Generalized conformal charges and supercharges define a generalization of Super Virasoro algebra of string models. Also Kac-Moody algebra assignable to the isometries of  $\delta M_+^4 \times CP_2$  and light  $H$  generalizes trivially.

6. An absolutely essential point is that generalized holomorphisms are *not* symmetries of Kähler function since otherwise Kähler metric involving second derivatives of type (1,1) with respect to complex coordinates of WCW is non-trivial if defined by these symmetry generators as differential operators. If Kähler function is equal to Kähler action, as it seems, Kähler action cannot be invariant under generalized holomorphies.

Noether's theorem states that the invariance of the action under a symmetry implies the conservation of corresponding charge but does *not* claim that the existence of conserved Noether currents implies invariance of the action. Since Noether currents are conserved now, one would have a concrete example about the situation in which the inverse of Noether's theorem does not hold true. In a string model based on area action, conformal transformations of complex string coordinates give rise to conserved Noether currents as one easily checks. The area element defined by the induced metric suffers a conformal scaling so that the action is not invariant in this case.

### 3.2.5 Challenging the existing picture of WCW geometry

These findings make it possible to challenge and perhaps sharpen the existing speculations concerning the metric and isometries of WCW.

I have considered the possibility that also the symplectomorphisms of  $\delta M^4 + \times CP_2$  could define WCW isometries. This actually the original proposal. One can imagine two options.

1. The continuation of symplectic transformations to transformations of the space-time surface from the boundary of light-cone or from the orbits partonic 2-surfaces should give rise to conserved Noether currents but it is not at all obvious whether this is the case.

2. One can assign conserved charges to the time evolution of the 3-D boundary data defining the holographic data: the time coordinate for the evolution would correspond to the light-like coordinate of light-cone boundary or partonic orbit. This option I have not considered hitherto. It turns out that this option works!

The conclusion would be that generalized holomorphies give rise to conserved charges for 4-D time evolution and symplectic transformations give rise to conserved charges for 3-D time evolution associated with the holographic data.

### 3.2.6 About extremals of Chern-Simons-Kähler action

Let us look first the general nature of the solutions to the extremization of Chern-Simons-Kähler action.

1. The light-likeness of the partonic orbits requires Chern-Simons action, which is equivalent to the topological action  $J \wedge J$ , which is total divergence and is a symplectic invariant. The field equations at the boundary cannot involve induced metric so that only induced symplectic structure remains. The 3-D holographic data at partonic orbits would extremize Chern-Simons-Kähler action. Note that at the ends of the space-time surface about boundaries of CD one cannot pose any dynamics.
2. If the induced Kähler form has only the  $CP_2$  part, the variation of Chern-Simons-Kähler form would give equations satisfied if the  $CP_2$  projection is at most 2-dimensional and Chern-Simons action would vanish and imply that instanton number vanishes.
3. If the action is the sum of  $M^4$  and  $CP_2$  parts, the field equations in  $M^4$  and  $CP_2$  degrees of freedom would give the same result. If the induced Kähler form is identified as the sum of the  $M^4$  and  $CP_2$  parts, the equations also allow solutions for which the induced  $M^4$  and  $CP_2$  Kähler forms sum up to zero. This phase would involve a map identifying  $M^4$  and  $CP_2$  projections and force induce Kähler forms to be identical. This would force magnetic charge in  $M^4$  and the question is whether the line connecting the tips of the CD makes non-trivial homology possible. The homology charges and the 2-D ends of the partonic orbit cancel each other so that partonic surfaces can have monopole charge.

The conditions at the partonic orbits do not pose conditions on the interior and should allow generalized holomorphy. The following considerations show that besides homology charges as Kähler magnetic fluxes also Hamiltonian fluxes are conserved in Chern-Simons-Kähler dynamics.

### 3.2.7 Can one assign conserved charges with symplectic transformations or partonic orbits and 3-surfaces at light-cone boundary?

The geometric picture is that symplectic symmetries are Hamiltonian flows along the light-like partonic orbits generated by the projection  $A_t$  of the Kähler gauge potential in the direction of the light-like time coordinate. The physical picture

is that the partonic 2-surface is a Kähler charged particle that couples to the Hamilton  $H = A_t$ . The Hamiltonians  $H_A$  are conserved in this time evolution and give rise to conserved Noether currents. The corresponding conserved charge is integral over the 2-surface defined by the area form defined by the induced Kähler form.

Let's examine the change of the Chern-Simons-Kähler action in a deformation that corresponds, for example, to the  $CP_2$  symplectic transformation generated by Hamilton  $H_A$ .  $M^4$  symplectic transformations can be treated in the same way: here however  $M^4$  Kähler form would be involved, assumed to accompany Hamilton-Jacobi structure as a dynamically generated structure.

1. Instanton density for the induced Kähler form reduces to a total divergence and gives Chern-Simons-Kähler action, which is TGD analog of topological action. This action should change in infinitesimal symplectic transformations by a total divergence, which should vanish for extremals and give rise to a conserved current. The integral of the divergence gives a vanishing charge difference between the ends of the partonic orbit. If the symplectic transformations define symmetries, it should be possible to assign to each Hamiltonian  $H_A$  a conserved charge. The corresponding quantal charge would be associated with the modified Dirac action.
2. The conserved charge would be an integral over  $X^2$ . The surface element is not given by the metric but by the symplectic structure, so that it is preserved in symplectic transformations. The 2-surface of the time evolution should correspond to the Hamiltonian time transformation generated by the projection  $A_\alpha = A_k \partial_\alpha s^k$  of the Kähler gauge potential  $A_k$  to the direction of light-like time coordinate  $x^\alpha \equiv t$ .
3. The effect of the generator  $j_A^k = J^{kl} \partial_l H_A$  on the Kähler potential  $A_l$  is given by  $j_A^k \partial_k A_l$ . This can be written as  $\partial_k A_l = J_{kl} + \partial_l A_k$ . The first term gives the desired total divergence  $\partial_\alpha (\epsilon^{\alpha\beta\gamma} J_{\beta\gamma} H_A)$ .

The second term is proportional to the term  $\partial_\alpha H_A - \{A_\alpha, H\}$ . Suppose that the induced Kähler form is transversal to the light-like time coordinate  $t$ , i.e. the induced Kähler form does not have components of form  $J_{t\mu}$ . In this kind of situation the only possible choice for  $\alpha$  corresponds to the time coordinate  $t$ . In this situation one can perform the replacement  $\partial_\alpha H_A - \{A_\alpha, H\} \rightarrow dH_A/dt - \{A_t, H\}$ . This corresponds to a Hamiltonian time evolution generated by the projection  $A_t$  acting as a Hamiltonian. If this is really a Hamiltonian time evolution, one has  $dH_A/dt - \{A, H\} = 0$ . Because the Poisson bracket represents a commutator, the Hamiltonian time evolution equation is analogous to the vanishing of a covariant derivative of  $H_A$  along light-like curves:  $\partial_t H_A + [A, H_A] = 0$ . The physical interpretation is that the partonic surface develops like a particle with a Kähler charge. As a consequence the change of the action reduces to a total divergence.

An explicit expression for the conserved current  $J_A^\alpha = H_A \epsilon^{\alpha\beta\gamma} J_{\beta\gamma}$  can be derived from the vanishing of the total divergence. Symplectic transformations on  $X^2$  generate an infinite-dimensional symplectic algebra. The charge is given by the Hamiltonian flux  $Q_A = \int H_A J_{\beta\gamma} dx^\alpha \wedge dx^\beta$ .

4. If the projection of the partonic path  $CP_2$  or  $M^4$  is 2-D, then the light-like geodesic line corresponds to the path of the parton surface. If  $A_t$  can be chosen parallel to the surface, its projection in the direction of time disappears and one has  $A_t = 0$ . In the more general case,  $X^2$  could, for example, rotate in  $CP_2$ . In this case  $A_t$  is nonvanishing. If  $J$  is transversal (no Kähler electric field), charge conservation is obtained.

Do the above observations apply at the boundary of the light-cone?

1. Now the 3-surface is space-like and Chern-Simons-Kähler action makes sense. It is not necessary but emerges from the "instanton density" for the Kähler form. The symplectic transformations of  $\delta M_+^4 \times CP_2$  are the symmetries. The most time evolution associated with the radial light-like coordinate would be from the tip of the light-cone boundary to the boundary of CD. Conserved charges as homological invariants defining symplectic algebra would be associated with the 2-D slices of 3-surfaces. For closed 3-surfaces the total charges from the sheets of 3-space as covering of  $\delta M_+^4$  must sum up to zero.
2. Interestingly, the original proposal [K7] for the isometries of WCW was that the Hamiltonian fluxes assignable to  $M^4$  and  $CP_2$  degrees of freedom at light-like boundary act define the charges associated with the WCW isometries as symplectic transformations so that a strong form of holography would have been realized and space-time surface would have been effectively 2-dimensional. The recent view is that these symmetries pose conditions only on the 3-D holographic data. The holographic charges would correspond to additional isometries of WCW and would be well-defined for the 3-surfaces at the light-cone boundary.

To sum up, one can imagine many options but the following picture is perhaps the simplest one and is supported by physical intuition and mathematical facts. The isometry algebra of  $\delta M_+^4 \times CP_2$  consists of generalized conformal and KM algebras at 3-surfaces in  $\delta M_+^4 \times CP_2$  and symplectic algebras at the light cone boundary and 3-D light-like partonic orbits. The latter symmetries give constraints on the 3-D holographic data. It is still unclear whether one can assign generalized conformal and Kac-Moody charges to Chern-Simons-Kähler action. The isomorphic subalgebras labelled by a positive integer and their commutators with the entire algebra would annihilate the physical states. The isomorphic subalgebras labelled by a positive integer and their commutators with the entire algebra would annihilate the physical states. These two representations would generalize the notions of inertial and gravitational mass and their equivalence would generalize the Equivalence Principle.

### 3.2.8 Objection against the idea about theoretician friendly Mother Nature

One of the key ideas behind the TGD view of dark matter is that Nature is theoretician friendly [?]. When the coupling strength proportional to  $\hbar_{eff}$  becomes so large that perturbation series ceases to converge, a phase transition increasing the value of  $\hbar_{eff}$  takes place so that the perturbation series converges.

One can however argue that this argument is quantum field-theoretic and does not apply in TGD since holography changes the very concept of perturbation theory. There is no path integral to worry about. Path integral is indeed such a fundamental concept that one expects it to have some approximate counterpart also in the TGD Universe. Bohr orbits are not completely deterministic: could the sum over the Bohr orbits however translate to an approximate description as a path integral at the QFT limit? The dynamics of light-like partonic orbits is indeed non-deterministic and could give rise to an analog of path integral as a finite sum.

1. The dynamics implied by Chern-Simons-Kähler action assignable to the partonic 3-surface with light-like coordinate in the role of time, is very topological in that the partonic orbits is light-like 3-surface and has 2-D  $CP_2$  and  $M^4$  projections unless the induced  $M^4$  and  $CP_2$  Kähler forms sum up to zero. The light-likeness of the projection is a very loose condition and the sum over partonic orbits as possible representation of holographic data analogous to initial values (light-likeness!) is therefore analogous to the sum over all paths appearing as a representation of Schrödinger equation in wave mechanics.

One would have an analog of 1-D QFT. This means that the infinities of quantum field theories are absent but for a large enough coupling strength  $g^2/4\pi\hbar$  the perturbation series fails to converge. The increase of  $h_{eff}$  would resolve the problem. For instance, Dirac equation in atomic physics makes unphysical predictions when the value of nuclear charge is larger than  $Z \sim 137$ .

2. I have also considered a discrete variant of this picture motivated by the fact that the presence of the volume term in the action implies that the  $M^4$  projection of the  $CP_2$  type extremal is a light-like geodesic line. The light-like orbits would consist of pieces of light-like geodesics implying that the average velocity would be smaller than  $c$ : this could be seen as a correlate for massivation.

The points at which the direction of segment changes would correspond to points at which energy and momentum transfer between the partonic orbit and environment takes place. This kind of quantum number transfer might occur at least for the fermionic lines as boundaries of string world sheets. They could be described quantum mechanically as interactions with classical fields in the same way as the creation of fermion pairs as a fundamental vertex [?]. The same universal 2-vertex would be in question.

At these points the minimal surface property would fail and the trace of the second fundamental form would not vanish but would have a delta function-like singularity. The  $CP_2$  part of the second fundamental form has quantum numbers of Higgs so that there would be an analogy with the standard description of massivation by the Higgs mechanism. Higgs would be only where the vertices are.

3. What is intriguing, that the light-likeness of the projection of the  $CP_2$  type extremals in  $M^4$  leads to Virasoro conditions assignable to  $M^4$  coordinates and this eventually led to the idea of conformal symmetries as isometries as WCW. In the case of the partonic orbits, the light-like curve would be in  $M^4 \times CP_2$

but it would not be surprising if the generalization of the Virasoro conditions would emerge also now.

One can write  $M^4$  and  $CP_2$  coordinates for the light-like curve as Fourier expansion in powers of  $\exp(it)$ , where  $t$  is the light-like coordinate. This gives  $h^k = \sum h_n^k \exp(int)$ . If the  $CP_2$  projection of the orbits of the partonic 2-surface is geodesic circle,  $CP_2$  metric  $s_{kl}$  is constant, the light-likeness condition  $h_{kl} \partial_t h^k \partial_t h^l = 0$  gives  $Re[h_{kl} \sum_m h_{n-m}^k \bar{h}_m^l] = 0$ . This does not give Virasoro conditions.

The condition  $d/dt(h_{kl} \partial_t h^k \partial_t h^l = 0) = 0$  however gives the standard Virasoro condition in quantization condition stating that the operator counterparts of quantities  $L_n = Re[h_{kl} \sum_m (n-m) h_{n-m}^k \bar{h}_m^l]$  annihilate the physical states. What is interesting is that the latter condition also allows time-like (and even space-like) geodesics.

Could massivation mean a failure of light-likeness? For piecewise light-like geodesics the light-likeness condition would be true only inside the segments. By taking Fourier transform one expects to obtain Virasoro conditions with a cutoff analogous to the momentum cutoff in condensed matter physics for crystals.

4. In TGD the Virasoro, Kac-Moody algebras and symplectic algebras are replaced by half-algebras and the gauge conditions are satisfied for conformal weights which are  $n$ -multiples of fundamentals with  $n$  larger than some minimal value. This would dramatically reduce the effects of the non-determinism and could make the sum over all paths allowed by the light-likeness manifestly finite and reduce it to a sum with a finite number of terms. This cutoff in degrees of freedom would correspond to a genuinely physical cutoff due to the finite measurement resolution coded to the number theoretical anatomy of the space-time surfaces. This cutoff is analogous to momentum cutoff and could at the space-time picture correspond to finite minimum length for the light-like segments of the orbit of the partonic 2-surface.

### 3.2.9 Boundary conditions at partonic orbits and holography

TGD reduces coupling constant evolution to a number theoretical evolution of the coupling parameters of the action identified as Kähler function for WCW. An interesting question is how the 3-D holographic data at the partonic orbits relates to the corresponding 3-D data at the ends of space-time surfaces at the boundary of CD, and how it relates to coupling constant evolution.

1. The twistor lift of TGD strongly favours 6-D Kähler action, which dimensionally reduces to Kähler action plus volume term plus topological  $\int J \wedge J$  term reducing to Chern Simons-Kähler action. The coefficients of these terms are proposed to be expressible in terms of number theoretical invariants characterizing the algebraic extensions of rationals and polynomials determining the space-time surfaces by  $M^8 - H$  duality.

Number theoretical coupling constant evolution would be discrete. Each extension of rationals would give rise to its own coupling parameters involving



also the ramified primes characterizing the polynomials involved and identified as p-adic length scales.

2. The time evolution of the partonic orbit would be non-deterministic but subject to the light-likeness constraint and boundary conditions guaranteeing conservation laws. The natural expectation is that the boundary/interface conditions for a given action cannot be satisfied for all partonic orbits (and other singularities). The deformation of the partonic orbit requiring that boundary conditions are satisfied, does not affect  $X^3$  but the time derivatives  $\partial_t h^k$  at  $X^3$  are affected since the form of the holomorphic functions defining the space-time surface would change. The interpretation would be in terms of duality of the holographic data associated with the partonic orbits *resp.*  $X^3$ .

There can of course exist deformations, which require the change of the coupling parameters of the action to satisfy the boundary conditions. One can consider an analog of renormalization group equations in which the deformation corresponds to a modification of the coupling parameters of the action, most plausibly determined by the twistor lift. Coupling parameters would label different regions of WCW and the space-time surfaces possible for two different sets of coupling parameters would define interfaces between these regions.

In order to build a more detailed view one must fix the details related to the action whose value defines the WCW Kähler function.

1. If Kähler action is identified as Kähler action, the identification is unique. There is however the possibility that the imaginary exponent of the instanton term or the contribution from the Euclidean region is not included in the definition of Kähler function. For instance instanton term could be interpreted as a phase of quantum state and would not contribute.
2. Both Minkowskian and Euclidean regions are involved and the Euclidean signature poses problems. The definition of the determinant as  $\sqrt{-g_4}$  is natural in Minkowskian regions but gives an imaginary contribution in Euclidean regions.  $\sqrt{|g_4|}$  is real in both regions.  $i\sqrt{g_4}$  is real in Minkowskian regions but imaginary in the Euclidean regions.

There is also a problem related to the instanton term, which does not depend on the metric determinant at all. In QFT context the instanton term is imaginary and this is important for instance in QCD in the definition of CP breaking vacuum functional. Should one include only the 4-D or possibly only Minkowskian contribution to the Kähler function imaginary coefficient for the instanton/Euclidian term would be possible?

3. Boundary conditions guaranteeing the conservation laws at the partonic orbits must be satisfied. Consider the  $\sqrt{|g_4|}$  case. Charge transfer between Euclidean and Minkowskian regions. If the C-S-K term is real, also the charge transfer between partonic orbit and 4-D regions is possible. The boundary conditions at the partonic orbit fix it to a high degree and also affect the time

derivatives  $\partial_t h^k$  at  $X^3$ . This option looks physically rather attractive because classical conserved charges would be real.

If the C-S-K term is imaginary it behaves like a free particle since charge exchange with Minkowskian and Euclidean regions is not possible. A possible interpretation of the possible  $M^4$  contribution to momentum could be in terms of decay width. The symplectic charges do not however involve momentum. The imaginary contribution to momentum could therefore come only from the Euclidean region.

4. If the Euclidean contribution is imaginary, it seems that it cannot be included in the Kähler function. Since in  $M^8$  picture the momenta of virtual fermions are in general complex, one could consider the possibility that Euclidean contribution to the momentum is imaginary and allows an interpretation as a decay width.

### 3.2.10 The TGD counterparts of the gauge conditions of string models

The string model picture forces to ask whether the symplectic algebras and the generalized conformal and Kac-Moody algebras could act as gauge symmetries.

1. In string model picture conformal invariance would suggest that the generators of the generalized conformal and KM symmetries act as gauge transformations annihilate the physical states. In the TGD framework, this does not however make sense physically. This also suggests that the components of the metric defined by supergenerators of generalized conformal and Kac Moody transformations vanish. If so, the symplectomorphisms  $\delta M_+^4 \times CP_2$  localized with respect to the light-like radial coordinate acting as isometries would be needed. The half-algebras of both symplectic and conformal generators are labelled by a non-negative integer defining an analog of conformal weight so there is a fractal hierarchy of isomorphic subalgebras in both cases.
2. TGD forces to ask whether only subalgebras of both conformal and Kac-Moody half algebras, isomorphic to the full algebras, act as gauge algebras. This applies also to the symplectic case. Here it is essential that only the half algebra with non-negative multiples of the fundamental conformal weights is allowed. For the subalgebra annihilating the states the conformal weights would be fixed integer multiples of those for the full algebra. The gauge property would be true for all algebras involved. The remaining symmetries would be genuine dynamical symmetries of the reduced WCW and this would reflect the number theoretically realized finite measurement resolution. The reduction of degrees of freedom would also be analogous to the basic property of hyperfinite factors assumed to play a key role in the definition of finite measurement resolution.
3. For strong holography, the orbits of partonic 2-surfaces and boundaries of the spacetime surface at  $\delta M_+^4$  would be dual in the information theoretic sense. Either would be enough to determine the space-time surface.

### 3.3 Interior Dynamics For Fermions, The Role Of Vacuum Extremals, And Dark Matter

The key role of  $CP_2$ -type and  $M^4$ -type vacuum extremals has been rather obvious from the beginning but the detailed understanding has been lacking. Both kinds of extremals are invariant under symplectic transformations of  $\delta M^4 \times CP_2$ , which inspires the idea that they give rise to isometries of WCW. The deformations  $CP_2$ -type extremals correspond to lines of generalized Feynman diagrams.  $M^4$  type vacuum extremals in turn are excellent candidates for the building bricks of many-sheeted space-time giving rise to GRT space-time as approximation. For  $M^4$  type vacuum extremals  $CP_2$  projection is (at most 2-D) Lagrangian manifold so that the induced Kähler form vanishes and the action is fourth-order in small deformations. This implies the breakdown of the path integral approach and of canonical quantization, which led to the notion of WCW.

If the action in Minkowskian regions contains also string area, the situation changes dramatically since strings dominate the dynamics in excellent approximation and string theory should give an excellent description of the situation: this of course conforms with the dominance of gravitation.

String tension would be proportional to  $1/\hbar G$  and this raises a grave classical counter argument. In string model massless particles are regarded as strings, which have contracted to a point in excellent approximation and cannot have length longer than Planck length. How this can be consistent with the formation of gravitationally bound states is however not understood since the required non-perturbative formulation of string model required by the large valued of the coupling parameter  $GMm$  is not known.

In TGD framework strings would connect even objects with macroscopic distance and would obviously serve as correlates for the formation of bound states in quantum level description. The classical energy of string connecting say the two wormhole contacts defining elementary particle is gigantic for the ordinary value of  $\hbar$  so that something goes wrong.

I have however proposed [K20, K16, K17] that gravitons - at least those mediating interaction between dark matter have large value of Planck constant. I talk about gravitational Planck constant and one has  $\hbar_{eff} = \hbar_{gr} = GMm/v_0$ , where  $v_0/c < 1$  ( $v_0$  has dimensions of velocity). This makes possible perturbative approach to quantum gravity in the case of bound states having mass larger than Planck mass so that the parameter  $GMm$  analogous to coupling constant is very large. The velocity parameter  $v_0/c$  becomes the dimensionless coupling parameter. This reduces the string tension so that for string world sheets connecting macroscopic objects one would have  $T \propto v_0/G^2 Mm$ . For  $v_0 = GMm/\hbar$ , which remains below unity for  $Mm/m_{pl}^2$  one would have  $\hbar_{gr}/\hbar = 1$ . Hence action remains small and its imaginary exponent does not fluctuate wildly to make the bound state forming part of gravitational interaction short ranged. This is expected to hold true for ordinary matter in elementary particle scales. The objects with size scale of large neutron (100  $\mu\text{m}$  in the density of water) - probably not an accident - would have mass above Planck mass so that dark gravitons and also life would emerge as massive enough gravitational bound states are formed.  $\hbar_{gr} = \hbar_{eff}$  hypothesis is indeed central in TGD based view about living matter.

If one assumes that for non-standard values of Planck constant only  $n$ -multiples of super-conformal algebra in interior annihilate the physical states, interior conformal gauge degrees of freedom become partly dynamical. The identification of dark matter as macroscopic quantum phases labeled by  $h_{eff}/h = n$  conforms with this.

The emergence of dark matter corresponds to the emergence of interior dynamics via breaking of super-conformal symmetry. The induced spinor fields in the interior of flux tubes obeying Kähler Dirac action should be highly relevant for the understanding of dark matter. The assumption that dark particles have essentially same masses as ordinary particles suggests that dark fermions correspond to induced spinor fields at both string world sheets and in the space-time interior: the spinor fields in the interior would be responsible for the long range correlations characterizing  $h_{eff}/h = n$ . Magnetic flux tubes carrying dark matter are key entities in TGD inspired quantum biology. Massless extremals represent second class of  $M^4$  type non-vacuum extremals.

This view forces once again to ask whether space-time SUSY is present in TGD and how it is realized. With a motivation coming from the observation that the mass scales of particles and sparticles most naturally have the same p-adic mass scale as particles in TGD Universe I have proposed that sparticles might be dark in TGD sense. The above argument leads to ask whether the dark variants of particles correspond to states in which one has ordinary fermion at string world sheet and 4-D fermion in the space-time interior so that dark matter in TGD sense would almost by definition correspond to sparticles!

### 3.4 Classical Number Fields And Associativity And Commutativity As Fundamental Law Of Physics

The dimensions of classical number fields appear as dimensions of basic objects in quantum TGD. Embedding space has dimension 8, space-time has dimension 4, light-like 3-surfaces are orbits of 2-D partonic surfaces. If conformal QFT applies to 2-surfaces (this is questionable), one-dimensional structures would be the basic objects. The lowest level would correspond to discrete sets of points identifiable as intersections of real and p-adic space-time sheets. This suggests that besides p-adic number fields also classical number fields (reals, complex numbers, quaternions, octonions [A10]) are involved [K22] and the notion of geometry generalizes considerably. In the recent view about quantum TGD the dimensional hierarchy defined by classical number field indeed plays a key role.  $H = M^4 \times CP_2$  has a number theoretic interpretation and standard model symmetries can be understood number theoretically as symmetries of hyper-quaternionic planes of hyper-octonionic space.

The associativity condition  $A(BC) = (AB)C$  suggests itself as a fundamental physical law of both classical and quantum physics. Commutativity can be considered as an additional condition. In conformal field theories associativity condition indeed fixes the n-point functions of the theory. At the level of classical TGD space-time surfaces could be identified as maximal associative (hyper-quaternionic) sub-manifolds of the embedding space whose points contain a preferred hyper-complex plane  $M^2$  in their tangent space and the hierarchy finite fields-rationals-reals-complex numbers-quaternions-octonions could have direct quantum physical counterpart [K22]. This leads to the notion of number theoretic compactification

analogous to the dualities of M-theory: one can interpret space-time surfaces either as hyper-quaternionic 4-surfaces of  $M^8$  or as 4-surfaces in  $M^4 \times CP_2$ . As a matter fact, commutativity in number theoretic sense is a further natural condition and leads to the notion of number theoretic braid naturally as also to direct connection with super string models.

At the level of Kähler-Dirac action the identification of space-time surface as a hyper-quaternionic sub-manifold of  $H$  means that the modified gamma matrices of the space-time surface defined in terms of canonical momentum currents of Kähler action using octonionic representation for the gamma matrices of  $H$  span a hyper-quaternionic sub-space of hyper-octonions at each point of space-time surface (hyper-octonions are the subspace of complexified octonions for which imaginary units are octonionic imaginary units multiplied by commuting imaginary unit). Hyper-octonionic representation leads to a proposal for how to extend twistor program to TGD framework [K29, K23].

### 3.4.1 How to achieve associativity in the fermionic sector?

In the fermionic sector an additional complication emerges. The associativity of the tangent- or normal space of the space-time surface need not be enough to guarantee the associativity at the level of Kähler-Dirac or Dirac equation. The reason is the presence of spinor connection. A possible cure could be the vanishing of the components of spinor connection for two conjugates of quaternionic coordinates combined with holomorphy of the modes.

1. The induced spinor connection involves sigma matrices in  $CP_2$  degrees of freedom, which for the octonionic representation of gamma matrices are proportional to octonion units in Minkowski degrees of freedom. This corresponds to a reduction of tangent space group  $SO(1,7)$  to  $G_2$ . Therefore octonionic Dirac equation identifying Dirac spinors as complexified octonions can lead to non-associativity even when space-time surface is associative or co-associative.
2. The simplest manner to overcome these problems is to assume that spinors are localized at 2-D string world sheets with 1-D  $CP_2$  projection and thus possible only in Minkowskian regions. Induced gauge fields would vanish. String world sheets would be minimal surfaces in  $M^4 \times D^1 \subset M^4 \times CP_2$  and the theory would simplify enormously. String area would give rise to an additional term in the action assigned to the Minkowskian space-time regions and for vacuum extremals one would have only strings in the first approximation, which conforms with the success of string models and with the intuitive view that vacuum extremals of Kähler action are basic building bricks of many-sheeted space-time. Note that string world sheets would be also symplectic covariants.

Without further conditions gauge potentials would be non-vanishing but one can hope that one can gauge transform them away in associative manner. If not, one can also consider the possibility that  $CP_2$  projection is geodesic circle  $S^1$ : symplectic invariance is considerably reduces for this option since symplectic transformations must reduce to rotations in  $S^1$ .

3. The first heavy objection is that action would contain Newton's constant  $G$  as a fundamental dynamical parameter: this is a standard recipe for building a non-renormalizable theory. The very idea of TGD indeed is that there is only single dimensionless parameter analogous to critical temperature. One can of course argue that the dimensionless parameter is  $\hbar G/R^2$ , "radius".

Second heavy objection is that the Euclidian variant of string action exponentially damps out all string world sheets with area larger than  $\hbar G$ . Note also that the classical energy of Minkowskian string would be gigantic unless the length of string is of order Planck length. For Minkowskian signature the exponent is oscillatory and one can argue that wild oscillations have the same effect.

The hierarchy of Planck constants would allow the replacement  $\hbar \rightarrow \hbar_{eff}$  but this is not enough. The area of typical string world sheet would scale as  $\hbar_{eff}$  and the size of CD and gravitational Compton lengths of gravitationally bound objects would scale as  $\sqrt{\hbar_{eff}}$  rather than  $\hbar_{eff} = GMm/v_0$ , which one wants. The only way out of problem is to assume  $T \propto (\hbar/\hbar_{eff})^2 \times (1/\hbar_{bar}G)$ . This is however un-natural for genuine area action. Hence it seems that the visit of the basic assumption of superstring theory to TGD remains very short.

#### 3.4.2 Is super-symmetrized Kähler-Dirac action enough?

Could one do without string area in the action and use only K-D action, which is in any case forced by the super-conformal symmetry? This option I have indeed considered hitherto. K-D Dirac equation indeed tends to reduce to a lower-dimensional one: for massless extremals the K-D operator is effectively 1-dimensional. For cosmic strings this reduction does not however take place. In any case, this leads to ask whether in some cases the solutions of Kähler-Dirac equation are localized at lower-dimensional surfaces of space-time surface.

1. The proposal has indeed been that string world sheets carry vanishing  $W$  and possibly even  $Z$  fields: in this manner the electromagnetic charge of spinor mode could be well-defined. The vanishing conditions force in the generic case 2-dimensionality.

Besides this the canonical momentum currents for Kähler action defining 4 embedding space vector fields must define an integrable distribution of two planes to give string world sheet. The four canonical momentum currents  $\Pi_k \alpha = \partial L_K / \partial_{\partial_\alpha h^k}$  identified as embedding 1-forms can have only two linearly independent components parallel to the string world sheet. Also the Frobenius conditions stating that the two 1-forms are proportional to gradients of two embedding space coordinates  $\Phi_i$  defining also coordinates at string world sheet, must be satisfied. These conditions are rather strong and are expected to select some discrete set of string world sheets.

2. To construct preferred extremal one should fix the partonic 2-surfaces, their light-like orbits defining boundaries of Euclidian and Minkowskian space-time regions, and string world sheets. At string world sheets the boundary condition would be that the normal components of canonical momentum currents for

Kähler action vanish. This picture brings in mind strong form of holography and this suggests that might make sense and also solution of Einstein equations with point like sources.

3. The localization of spinor modes at 2-D surfaces would follow from the well-definedness of em charge and one could have situation is which the localization does not occur. For instance, covariantly constant right-handed neutrinos spinor modes at cosmic strings are completely de-localized and one can wonder whether one could give up the localization inside wormhole contacts.
4. String tension is dynamical and physical intuition suggests that induced metric at string world sheet is replaced by the anti-commutator of the K-D gamma matrices and by conformal invariance only the conformal equivalence class of this metric would matter and it could be even equivalent with the induced metric. A possible interpretation is that the energy density of Kähler action has a singularity localized at the string world sheet.

Another interpretation that I proposed for years ago but gave up is that in spirit with the TGD analog of AdS/CFT duality the Noether charges for Kähler action can be reduced to integrals over string world sheet having interpretation as area in effective metric. In the case of magnetic flux tubes carrying monopole fluxes and containing a string connecting partonic 2-surfaces at its ends this interpretation would be very natural, and string tension would characterize the density of Kähler magnetic energy. String model with dynamical string tension would certainly be a good approximation and string tension would depend on scale of CD.

5. There is also an objection. For  $M^4$  type vacuum extremals one would not obtain any non-vacuum string world sheets carrying fermions but the successes of string model strongly suggest that string world sheets are there. String world sheets would represent a deformation of the vacuum extremal and far from string world sheets one would have vacuum extremal in an excellent approximation. Situation would be analogous to that in general relativity with point particles.
6. The hierarchy of conformal symmetry breakings for K-D action should make string tension proportional to  $1/h_{eff}^2$  with  $h_{eff} = h_{gr}$  giving correct gravitational Compton length  $\Lambda_{gr} = GM/v_0$  defining the minimal size of CD associated with the system. Why the effective string tension of string world sheet should behave like  $(\hbar/\hbar_{eff})^2$ ?

The first point to notice is that the effective metric  $G^{\alpha\beta}$  defined as  $h^{kl}\Pi_k^\alpha\Pi_l^\beta$ , where the canonical momentum current  $\Pi_k^\alpha = \partial L_K/\partial_{\partial_\alpha h^k}$  has dimension  $1/L^2$  as required. Kähler action density must be dimensionless and since the induced Kähler form is dimensionless the canonical momentum currents are proportional to  $1/\alpha_K$ .

Should one assume that  $\alpha_K$  is fundamental coupling strength fixed by quantum criticality to  $\alpha_K = 1/137$ ? Or should one regard  $g_K^2$  as fundamental parameter

so that one would have  $1/\alpha_K = \hbar_{eff}/4\pi g_K^2$  having spectrum coming as integer multiples (recall the analogy with inverse of critical temperature)?

The latter option is the in spirit with the original idea stating that the increase of  $\hbar_{eff}$  reduces the values of the gauge coupling strengths proportional to  $\alpha_K$  so that perturbation series converges (Universe is theoretician friendly). The non-perturbative states would be critical states. The non-determinism of Kähler action implying that the 3-surfaces at the boundaries of CD can be connected by large number of space-time sheets forming  $n$  conformal equivalence classes. The latter option would give  $G^{\alpha\beta} \propto \hbar_{eff}^2$  and  $\det(G) \propto 1/\hbar_{eff}^2$  as required.

7. It must be emphasized that the string tension has interpretation in terms of gravitational coupling on only at the GRT limit of TGD involving the replacement of many-sheeted space-time with single sheeted one. It can have also interpretation as hadronic string tension or effective string tension associated with magnetic flux tubes and telling the density of Kähler magnetic energy per unit length.

Superstring models would describe only the perturbative Planck scale dynamics for emission and absorption of  $\hbar_{eff}/h = 1$  on mass shell gravitons whereas the quantum description of bound states would require  $\hbar_{eff}/n > 1$  when the masses. Also the effective gravitational constant associated with the strings would differ from  $G$ .

The natural condition is that the size scale of string world sheet associated with the flux tube mediating gravitational binding is  $G(M + m)/v_0$ . By expressing string tension in the form  $1/T = n^2 \hbar G_1$ ,  $n = \hbar_{eff}/h$ , this condition gives  $\hbar G_1 = \hbar^2/M_{red}^2$ ,  $M_{red} = Mm/(M + m)$ . The effective Planck length defined by the effective Newton's constant  $G_1$  analogous to that appearing in string tension is just the Compton length associated with the reduced mass of the system and string tension equals to  $T = [v_0/G(M + m)]^2$  apart from a numerical constant ( $2G(M + m)$  is Schwarzschild radius for the entire system). Hence the macroscopic stringy description of gravitation in terms of string differs dramatically from the perturbative one. Note that one can also understand why in the Bohr orbit model of Nottale [?] for the planetary system and in its TGD version [K20]  $v_0$  must be by a factor 1/5 smaller for outer planets rather than inner planets.

### 3.4.3 Are 4-D spinor modes consistent with associativity?

The condition that octonionic spinors are equivalent with ordinary spinors looks rather natural but in the case of Kähler-Dirac action the non-associativity could leak in. One could of course give up the condition that octonionic and ordinary K-D equation are equivalent in 4-D case. If so, one could see K-D action as related to non-commutative and maybe even non-associative fermion dynamics. Suppose that one does not.

1. K-D action vanishes by K-D equation. Could this save from non-associativity? If the spinors are localized to string world sheets, one obtains just the standard stringy construction of conformal modes of spinor field. The induce



spinor connection would have only the holomorphic component  $A_z$ . Spinor mode would depend only on  $z$  but K-D gamma matrix  $\Gamma^z$  would annihilate the spinor mode so that K-D equation would be satisfied. There are good hopes that the octonionic variant of K-D equation is equivalent with that based on ordinary gamma matrices since quaternionic coordinated reduces to complex coordinate, octonionic quaternionic gamma matrices reduce to complex gamma matrices, sigma matrices are effectively absent by holomorphy.

2. One can consider also 4-D situation (maybe inside wormhole contacts). Could some form of quaternion holomorphy [A14] [K23] allow to realize the K-D equation just as in the case of super string models by replacing complex coordinate and its conjugate with quaternion and its 3 conjugates. Only two quaternion conjugates would appear in the spinor mode and the corresponding quaternionic gamma matrices would annihilate the spinor mode. It is essential that in a suitable gauge the spinor connection has non-vanishing components only for two quaternion conjugate coordinates. As a special case one would have a situation in which only one quaternion coordinate appears in the solution. Depending on the character of quaternion holomorphy the modes would be labelled by one or two integers identifiable as conformal weights.

Even if these octonionic 4-D modes exists (as one expects in the case of cosmic strings), it is far from clear whether the description in terms of them is equivalent with the description using K-D equation based ordinary gamma matrices. The algebraic structure however raises hopes about this. The quaternion coordinate can be represented as sum of two complex coordinates as  $q = z_1 + Jz_2$  and the dependence on two quaternion conjugates corresponds to the dependence on two complex coordinates  $z_1, z_2$ . The condition that two quaternion complexified gammas annihilate the spinors is equivalent with the corresponding condition for Dirac equation formulated using 2 complex coordinates. This for wormhole contacts. The possible generalization of this condition to Minkowskian regions would be in terms Hamilton-Jacobi structure.

Note that for cosmic strings of form  $X^2 \times Y^2 \subset M^4 \times CP_2$  the associativity condition for  $S^2$  sigma matrix and without assuming localization demands that the commutator of  $Y^2$  imaginary units is proportional to the imaginary unit assignable to  $X^2$  which however depends on point of  $X^2$ . This condition seems to imply correlation between  $Y^2$  and  $S^2$  which does not look physical.

To summarize, the minimal and mathematically most optimistic conclusion is that Kähler-Dirac action is indeed enough to understand gravitational binding without giving up the associativity of the fermionic dynamics. Conformal spinor dynamics would be associative if the spinor modes are localized at string world sheets with vanishing  $W$  (and maybe also  $Z$ ) fields guaranteeing well-definedness of em charge and carrying canonical momentum currents parallel to them. It is not quite clear whether string world sheets are present also inside wormhole contacts: for  $CP_2$  type vacuum extremals the Dirac equation would give only right-handed neutrino as a solution (could they give rise to  $N = 2$  SUSY?).

The construction of preferred extremals would realize strong form of holography. By conformal symmetry the effective metric at string world sheet could be conformally equivalent with the induced metric at string world sheets.

Dynamical string tension would be proportional to  $\hbar/h_{eff}^2$  due to the proportionality  $\alpha_K \propto 1/h_{eff}$  and predict correctly the size scales of gravitationally bound states for  $\hbar_{gr} = \hbar_{eff} = GMm/v_0$ . Gravitational constant would be a prediction of the theory and be expressible in terms of  $\alpha_K$  and  $R^2$  and  $\hbar_{eff}$  ( $G \propto R^2/g_K^2$ ).

In fact, all bound states - elementary particles as pairs of wormhole contacts, hadronic strings, nuclei [K12], molecules, etc. - are described in the same manner quantum mechanically. This is of course nothing new since magnetic flux tubes associated with the strings provide a universal model for interactions in TGD Universe. This also conforms with the TGD counterpart of AdS/CFT duality.

## 4 About some unclear issues of TGD

TGD has been in the middle of palace revolution during last two years and it is almost impossible to keep the chapters of the books updated. Adelic vision and twistor lift of TGD are the newest developments and there are still many details to be understood and errors to be corrected. The description of fermions in TGD framework has contained some unclear issues. Hence the motivation for the following brief comments.

There questions about the adelic vision about symmetries. Do the cognitive representations implying number theoretic discretization of the space-time surface lead to the breaking of the basic symmetries and are preferred embedding space coordinates actually necessary?

In the fermionic sector there are many questions deserving clarification. How quantum classical correspondence (QCC) is realized for fermions? How is SH realized for fermions and how does it lead to the reduction of dimension  $D = 4$  to  $D = 2$  (apart from number theoretical discretization)? Can scattering amplitudes be really formulated by using only the data at the boundaries of string sheets and what does this mean from the point of view of the modified Dirac equation? Are the spinors at light-like boundaries limiting values or sources? A long-standing issue concerns the fermionic anti-commutation relations: what motivated this article was the solution of this problem. There is also the general problem about whether statistical entanglement is “real”.

### 4.1 Adelic vision and symmetries

In the adelic TGD SH is weakened: also the points of the space-time surface having embedding space coordinates in an extension of rationals (cognitive representation) are needed so that data are not precisely 2-D. I have believed hitherto that one must use preferred coordinates for the embedding space  $H$  - a subset of these coordinates would define space-time coordinates. These coordinates are determined apart from isometries. Does the number theoretic discretization imply loss of general coordinate invariance and also other symmetries?

The reduction of symmetry groups to their subgroups (not only algebraic since

powers of  $e$  define finite-dimensional extension of  $p$ -adic numbers since  $e^p$  is ordinary  $p$ -adic number) is genuine loss of symmetry and reflects finite cognitive resolution. The physics itself has the symmetries of real physics.

The assumption about preferred embedding space coordinates is actually not necessary. Different choices of  $H$ -coordinates means only different and non-equivalent cognitive representations. Spherical and linear coordinates in finite accuracy do not provide equivalent representations.

## 4.2 Quantum-classical correspondence for fermions

Quantum-classical correspondence (QCC) for fermions is rather well-understood but deserves to be mentioned also here.

QCC for fermions means that the space-time surface as preferred extremal should depend on fermionic quantum numbers. This is indeed the case if one requires QCC in the sense that the fermionic representations of Noether charges in the Cartan algebras of symmetry algebras are equal to those to the classical Noether charges for preferred extremals.

Second aspect of QCC becomes visible in the representation of fermionic states as point like particles moving along the light-like curves at the light-like orbits of the partonic 2-surfaces (curve at the orbit can be locally only light-like or space-like). The number of fermions and antifermions dictates the number of string world sheets carrying the data needed to fix the preferred extremal by SH. The complexity of the space-time surface increases as the number of fermions increases.

## 4.3 Strong form of holography for fermions

It seems that scattering amplitudes can be formulated by assigning fermions with the boundaries of strings defining the lines of twistor diagrams [K9, K19]. This information theoretic dimensional reduction from  $D = 4$  to  $D = 2$  for the scattering amplitudes can be partially understood in terms of strong form of holography (SH): one can construct the theory by using the data at string worlds sheets and/or partonic 2-surfaces at the ends of the space-time surface at the opposite boundaries of causal diamond (CD).

4-D modified Dirac action would appear at fundamental level as supersymmetry demands but would be reduced for preferred extremals to its 2-D stringy variant serving as effective action. Also the value of the 4-D action determining the space-time dynamics would reduce to effective stringy action containing area term, 2-D Kähler action, and topological Kähler magnetic flux term. This reduction would be due to the huge gauge symmetries of preferred extremals. Sub-algebra of supersymplectic algebra with conformal weights coming as  $n$ -multiples of those for the entire algebra and the commutators of this algebra with the entire algebra would annihilate the physical states, and the corresponding classical Noether charges would vanish.

One still has the question why not the data at the entire string world sheets is not needed to construct scattering amplitudes. Scattering amplitudes of course need not code for the entire physics. QCC is indeed motivated by the fact that

quantum experiments are always interpreted in terms of classical physics, which in TGD framework reduces to that for space-time surface.

#### 4.4 The relationship between spinors in space-time interior and at boundaries between Euclidian and Minkoskian regions

Space-time surface decomposes to interiors of Minkowskian and Euclidian regions. At light-like 3-surfaces at which the four-metric changes, the 4-metric is degenerate. These metrically singular 3-surfaces - partonic orbits- carry the boundaries of string world sheets identified as carriers of fermionic quantum numbers. The boundaries define fermion lines in the twistor lift of TGD [K9, K19]. The relationship between fermions at the partonic orbits and interior of the space-time surface has however remained somewhat enigmatic.

So: What is the precise relationship between induced spinors  $\Psi_B$  at light-like partonic 3-surfaces and  $\Psi_I$  in the interior of Minkowskian and Euclidian regions? Same question can be made for the spinors  $\Psi_B$  at the boundaries of string world sheets and  $\Psi_I$  in interior of the string world sheets. There are two options to consider:

- Option I:  $\Psi_B$  is the limiting value of  $\Psi_I$  .
- Option II:  $\Psi_B$  serves as a source of  $\Psi_I$  .

For the Option I it is difficult to understand the preferred role of  $\Psi_B$ . I have considered Option II already years ago but have not been able to decide.

1. That scattering amplitudes could be formulated only in terms of sources only, would fit nicely with SH, twistorial amplitude construction, and also with the idea that scattering amplitudes in gauge theories can be formulated in terms of sources of boson fields assignable to vertices and propagators. Now the sources would become fermionic.
2. One can take gauge theory as a guideline. One adds to free Dirac equation source term  $\gamma^k A_k \Psi$ . Therefore the natural boundary term in the action would be of the form (forgetting overall scale factor)

$$S_B = \bar{\Psi}_I \Gamma^\alpha (C - S) A_\alpha \Psi_B + c.c .$$

Here the modified gamma matrix is  $\Gamma^\alpha (C - S)$  (contravariant form is natural for light-like 3-surfaces) is most naturally defined by the boundary part of the action - naturally Chern-Simons term for Kähler action.  $A$  denotes the Kähler gauge potential.

3. The variation with respect to  $\Psi_B$  gives

$$G^\alpha (C - S) A_\alpha \Psi_I = 0$$

at the boundary so that the C-S term and interaction term vanish. This does not however imply vanishing of the source term! This condition can be seen as a boundary condition.

The same argument applies also to string world sheets.

## 4.5 About second quantization of the induced spinor fields

The anti-commutation relations for the induced spinors have been a long-standing issue and during years I have considered several options. The solution of the problem looks however stupifyingly simple. The conserved fermion currents are accompanied by super-currents obtained by replacing  $\Psi$  with a mode of the induced spinor field to get  $\bar{u}_n \Gamma^\alpha \Psi$  or  $\bar{\Psi} \Gamma^\alpha u_n$  with the conjugate of the mode. One obtains infinite number of conserved super currents. One can also replace both  $\Psi$  and  $\bar{\Psi}$  in this manner to get purely bosonic conserved currents  $\bar{u}_m \Gamma^\alpha u_n$  to which one can assign a conserved bosonic charges  $Q_{mn}$ .

I noticed this years ago but did not realize that these bosonic charges define naturally anti-commutators of fermionic creation and annihilation operators! The ordinary anti-commutators of quantum field theory follow as a special case! By a suitable unitary transformation of the spinor basis one can diagonalize the hermitian matrix defined by  $Q_{mn}$  and by performing suitable scalings one can transform anti-commutation relations to the standard form. An interesting question is whether the diagonalization is needed, and whether the deviation of the diagonal elements from unity could have some meaning and possibly relate to the hierarchy  $h_{eff} = n \times h$  of Planck constants - probably not.

## 4.6 Is statistical entanglement “real” entanglement?

The question about the “reality” of statistical entanglement has bothered me for years. This entanglement is maximal and it cannot be reduced by measurement so that one can argue that it is not “real”. Quite recently I learned that there has been a longstanding debate about the statistical entanglement and that the issue still remains unresolved.

The idea that all electrons of the Universe are maximally entangled looks crazy. TGD provides several variants for solutions of this problem. It could be that only the fermionic oscillator operators at partonic 2-surfaces associated with the space-time surface (or its connected component) inside given CD anti-commute and the fermions are thus indistinguishable. The extremist option is that the fermionic oscillator operators belonging to a network of partonic 2-surfaces connected by string world sheets anti-commute: only the oscillator operators assignable to the same scattering diagram would anti-commute.

What about QCC in the case of entanglement. ER-EPR correspondence introduced by Maldacena and Susskind for 4 years ago proposes that blackholes (maybe even elementary particles) are connected by wormholes. In TGD the analogous statement emerged for more than decade ago - magnetic flux tubes take the role of wormholes in TGD. Magnetic flux tubes were assumed to be accompanied by string world sheets. I did not consider the question whether string world sheets are always accompanied by flux tubes.

What could be the criterion for entanglement to be “real”? “Reality” of entanglement demands some space-time correlate. Could the presence of the flux tubes make the entanglement “real”? If statistical entanglement is accompanied by string connections without magnetic flux tubes, it would not be “real”: only the presence of flux tubes would make it “real”. Or is the presence of strings enough to make the statistical entanglement “real”. In both cases the fermions associated with disjoint space-time surfaces or with disjoint CDs would not be indistinguishable. This looks rather sensible.

The space-time correlate for the reduction of entanglement would be the splitting of a flux tube and fermionic strings inside it. The fermionic strings associated with flux tubes carrying monopole flux are closed and the return flux comes back along parallel space-time sheet. Also fermionic string has similar structure. Reconnection of this flux tube with shape of very long flattened square splitting it to two pieces would be the correlate for the state function reduction reducing the entanglement with other fermions and would indeed decouple the fermion from the network.

## 5 About The Notion Of Four-Momentum In TGD Framework

The starting point of TGD was the energy problem of General Relativity [K25]. The solution of the problem was proposed in terms of sub-manifold gravity and based on the lifting of the isometries of space-time surface to those of  $M^4 \times CP_2$  in which space-times are realized as 4-surfaces so that Poincare transformations act on space-time surface as an 4-D analog of rigid body rather than moving points at space-time surface. It however turned out that the situation is not at all so simple.

There are several conceptual hurdles and I have considered several solutions for them. The basic source of problems has been Equivalence Principle (EP): what does EP mean in TGD framework [K25] ? A related problem has been the interpretation of gravitational and inertial masses, or more generally the corresponding 4-momenta. In General Relativity based cosmology gravitational mass is not conserved and this seems to be in conflict with the conservation of Noether charges. The resolution is in terms of zero energy ontology (ZEO), which however forces to modify slightly the original view about the action of Poincare transformations.

A further problem has been quantum classical correspondence (QCC): are quantal four-momenta associated with super conformal representations and classical four-momenta associated as Noether charges with Kähler action for preferred extremals identical? Could inertial-gravitational duality - that is EP - be actually equivalent with QCC? Or are EP and QCC independent dualities. A powerful experimental input comes p-adic mass calculations [K14] giving excellent predictions provided the number of tensor factors of super-Virasoro representations is five, and this input together with Occam’s razor strongly favors QCC=EP identification.

There is also the question about classical realization of EP and more generally, TGD-GRT correspondence.

Twistor Grassmannian approach has meant a technical revolution in quantum field theory (for attempts to understand and generalize the approach in TGD framework see [K23]). This approach seems to be extremely well suited to TGD and I have

considered a generalization of this approach from  $\mathcal{N} = 4$  SUSY to TGD framework by replacing point like particles with string world sheets in TGD sense and super-conformal algebra with its TGD version: the fundamental objects are now massless fermions which can be regarded as on mass shell particles also in internal lines (but with unphysical helicity). The approach solves old problems related to the realization of stringy amplitudes in TGD framework, and avoids some problems of twistorial QFT (IR divergences and the problems due to non-planar diagrams). The Yangian [A5] [?, ?, ?] variant of 4-D conformal symmetry is crucial for the approach in  $\mathcal{N} = 4$  SUSY, and implies the recently introduced notion of amplituhedron [?]. A Yangian generalization of various super-conformal algebras seems more or less a “must” in TGD framework. As a consequence, four-momentum is expected to have characteristic multilocal contributions identifiable as multipart on contributions now and possibly relevant for the understanding of bound states such as hadrons.

## 5.1 Scale Dependent Notion Of Four-Momentum In Zero Energy Ontology

Quite generally, General Relativity does not allow to identify four-momentum as Noether charges but in GRT based cosmology one can speak of non-conserved mass [K21], which seems to be in conflict with the conservation of four-momentum in TGD framework. The solution of the problem comes in terms of zero energy ontology (ZEO) [K1, K26], which transforms four-momentum to a scale dependent notion: to each causal diamond (CD) one can assign four-momentum assigned with say positive energy part of the quantum state defined as a quantum superposition of 4-surfaces inside CD.

ZEO is necessary also for the fusion of real and various p-adic physics to single coherent whole. ZEO also allows maximal “free will” in quantum jump since every zero energy state can be created from vacuum and at the same time allows consistency with the conservation laws. ZEO has rather dramatic implications: in particular the arrow of thermodynamical time is predicted to vary so that second law must be generalized. This has especially important implications in living matter, where this kind of variation is observed.

More precisely, this superposition corresponds to a spinor field in the “world of classical worlds” ( WCW ) [K26]: its components - WCW spinors - correspond to elements of fermionic Fock basis for a given 4-surface - or by holography implied by general coordinate invariance (GCI) - for 3-surface having components at both ends of CD. Strong form of GGI implies strong form of holography (SH) so that partonic 2-surfaces at the ends of space-time surface plus their 4-D tangent space data are enough to fix the quantum state. The classical dynamics in the interior is necessary for the translation of the outcomes of quantum measurements to the language of physics based on classical fields, which in turn is reduced to sub-manifold geometry in the extension of the geometrization program of physics provided by TGD.

Holography is very much reminiscent of QCC suggesting trinity: GCI-holography-QCC. Strong form of holography has strongly stringy flavor: string world sheets connecting the wormhole throats appearing as basic building bricks of particles emerge from the dynamics of induced spinor fields if one requires that the fermionic mode carries well-defined electromagnetic charge [K29].

## 5.2 Are The Classical And Quantal Four-Momenta Identical?

One key question concerns the classical and quantum counterparts of four-momentum. In TGD framework classical theory is an exact part of quantum theory. Classical four-momentum corresponds to Noether charge for preferred extremals of Kähler action. Quantal four-momentum in turn is assigned with the quantum superposition of space-time sheets assigned with CD - actually WCW spinor field analogous to ordinary spinor field carrying fermionic degrees of freedom as analogs of spin. Quantal four-momentum emerges just as it does in super string models - that is as a parameter associated with the representations of super-conformal algebras. The precise action of translations in the representation remains poorly specified. Note that quantal four-momentum does not emerge as Noether charge: at least it is not at all obvious that this could be the case.

Are these classical and quantal four-momenta identical as QCC would suggest? If so, the Noether four-momentum should be same for all space-time surfaces in the superposition. QCC suggests that also the classical correlation functions for various general coordinate invariant local quantities are same as corresponding quantal correlation functions and thus same for all 4-surfaces in quantum superposition - this at least in the measurement resolution used. This would be an extremely powerful constraint on the quantum states and to a high extend could determined the U-, M-, and S-matrices.

QCC seems to be more or less equivalent with SH stating that in some respects the descriptions based on classical physics defined by Kähler action in the interior of space-time surface and the quantal description in terms of quantum states assignable to the intersections of space-like 3-surfaces at the boundaries of CD and light-like 3-surfaces at which the signature of induced metric changes. SH means effective 2-dimensionality since the four-dimensional tangent space data at partonic 2-surfaces matters. SH could be interpreted as Kac-Mody and symplectic symmetries meaning that apart from central extension they act almost like gauge symmetries in the interiors of space-like 3-surfaces at the ends of CD and in the interiors of light-like 3-surfaces representing orbits of partonic 2-surfaces. Gauge conditions are replaced with Super Virasoro conditions. The word “almost” is of course extremely important.

## 5.3 What Equivalence Principle (EP) Means In Quantum TGD?

EP states the equivalence of gravitational and inertial masses in Newtonian theory. A possible generalization would be equivalence of gravitational and inertial four-momenta. In GRT this correspondence cannot be realized in mathematically rigorous manner since these notions are poorly defined and EP reduces to a purely local statement in terms of Einstein's equations.

What about TGD? What could EP mean in TGD framework?

1. Is EP realized at both quantum and space-time level? This option requires the identification of inertial and gravitational four-momenta at both quantum



and classical level. It is now clear that at classical level EP follows from very simple assumption that GRT space-time is obtained by lumping together the space-time sheets of the many-sheeted space-time and by the identification the effective metric as sum of  $M^4$  metric and deviations of the induced metrics of space-time sheets from  $M^2$  metric: the deviations indeed define the gravitational field defined by multiply topologically condensed test particle. Similar description applies to gauge fields. EP as expressed by Einstein's equations would follow from Poincare invariance at microscopic level defined by TGD space-time. The effective fields have as sources the energy momentum tensor and YM currents defined by topological inhomogenities smaller than the resolution scale.

2. QCC would require the identification of quantal and classical counterparts of both gravitational and inertial four-momenta. This would give three independent equivalences, say  $P_{I,class} = P_{I,quant}$ ,  $P_{gr,class} = P_{gr,quant}$ ,  $P_{gr,class} = P_{I,quant}$ , which imply the remaining ones.

Consider the condition  $P_{gr,class} = P_{I,class}$ . At classical level the condition that the standard energy momentum tensor associated with Kähler action has a vanishing divergence is guaranteed if Einstein's equations with cosmological term are satisfied. If preferred extremals satisfy this condition they are constant curvature spaces for non-vanishing cosmological constant. It must be emphasized that field equations are extremely non-linear and one must also consider preferred extremals (which could be identified in terms of space-time regions having so called Hamilton-Jacobi structure): hence these proposals are guesses motivated by what is known about exact solutions of field equations.

Consider next  $P_{gr,class} = P_{I,class}$ . At quantum level I have proposed coset representations for the pair of super conformal algebras  $g$  and  $h \subset g$  which correspond to the coset space decomposition of a given sector of WCW with constant values of zero modes. The coset construction would state that the differences of super-Virasoro generators associated with  $g$  resp.  $h$  annihilate physical states.

The identification of the algebras  $g$  and  $h$  is not straightforward. The algebra  $g$  could be formed by the direct sum of super-symplectic and super Kac-Moody algebras and its sub-algebra  $h$  for which the generators vanish at partonic 2-surface considered. This would correspond to the idea about WCW as a coset space  $G/H$  of corresponding groups (consider as a model  $CP_2 = SU(3)/U(2)$  with  $U(2)$  leaving preferred point invariant). The sub-algebra  $h$  in question includes or equals to the algebra of Kac-Moody generators vanishing at the partonic 2-surface. A natural choice for the preferred WCW point would be as maximum of Kähler function in Euclidian regions: positive definiteness of Kähler function allows only single maximum for fixed values of zero modes). Coset construction states that differences of super Virasoro generators associated with  $g$  and  $h$  annihilate physical states. This implies that corresponding four-momenta are identical that is Equivalence Principle.

3. Does EP at quantum level reduce to one aspect of QCC? This would require that classical Noether four-momentum identified as inertial momentum equals

to the quantal four-momentum assignable to the states of super-conformal representations and identifiable as gravitational four-momentum. There would be only one independent condition:  $P_{class} \equiv P_{I,class} = P_{gr,quant} \equiv P_{quant}$ .

Holography realized as AdS/CFT correspondence states the equivalence of descriptions in terms of gravitation realized in terms of strings in 10-D space-time and gauge fields at the boundary of AdS. What is disturbing is that this picture is not completely equivalent with the proposed one. In this case the super-conformal algebra would be direct sum of super-symplectic and super Kac-Moody parts.

Which of the options looks more plausible? The success of p-adic mass calculations [K14] have motivated the use of them as a guideline in attempts to understand TGD. The basic outcome was that elementary particle spectrum can be understood if Super Virasoro algebra has five tensor factors. Can one decide the fate of the two approaches to EP using this number as an input?

This is not the case. For both options the number of tensor factors is five as required. Four tensor factors come from Super Kac-Moody and correspond to translational Kac-Moody type degrees of freedom in  $M^4$ , to color degrees of freedom and to electroweak degrees of freedom ( $SU(2) \times U(1)$ ). One tensor factor comes from the symplectic degrees of freedom in  $\Delta CD \times CP_2$  (note that Hamiltonians include also products of  $\delta CD$  and  $CP_2$  Hamiltonians so that one does not have direct sum!).

The reduction of EP to the coset structure of WCW sectors is extremely beautiful property. But also the reduction of EP to QCC looks very nice and deep. It is of course possible that the two realizations of EP are equivalent and the natural conjecture is that this is the case.

For QCC option the GRT inspired interpretation of Equivalence Principle at space-time level remains to be understood. Is it needed at all? The condition that the energy momentum tensor of Kähler action has a vanishing divergence leads in General Relativity to Einstein equations with cosmological term. In TGD framework preferred extremals satisfying the analogs of Einstein's equations with several cosmological constant like parameters can be considered.

Should one give up this idea, which indeed might be wrong? Could the divergence of energy momentum tensor vanish only asymptotically as was the original proposal? Or should one try to generalize the interpretation? QCC states that quantum physics has classical correlate at space-time level and implies EP. Could also quantum classical correspondence itself have a correlate at space-time level. If so, space-time surface would be able to represent abstractions as statements about statements about.... as the many-sheeted structure and the vision about TGD physics as analog of Turing machine able to mimic any other Turing machine suggest.

## 5.4 TGD-GRT Correspondence And Equivalence Principle

One should also understand how General Relativity and EP emerge at classical level. The understanding comes from the realization that GRT is only an effective theory obtained by endowing  $M^4$  with effective metric.

1. The replacement of superposition of fields with superposition of their effects means replacing superposition of fields with the set-theoretic union of space-

time surfaces. Particle experiences sum of the effects caused by the classical fields at the space-time sheets (see **Fig.** <http://tgdtheory.fi/appfigures/fieldsuperpose.jpg> or **Fig. ??** in the appendix of this book).

2. This is true also for the classical gravitational field defined by the deviation from flat Minkowski metric in standard  $M^4$  coordinates for the space-time sheets. One can define effective metric as sum of  $M^4$  metric and deviations. This effective metric would correspond to that of General Relativity. This resolves long standing issues relating to the interpretation of TGD.
3. Einstein's equations could hold true for the effective metric. They are motivated by the underlying Poincare invariance which cannot be realized as global conservation laws for the effective metric. The conjecture vanishing of divergence of Kähler energy momentum tensor can be seen as the microscopic justification for the claim that Einstein's equations hold true for the effective space-time.
4. The breaking of Poincare invariance could have interpretation as effective breaking in zero energy ontology (ZEO), in which various conserved charges are length dependent and defined separately for each causal diamond (CD).

One can of course consider the possibility that Einstein's equations generalize for preferred extremals of Kähler action. This would actually represent at space-time level the notion of QCC rather than realise QCC interpreted as EP. The condition that the energy momentum tensor for Kähler action has vanishing covariant divergence would be satisfied in GRT if Einstein's equations with cosmological term hold true. This is the case also now but one can consider also more general solutions in which one has two cosmological constants which are not genuine constants anymore. It has turned out that this line of approach is too adhoc to be taken seriously.

## 5.5 How Translations Are Represented At The Level Of WCW ?

The four-momentum components appearing in the formulas of super conformal generators correspond to infinitesimal translations. In TGD framework one must be able to identify these infinitesimal translations precisely. As a matter of fact, finite measurement resolution implies that it is probably too much to assume infinitesimal translations. Rather, finite exponentials of translation generators are involved and translations are discretized. This does not have practical significance since for optimal resolution the discretization step is about  $CP_2$  length scale.

Where and how do these translations act at the level of WCW ? ZEO provides a possible answer to this question.

### 5.5.1 Discrete Lorentz transformations and time translations act in the space of CDs: inertial four-momentum

Quantum state corresponds also to wave function in moduli space of CDs. The moduli space is obtained from given CD by making all boosts for its non-fixed

boundary: boosts correspond to a discrete subgroup of Lorentz group and define a lattice-like structure at the hyperboloid for which proper time distance from the second tip of CD is fixed to  $T_n = n \times T(CP_2)$ . The quantization of cosmic redshift for which there is evidence, could relate to this lattice generalizing ordinary 3-D lattices from Euclidian to hyperbolic space by replacing translations with boosts (velocities).

The additional degree of freedom comes from the fact that the integer  $n > 0$  obtains all positive values. One has wave functions in the moduli space defined as a pile of these lattices defined at the hyperboloid with constant value of  $T(CP_2)$ : one can say that the points of this pile of lattices correspond to Lorentz boosts and scalings of CDs defining sub- WCW : s.

The interpretation in terms of group which is product of the group of shifts  $T_n(CP_2) \rightarrow T_{n+m}(CP_2)$  and discrete Lorentz boosts is natural. This group has same Cartesian product structure as Galilean group of Newtonian mechanics. This would give a discrete rest energy and by Lorentz boosts discrete set of four-momenta giving a contribution to the four-momentum appearing in the super-conformal representation.

What is important that each state function reduction would mean localisation of either boundary of CD (that is its tip). This localization is analogous to the localization of particle in position measurement in  $E^3$  but now discrete Lorentz boosts and discrete translations  $T_n \rightarrow T_{n+m}$  replace translations. Since the second end of CD is necessary del-ocalized in moduli space, one has kind of flip-flop: localization at second end implies de-localization at the second end. Could the localization of the second end (tip) of CD in moduli space correspond to our experience that momentum and position can be measured simultaneously? This apparent classicality would be an illusion made possible by ZEO.

The flip-flop character of state function reduction process implies also the alternation of the direction of the thermodynamical time: the asymmetry between the two ends of CDs would induce the quantum arrow of time. This picture also allows to understand what the experience growth of geometric time means in terms of CDs.

### 5.5.2 The action of translations at space-time sheets

The action of embedding space translations on space-time surfaces possibly becoming trivial at partonic 2-surfaces or reducing to action at  $\delta CD$  induces action on space-time sheet which becomes ordinary translation far enough from end end of space-time surface. The four-momentum in question is very naturally that associated with Kähler action and would therefore correspond to inertial momentum for  $P_{I,class} = P_{quant,gr}$  option. Indeed, one cannot assign quantal four-momentum to Kähler action as an operator since canonical quantization badly fails. In finite measurement infinitesimal translations are replaced with their exponentials for  $P_{I,class} = P_{quant,gr}$  option.

What looks like a problem is that ordinary translations in the general case lead out from given CD near its boundaries. In the interior one expects that the translation acts like ordinary translation. The Lie-algebra structure of Poincare algebra including sums of translation generators with positive coefficient for time translation is preserved if only time-like superpositions if generators are allowed also the

commutators of time-like translation generators with boost generators give time like translations. This defines a Lie-algebraic formulation for the arrow of geometric time. The action of time translation on preferred extremal would be ordinary translation plus continuation of the translated preferred extremal backwards in time to the boundary of CD. The transversal space-like translations could be made Kac-Moody algebra by multiplying them with functions which vanish at  $\delta CD$ .

A possible interpretation would be that  $P_{quant,gr}$  corresponds to the momentum assignable to the moduli degrees of freedom and  $P_{cl,I}$  to that assignable to the time like translations.  $P_{quant,gr} = P_{cl,I}$  would code for QCC. Geometrically quantum classical correspondence would state that time-like translation shift both the interior of space-time surface and second boundary of CD to the geometric future/past while keeping the second boundary of space-time surface and CD fixed.

## 5.6 Yangian And Four-Momentum

Yangian symmetry implies the marvellous results of twistor Grassmannian approach to  $\mathcal{N} = 4$  SUSY culminating in the notion of amplituhedron which promises to give a nice projective geometry interpretation for the scattering amplitudes [?]. Yangian symmetry is a multilocal generalization of ordinary symmetry based on the notion of co-product and implies that Lie algebra generates receive also multilocal contributions. I have discussed these topics from slightly different point of view in [K23], where also references to the work of pioneers can be found.

### 5.6.1 Yangian symmetry

The notion equivalent to that of Yangian was originally introduced by Faddeev and his group in the study of integrable systems. Yangians are Hopf algebras which can be assigned with Lie algebras as the deformations of their universal enveloping algebras. The elegant but rather cryptic looking definition is in terms of the modification of the relations for generating elements [K23]. Besides ordinary product in the enveloping algebra there is co-product  $\Delta$  which maps the elements of the enveloping algebra to its tensor product with itself. One can visualize product and co-product is in terms of particle reactions. Particle annihilation is analogous to annihilation of two particle so single one and co-product is analogous to the decay of particle to two.  $\Delta$  allows to construct higher generators of the algebra.

Lie-algebra can mean here ordinary finite-dimensional simple Lie algebra, Kac-Moody algebra or Virasoro algebra. In the case of SUSY it means conformal algebra of  $M^4$ - or rather its super counterpart. Witten, Nappi and Dolan have described the notion of Yangian for super-conformal algebra in very elegant and concrete manner in the article *Yangian Symmetry in  $D=4$  superconformal Yang-Mills theory* [?]. Also Yangians for gauge groups are discussed.

In the general case Yangian resembles Kac-Moody algebra with discrete index  $n$  replaced with a continuous one. Discrete index poses conditions on the Lie group and its representation (adjoint representation in the case of  $\mathcal{N} = 4$  SUSY). One of the conditions conditions is that the tensor product  $R \otimes R^*$  for representations involved contains adjoint representation only once. This condition is non-trivial. For  $SU(n)$  these conditions are satisfied for any representation. In the case of  $SU(2)$

the basic branching rule for the tensor product of representations implies that the condition is satisfied for the product of any representations.

Yangian algebra with a discrete basis is in many respects analogous to Kac-Moody algebra. Now however the generators are labelled by non-negative integers labeling the light-like incoming and outgoing momenta of scattering amplitude whereas in the case of Kac-Moody algebra also negative values are allowed. Note that only the generators with non-negative conformal weight appear in the construction of states of Kac-Moody and Virasoro representations so that the extension to Yangian makes sense.

The generating elements are labelled by the generators of ordinary conformal transformations acting in  $M^4$  and their duals acting in momentum space. These two sets of elements can be labelled by conformal weights  $n = 0$  and  $n = 1$  and their mutual commutation relations are same as for Kac-Moody algebra. The commutators of  $n = 1$  generators with themselves are however something different for a non-vanishing deformation parameter  $h$ . Serre's relations characterize the difference and involve the deformation parameter  $h$ . Under repeated commutations the generating elements generate infinite-dimensional symmetric algebra, the Yangian. For  $h = 0$  one obtains just one half of the Virasoro algebra or Kac-Moody algebra. The generators with  $n > 0$  are  $n + 1$ -local in the sense that they involve  $n + 1$ -forms of local generators assignable to the ordered set of incoming particles of the scattering amplitude. This non-locality generalizes the notion of local symmetry and is claimed to be powerful enough to fix the scattering amplitudes completely.

### 5.6.2 How to generalize Yangian symmetry in TGD framework?

As far as concrete calculations are considered, it is not much to say. It is however possible to keep discussion at general level and still say something interesting (as I hope!). The key question is whether it could be possible to generalize the proposed Yangian symmetry and geometric picture behind it to TGD framework.

1. The first thing to notice is that the Yangian symmetry of  $\mathcal{N} = 4$  SUSY in question is quite too limited since it allows only single representation of the gauge group and requires massless particles. One must allow all representations and massive particles so that the representation of symmetry algebra must involve states with different masses, in principle arbitrary spin and arbitrary internal quantum numbers. The candidates are obvious: Kac-Moody algebras [A2] and Virasoro algebras [A3] and their super counterparts. Yangians indeed exist for arbitrary super Lie algebras. In TGD framework conformal algebra of Minkowski space reduces to Poincare algebra and its extension to Kac-Moody allows to have also massive states.
2. The formal generalization looks surprisingly straightforward at the formal level. In zero energy ontology one replaces point like particles with partonic two-surfaces appearing at the ends of light-like orbits of wormhole throats located to the future and past light-like boundaries of causal diamond ( $CD \times CP_2$  or briefly CD). Here CD is defined as the intersection of future and past directed light-cones. The polygon with light-like momenta is naturally replaced with a polygon with more general momenta in zero energy ontology and having

partonic surfaces as its vertices. Non-point-likeness forces to replace the finite-dimensional super Lie-algebra with infinite-dimensional Kac-Moody algebras and corresponding super-Virasoro algebras assignable to partonic 2-surfaces.

3. This description replaces disjoint holomorphic surfaces in twistor space with partonic 2-surfaces at the boundaries of  $CD \times CP_2$  so that there seems to be a close analogy with Cachazo-Svrcek-Witten picture. These surfaces are connected by either light-like orbits of partonic 2-surface or space-like 3-surfaces at the ends of CD so that one indeed obtains the analog of polygon.

What does this then mean concretely (if this word can be used in this kind of context)?

1. At least it means that ordinary Super Kac-Moody and Super Virasoro algebras associated with isometries of  $M^4 \times CP_2$  annihilating the scattering amplitudes must be extended to a co-algebras with a non-trivial deformation parameter. Kac-Moody group is thus the product of Poincare and color groups. This algebra acts as deformations of the light-like 3-surfaces representing the light-like orbits of particles which are extremals of Chern-Simon action with the constraint that weak form of electric-magnetic duality holds true. I know so little about the mathematical side that I cannot tell whether the condition that the product of the representations of Super-Kac-Moody and Super-Virasoro algebras contains adjoint representation only once, holds true in this case. In any case, it would allow all representations of finite-dimensional Lie group in vertices whereas  $\mathcal{N} = 4$  SUSY would allow only the adjoint.
2. Besides this ordinary kind of Kac-Moody algebra there is the analog of Super-Kac-Moody algebra associated with the light-cone boundary which is metrically 3-dimensional. The finite-dimensional Lie group is in this case replaced with infinite-dimensional group of symplectomorphisms of  $\delta M^4_{+/-}$  made local with respect to the internal coordinates of the partonic 2-surface. This picture also justifies p-adic thermodynamics applied to either symplectic or isometry Super-Virasoro and giving thermal contribution to the vacuum conformal and thus to mass squared.
3. The construction of TGD leads also to other super-conformal algebras and the natural guess is that the Yangians of all these algebras annihilate the scattering amplitudes.
4. Obviously, already the starting point symmetries look formidable but they still act on single partonic surface only. The discrete Yangian associated with this algebra associated with the closed polygon defined by the incoming momenta and the negatives of the outgoing momenta acts in multi-local manner on scattering amplitudes. It might make sense to speak about polygons defined also by other conserved quantum numbers so that one would have generalized light-like curves in the sense that state are massless in 8-D sense.

### 5.6.3 Could Yangian symmetry provide a new view about conserved quantum numbers?

The Yangian algebra has some properties which suggest a new kind of description for bound states. The Cartan algebra generators of  $n = 0$  and  $n = 1$  levels of Yangian algebra commute. Since the co-product  $\Delta$  maps  $n = 0$  generators to  $n = 1$  generators and these in turn to generators with high value of  $n$ , it seems that they commute also with  $n \geq 1$  generators. This applies to four-momentum, color isospin and color hyper charge, and also to the Virasoro generator  $L_0$  acting on Kac-Moody algebra of isometries and defining mass squared operator.

Could one identify total four momentum and Cartan algebra quantum numbers as sum of contributions from various levels? If so, the four momentum and mass squared would involve besides the local term assignable to wormhole throats also  $n$ -local contributions. The interpretation in terms of  $n$ -parton bound states would be extremely attractive.  $n$ -local contribution would involve interaction energy. For instance, string like object would correspond to  $n = 1$  level and give  $n = 2$ -local contribution to the momentum. For baryonic valence quarks one would have 3-local contribution corresponding to  $n = 2$  level. The Yangian view about quantum numbers could give a rigorous formulation for the idea that massive particles are bound states of massless particles.

## 6 Generalization Of Ads/CFT Duality To TGD Framework

AdS/CFT duality has provided a powerful approach in the attempts to understand the non-perturbative aspects of super-string theories. The duality states that conformal field theory in  $n$ -dimensional Minkowski space  $M^n$  identifiable as a boundary of  $n + 1$ -dimensional space  $AdS_{n+1}$  is dual to a string theory in  $AdS_{n+1} \times S^{9-n}$ .

As a mathematical discovery the duality is extremely interesting but it seems that it need not have much to do with physics. From TGD point of view the reason is obvious: the notion of conformal invariance is quite too limited. In TGD framework conformal invariance is extended to a super-symplectic symmetry in  $\delta M_{\pm}^4 \times CP_2$ , whose Lie-algebra has the structure of conformal algebra. Also ordinary super-conformal symmetries associated with string world sheets are present as well as generalization of 2-D conformal symmetries to their analogs at light-cone boundary and light-like orbits of partonic 2-surfaces. In this framework AdS/CFT duality is expected to be modified and this seems to be the case.

The matrix elements of Kähler metric of WCW can be expressed in two ways. As contractions of the derivatives  $\partial_K \partial_{\bar{L}} K$  of the Kähler function of WCW with isometry generators or as anticommutators of WCW gamma matrices identified as supersymplectic Noether super charges assignable to fermionic strings connecting partonic 2-surfaces. Kähler function is identified as Kähler action for the Euclidian space-time regions with 4-D  $CP_2$  projection. Kähler action defines the Kähler-Dirac gamma matrices appearing in K-D action as contractions of canonical momentum currents with embedding space gamma matrices. Bosonic and fermionic degrees of freedom are therefore dual in a well-defined sense.



This observation suggests various generalizations. There is super-symmetry between Kähler action and Kähler-Dirac action. The problem is that induced spinor fields are localized at 2-D string world sheets. Strong form of holography implying effective 2-dimensionality suggests the solution to the paradox. The paradox disappears if the Kähler action is expressible as string area for the effective metric defined by the anti-commutators of K-D gamma matrices at string world sheet. This expression allows to understand how strings connecting partonic 2-surfaces give rise to the formation of gravitationally bound states. Bound states of macroscopic size are however possible only if one allows hierarchy of Planck constants. This representation of Kähler action can be seen as one aspect of the analog of AdS/CFT duality in TGD framework.

One can imagine also other realizations. For instance, Dirac determinant for the spinors associated with string world sheets should reduce to the exponent of Kähler action.

## 6.1 Does The Exponent Of Chern-Simons Action Reduce To The Exponent Of The Area Of Minimal Surfaces?

As I scanned of hep-th I found an interesting article (see <http://tinyurl.com/yckprg4f>) by Giordano, Peschanski, and Seki [?] based on AdS/CFT correspondence. What is studied is the high energy behavior of the gluon-gluon and quark-quark scattering amplitudes of  $\mathcal{N} = 4$  SUSY.

1. The proposal made earlier by Aldaya and Maldacena (see <http://tinyurl.com/ybnk6kbs>) [?] is that gluon-gluon scattering amplitudes are proportional to the imaginary exponent of the area of a minimal surface in  $AdS_5$  whose boundary is identified as *momentum space*. The boundary of the minimal surface would be polygon with light-like edges: this polygon and its dual are familiar from twistor approach.
2. Giordano, Peschanski, and Seki claim that quark-quark scattering amplitude for heavy quarks corresponds to the exponent of the area for a minimal surface in the *Euclidian* version of  $AdS_5$  which is hyperbolic space (space with a constant negative curvature): it is interpreted as a counterpart of WCW rather than momentum space and amplitudes are obtained by analytic continuation. For instance, a universal Regge behavior is obtained. For general amplitudes the exponent of the area alone is not enough since it does not depend on gluon quantum numbers and vertex operators at the edges of the boundary polygon are needed.

In the following my intention is to consider the formulation of this conjecture in quantum TGD framework. I hasten to inform that I am not a specialist in AdS/CFT and can make only general comments inspired by analogies with TGD and the generalization of AdS/CFT duality to TGD framework based on the localization of induced spinors at string world sheets, super-symmetry between bosonic and fermionic degrees of freedom at the level of WCW , and the notion of effective metric at string world sheets.

## 6.2 Does Kähler Action Reduce To The Sum Of Areas Of Minimal Surfaces In Effective Metric?

Minimal surface conjectures are highly interesting from TGD point of view. The weak form of electric magnetic duality implies the reduction of Kähler action to 3-D Chern-Simons terms. Effective 2-dimensionality implied by the strong form of General Coordinate Invariance suggests a further reduction of Chern-Simons terms to 2-D terms and the areas of string world sheet and of partonic 2-surface are the only non-topological options that one can imagine. Skeptic could of course argue that the exponent of the minimal surface area results as a characterizer of the quantum state rather than vacuum functional. In the following I end up with the proposal that the Kähler action should reduce to the sum of string world sheet areas in the effective metric defines by the anticommutators of Kähler-Dirac gamma matrices at string world sheets.

Let us look this conjecture in more detail.

1. In zero energy ontology twistor approach is very natural since all physical states are bound states of massless particles. Also virtual particles are composites of massless states. The possibility to have both signs of energy makes possible space-like momenta for wormhole contacts. Mass shell conditions at internal lines imply extremely strong constraints on the virtual momenta and both UV and IR finiteness are expected to hold true.
2. The weak form of electric magnetic duality [K29] implies that the exponent of Kähler action reduces to the exponent of Chern-Simons term for 3-D space-like surfaces at the ends of space-time surface inside CD and for light-like 3-surfaces. The coefficient of this term is complex since the contribution of Minkowskian regions of the space-time surface is imaginary ( $\sqrt{g_4}$  is imaginary) and that of Euclidian regions (generalized Feynman diagrams) real. The Chern-Simons term from Minkowskian regions is like Morse function and that from Euclidian regions defines Kähler function and stationary phase approximation makes sense. The two contributions are different since the space-like 3-surfaces contributing to Kähler function and Morse function are different.
3. Electric magnetic duality [K29] leads also to the conclusion that wormhole throats carrying elementary particle quantum numbers are Kähler magnetic monopoles. This forces to identify elementary particles as string like objects with ends having opposite monopole charges. Also more complex configurations are possible.

It is not quite clear what the scale of the stringyness is. The natural first guess inspired by quantum classical correspondence is that it corresponds to the p-adic length scale of the particle characterizing its Compton length. Second possibility is that it corresponds to electroweak scale. For leptons stringyness in Compton length scale might not have any fatal implications since the second end of string contains only neutrinos neutralizing the weak isospin of the state. This kind of monopole pairs could appear even in condensed matter scales: in particular if the proposed hierarchy of Planck constants [K8] is realized.

4. Strong form of General Coordinate Invariance requires effective 2-dimensionality. In given UV and IR resolutions either partonic 2-surfaces or string world sheets form a finite hierarchy of CDs inside CDs with given CD characterized by a discrete scale coming as an integer multiple of a fundamental scale (essentially  $CP_2$  size). The string world sheets have boundaries consisting of either light-like curves in induced metric at light-like wormhole throats and space-like curves at the ends of CD whose  $M^4$  projections are light-like. These braids intersect partonic 2-surfaces at discrete points carrying fermionic quantum numbers.

This implies a rather concrete analogy with  $AdS_5 \times S_5$  duality, which describes gluons as open strings. In zero energy ontology (ZEO) string world sheets are indeed a fundamental notion and the natural conjecture is that these surfaces are minimal surfaces whose area by quantum classical correspondence depends on the quantum numbers of the external particles. String tension in turn should depend on gauge couplings -perhaps only Kähler coupling strength- and geometric parameters like the size scale of CD and the p-adic length scale of the particle.

5. One can of course ask whether the metric defining the string area is induced metric or possibly the metric defined by the anti-commutators of Kähler-Dirac gamma matrices. The recent view does not actually leave any other alternative. The analog of AdS/CFT duality together with supersymmetry demands that Kähler action is proportional to the sum of the areas of string world sheets in this effective metric. Whether the vanishing of induced  $W$  fields (and possibly also  $Z^0$ ) making possible well-defined em charge for the spinor nodes is realized by the condition that the string world sheet is a minimal surface in the effective metric remains an open question.

The assumption that ordinary minimal surfaces are in question is not consistent with the TGD view about the formation of gravitational bound states and if string tension is  $1/\hbar G$  as in string models, only bound states with size of order Planck length are possible. This strongly favors effective metric giving string tension proportional to  $1/\hbar_{eff}^2$ . How  $1/\hbar_{eff}^2$  proportionality might be understood is discussed in [?] in terms electric-magnetic duality.

6. One can of course still consider also the option that ordinary minimal surfaces are in question. Are the minimal surfaces in question minimal surfaces of the embedding space  $M^4 \times CP_2$  or of the space-time surface  $X^4$ ? All possible 2-surfaces at the boundary of CD must be allowed so that they cannot correspond to minimal surfaces in  $M^4 \times CP_2$  unless one assumes that they emerge in stationary phase approximation only. The boundary conditions at the ends of CD could however be such that *any* partonic 2-surface correspond to a minimal surfaces in  $X^4$ . Same applies to string world sheets. One might even hope that these conditions combined with the weak form of electric magnetic duality fixes completely the boundary conditions at wormhole throats and space-like ends of space-time surface.

The trace of the second fundamental form orthogonal to the string world sheet/partonic 2-surface as sub-manifold of space-time surface would vanish:

this is nothing but a generalization of the geodesic motion obtained by replacing word line with a 2-D surface. It does not imply the vanishing of the trace of the second fundamental form in  $M^4 \times CP_2$  having interpretation as a generalization of particle acceleration [K25]. Effective 2-dimensionality would be realized if Chern-Simons terms reduce to a sum of the areas of these minimal surfaces.

These arguments suggest that scattering amplitudes are proportional to the product of exponents of 2-dimensional actions which can be either imaginary or real. Imaginary exponent would be proportional to the total area of string world sheets and the imaginary unit would come naturally from  $\sqrt{g_2}$ , where  $g_2$  is effective metric most naturally. Teal exponent proportional to the total area of partonic 2-surfaces. The coefficient of these areas would not in general be same.

The reduction of the Kähler action from Minkowskian regions to Chern-Simons terms means that Chern-Simons terms reduce to actions assignable to string world sheets. The equality of the Minkowskian and Euclidian Chern-Simons terms is suggestive but not necessarily true since there could be also other Chern-Simons contributions than those assignable to wormhole throats and the ends of space-time. The equality would imply that the total area of string world sheets equals to the total area of partonic 2-surfaces suggesting strongly a duality meaning that either Euclidian or Minkowskian regions carry the needed information.

### **6.3 Surface Area As Geometric Representation Of Entanglement Entropy?**

I encountered a link to a talk by James Sully and having the title “Geometry of Compression” (see <http://tinyurl.com/ycuu8xcr>). I must admit that I understood very little about the talk. My not so educated guess is however that information is compressed: UV or IR cutoff eliminating entanglement in short length scales and describing its presence in terms of density matrix - that is thermodynamically - is another manner to say it. The TGD inspired proposal for the interpretation of the inclusions of hyper-finite factors of type  $II_1$  (HFFs) [K28] is in spirit with this.

The space-time counterpart for the compression would be in TGD framework discretization. Discretizations using rational points (or points in algebraic extensions of rationals) make sense also p-adically and thus satisfy number theoretic universality. Discretization would be defined in terms of intersection (rational or in algebraic extension of rationals) of real and p-adic surfaces. At the level of “world of classical worlds” the discretization would correspond to - say - surfaces defined in terms of polynomials, whose coefficients are rational or in some algebraic extension of rationals. Pinary UV and IR cutoffs are involved too. The notion of p-adic manifold allows to interpret the p-adic variants of space-time surfaces as cognitive representations of real space-time surfaces.

Finite measurement resolution does not allow state function reduction reducing entanglement totally. In TGD framework also negentropic entanglement stable under Negentropy Maximization Principle (NMP) is possible [K11]. For HFFs the projection into single ray of Hilbert space is indeed impossible: the reduction takes always to infinite-D sub-space.

The visit to the URL was however not in vain. There was a link to an article (see <http://tinyurl.com/y9h3qtr8>) [?] discussing the geometrization of entanglement entropy inspired by the AdS/CFT hypothesis.

Quantum classical correspondence is basic guiding principle of TGD and suggests that entanglement entropy should indeed have space-time correlate, which would be the analog of Hawking-Bekenstein entropy.

### 6.3.1 Generalization of AdS/CFT to TGD context

AdS/CFT generalizes to TGD context in non-trivial manner. There are two alternative interpretations, which both could make sense. These interpretations are not mutually exclusive. The first interpretation makes sense at the level of “world of classical worlds” ( WCW ) with symplectic algebra and extended conformal algebra associated with  $\delta M_{\pm}^4$  replacing ordinary conformal and Kac-Moody algebras. Second interpretation at the level of space-time surface with the extended conformal algebras of the light-likes orbits of partonic 2-surfaces replacing the conformal algebra of boundary of  $AdS^n$ .

#### 1. First interpretation

For the first interpretation 2-D conformal invariance is generalised to 4-D conformal invariance relying crucially on the 4-dimensionality of space-time surfaces and Minkowski space.

1. One has an extension of the conformal invariance provided by the symplectic transformations of  $\delta CD \times CP_2$  for which Lie algebra has the structure of conformal algebra with radial light-like coordinate of  $\delta M_{\pm}^4$  replacing complex coordinate  $z$ .
2. One could see the counterpart of  $AdS_n$  as embedding space  $H = M^4 \times CP_2$  completely unique by twistorial considerations and from the condition that standard model symmetries are obtained and its causal diamonds defined as sub-sets  $CD \times CP_2$ , where CD is an intersection of future and past directed light-cones. I will use the shorthand CD for  $CD \times CP_2$ . Strings in  $AdS_5 \times S^5$  are replaced with space-time surfaces inside 8-D CD.
3. For this interpretation 8-D CD replaces the 10-D space-time  $AdS_5 \times S^5$ . 7-D light-like boundaries of CD correspond to the boundary of say  $AdS_5$ , which is 4-D Minkowski space so that zero energy ontology (ZEO) allows rather natural formulation of the generalization of AdS/CFT correspondence since the positive and negative energy parts of zero energy states are localized at the boundaries of CD.

### 6.3.2 Second interpretation

For the second interpretation relies on the observation that string world sheets as carriers of induced spinor fields emerge in TGD framework from the condition that electromagnetic charge is well-defined for the modes of induced spinor field.

1. One could see the 4-D space-time surfaces  $X^4$  as counterparts of  $AdS_4$ . The boundary of  $AdS_4$  is replaced in this picture with 3-surfaces at the ends of

space-time surface at opposite boundaries of CD and by strong form of holography the union of partonic 2-surfaces defining the intersections of the 3-D boundaries between Euclidian and Minkowskian regions of space-time surface with the boundaries of CD. Strong form of holography in TGD is very much like ordinary holography.

2. Note that one has a dimensional hierarchy: the ends of the boundaries of string world sheets at boundaries of CD as point-like partices, boundaries as fermion number carrying lines, string world sheets, light-like orbits of partonic 2-surfaces, 4-surfaces, embedding space  $M^4 \times CP_2$ . Clearly the situation is more complex than for AdS/CFT correspondence.
3. One can restrict the consideration to 3-D sub-manifolds  $X^3$  at either boundary of causal diamond (CD): the ends of space-time surface. In fact, the position of the other boundary is not well-defined since one has superposition of CDs with only one boundary fixed to be piece of light-cone boundary. The delocalization of the other boundary is essential for the understanding of the arrow of time. The state function reductions at fixed boundary leave positive energy part (say) of the zero energy state at that boundary invariant (in positive energy ontology entire state would remain unchanged) but affect the states associated with opposite boundaries forming a superposition which also changes between reduction: this is analog for unitary time evolution. The average for the distance between tips of CDs in the superposition increases and gives rise to the flow of time.
4. One wants an expression for the entanglement entropy between  $X^3$  and its partner. Bekenstein area law allows to guess the general expression for the entanglement entropy: for the proposal discussed in the article the entropy would be the area of the boundary of  $X^3$  divided by gravitational constant:  $S = A/4G$ . In TGD framework gravitational constant might be replaced by the square of  $CP_2$  radius apart from numerical constant. How gravitational constant emerges in TGD framework is not completely understood although one can deduce for it an estimate using dimensional analyses. In any case, gravitational constant is a parameter which characterizes GRT limit of TGD in which many-sheeted space-time is in long scales replaced with a piece of Minkowski space such that the classical gravitational fields and gauge potentials for sheets are summed. The physics behind this relies on the generalization of linear superposition of fields: the effects of different space-time sheets particle touching them sum up rather than fields.
5. The counterpart for the boundary of  $X^3$  appearing in the proposal for the geometrization of the entanglement entropy naturally corresponds to partonic 2-surface or a collection of them if strong form of holography holds true.

There is however also another candidate to be considered! Partonic 2-surfaces are basic objects, and one expects that the entanglement between fundamental fermions associated with distinct partonic 2-surfaces has string world sheets as space-time correlates. Could the area of the string world sheet in the effective metric defined by the anti-commutators of K-D gamma matrices at string

world sheet provide a measure for entanglement entropy? If this conjecture is correct: the entanglement entropy would be proportional to Kähler action. Also negative values are possible for Kähler action in Minkowskian regions but in TGD framework number theoretic entanglement entropy having also negative values emerges naturally.

Which of these guesses is correct, if any? Or are they equivalent?

### 6.3.3 With what kind of systems 3-surfaces can entangle?

With what system  $X^3$  is entangled/can entangle? There are several options to consider and they could correspond to the two TGD variants for the AdS/CFT correspondence.

1.  $X^3$  could correspond to a wormhole contact with Euclidian signature of induced metric. The entanglement would be between it and the exterior region with Minkowskian signature of the induced metric.
2.  $X^3$  could correspond to single sheet of space-time surface connected by wormhole contacts to a larger space-time sheet defining its environment. More precisely,  $X^3$  and its complement would be obtained by throwing away the wormhole contacts with Euclidian signature of induce metric. Entanglement would be between these regions. In the generalization of the formula

$$S = \frac{A}{4\hbar G}$$

area  $A$  would be replaced by the total area of partonic 2-surfaces and  $G$  perhaps with  $CP_2$  length scale squared.

3. In ZEO the entanglement could also correspond to time-like entanglement between the 3-D ends of the space-time surface at opposite light-like boundaries of CD. M-matrix, which can be seen as the analog of thermal S-matrix, decomposes to a product of hermitian square root of density matrix and unitary S-matrix and this hermitian matrix could also define p-adic thermodynamics. Note that in ZEO quantum theory can be regarded as square root of thermodynamics.

### 6.3.4 Minimal surface property is not favored in TGD framework

Minimal surface property for the 3-surfaces  $X^3$  at the ends of space-time surface looks at first glance strange but a proper generalization of this condition makes sense if one assumes strong form of holography. Strong form of holography realizes General Coordinate Invariance (GCI) in strong sense meaning that light-like parton orbits and space-like 3-surfaces at the ends of space-time surfaces are equivalent physically. As a consequence, partonic 2-surfaces and their 4-D tangent space data must code for the quantum dynamics.

The mathematical realization is in terms of conformal symmetries accompanying the symplectic symmetries of  $\delta M_{\pm}^4 \times CP_2$  and conformal transformations of the light-like partonic orbit [K29]. The generalizations of ordinary conformal algebras

correspond to conformal algebra, Kac-Moody algebra at the light-like parton orbits and to symplectic transformations  $\delta M^4 \times CP_2$  acting as isometries of WCW and having conformal structure with respect to the light-like radial coordinate plus conformal transformations of  $\delta M_+^4$ , which is metrically 2-dimensional and allows extended conformal symmetries.

1. If the conformal realization of the strong form of holography works, conformal transformations act at quantum level as gauge symmetries in the sense that generators with non-vanishing conformal weight are zero or generate zero norm states. Conformal degeneracy can be eliminated by fixing the gauge somehow. Classical conformal gauge conditions analogous to Virasoro and Kac-Moody conditions satisfied by the 3-surfaces at the ends of CD are natural in this respect. Similar conditions would hold true for the light-like partonic orbits at which the signature of the induced metric changes.
2. What is also completely new is the hierarchy of conformal symmetry breakings associated with the hierarchy of Planck constants  $h_{eff}/h = n$  [K8]. The deformations of the 3-surfaces which correspond to non-vanishing conformal weight in algebra or any sub-algebra with conformal weights vanishing modulo  $n$  give rise to vanishing classical charges and thus do not affect the value of the Kähler action [K29].

The inclusion hierarchies of conformal sub-algebras are assumed to correspond to those for hyper-finite factors. There is obviously a precise analogy with quantal conformal invariance conditions for Virasoro algebra and Kac-Moody algebra. There is also hierarchy of inclusions which corresponds to hierarchy of measurement resolutions. An attractive interpretation is that singular conformal transformations relate to each other the states for broken conformal symmetry. Infinitesimal transformations for symmetry broken phase would carry fractional conformal weights coming as multiples of  $1/n$ .

3. Conformal gauge conditions need not reduce to minimal surface conditions holding true for all variations.
4. Note that Kähler action reduces to Chern-Simons term at the ends of CD if weak form of electric magnetic duality holds true. The conformal charges at the ends of CD cannot however reduce to Chern-Simons charges by this condition since only the charges associated with  $CP_2$  degrees of freedom would be non-trivial.

The way out of the problem is provided by the generalization of AdS/CFT conjecture. String area is estimated in the effective metric provided by the anti-commutator of K-D gamma matrices at string world sheet.

## 6.4 Related Ideas

p-Adic mass calculations led to the introduction of the p-adic variant of Bekenstein-Hawkin law in which Planck length is replaced by p-adic length scale. This generalization is in spirit with the idea that string world sheet area is estimated in effective rather than induced metric.



### 6.4.1 p-Adic variant of Bekenstein-Hawking law

When the 3-surface corresponds to elementary particle, a direct connection with p-adic thermodynamics suggests itself and allows to answer the questions above. p-Adic thermodynamics could be interpreted as a description of the entanglement with environment. In ZEO the entanglement could also correspond to time-like entanglement between the 3-D ends of the space-time surface at opposite light-like boundaries of CD. M-matrix, which can be seen as the analog of thermal S-matrix, decomposes to a product of hermitian square root of density matrix and unitary S-matrix and this hermitian matrix could also define p-adic thermodynamics.

1. p-Adic thermodynamics [K14] would not be for energy but for mass squared (or scaling generator  $L_0$ ) would describe the entanglement of the particle with environment defined by the larger space-time sheet. Conformal weights would come as positive powers of integers ( $p_0^L$  would replace  $\exp(-H/T)$  to guarantee the number theoretical existence and convergence of the Boltzmann weight: note that conformal invariance that is integer spectrum of  $L_0$  is also essential).
2. The interactions with environment would excite very massive  $CP_2$  mass scale excitations (mass scale is about  $10^{-4}$  times Planck mass) of the particle and give it thermal mass squared identifiable as the observed mass squared. The Boltzmann weights would be extremely small having p-adic norm about  $1/p^n$ ,  $p$  the p-adic prime:  $M_{127} = 2^{127} - 1$  for electron.
3. One of the first ideas inspired by p-adic vision was that p-adic entropy could be seen as a p-adic counterpart of Bekenstein-Hawking entropy [K15].  $S = (R^2/\hbar^2) \times M^2$  holds true identically apart from numerical constant. Note that one could interpret  $R^2 M/\hbar$  as the counterpart of Schwarzschild radius. Note that this radius is proportional to  $1/\sqrt{p}$  so that the area  $A$  would correspond to the area defined by Compton length. This is in accordance with the third option.

### 6.4.2 What is the space-time correlate for negentropic entanglement?

The new element brought in by TGD framework is that number theoretic entanglement entropy is negative for negentropic entanglement assignable to unitary entanglement (in the sense that entanglement matrix is proportional to a unitary matrix) and NMP states that this negentropy increases [K11]. Since entropy is essentially number of energy degenerate states, a good guess is that the number  $n = h_{eff}/h$  of space-time sheets associated with  $h_{eff}$  defines the negentropy. An attractive space-time correlate for the negentropic entanglement is braiding. Braiding defines unitary S-matrix between the states at the ends of braid and this entanglement is negentropic. This entanglement gives also rise to topological quantum computation.

## 6.5 The Importance Of Being Light-Like

The singular geometric objects associated with the space-time surface have become increasingly important in TGD framework. In particular, the recent progress has made clear that these objects might be crucial for the understanding of quantum

TGD. The singular objects are associated not only with the induced metric but also with the effective metric defined by the anti-commutators of the Kähler-Dirac gamma matrices appearing in the Kähler-Dirac equation and determined by the Kähler action.

### 6.5.1 The singular objects associated with the induced metric

Consider first the singular objects associated with the induced metric.

1. At light-like 3-surfaces defined by wormhole throats the signature of the induced metric changes from Euclidian to Minkowskian so that 4-metric is degenerate. These surfaces are carriers of elementary particle quantum numbers and the 4-D induced metric degenerates locally to 3-D one at these surfaces.
2. Braid strands at partonic orbits - fermion lines identified as boundaries of string world sheets in the more recent terminology - are most naturally light-like curves: this correspond to the boundary condition for open strings. One can assign fermion number to the braid strands. Braid strands allow an identification as curves along which the Euclidian signature of the string world sheet in Euclidian region transforms to Minkowskian one. Number theoretic interpretation would be as a transformation of complex regions to hyper-complex regions meaning that imaginary unit  $i$  satisfying  $i^2 = -1$  becomes hyper-complex unit  $e$  satisfying  $e^2 = 1$ . The complex coordinates  $(z, \bar{z})$  become hyper-complex coordinates  $(u = t + ex, v = t - ex)$  giving the standard light-like coordinates when one puts  $e = 1$ .

### 6.5.2 The singular objects associated with the effective metric

There are also singular objects assignable to the effective metric. According to the simple arguments already developed, string world sheets and possibly also partonic 2-surfaces are singular objects with respect to the effective metric defined by the anti-commutators of the Kähler-Dirac gamma matrices rather than induced gamma matrices. Therefore the effective metric might be more than a mere formal structure. The following is of course mere speculation and should be taken as such.

1. For instance, quaternionicity of the space-time surface *might* allow an elegant formulation in terms of the effective metric avoiding the problems due to the Minkowski signature. This is achieved if the effective metric has Euclidian signature  $\epsilon \times (1, 1, 1, 1)$ ,  $\epsilon = \pm 1$  or a complex counterpart of the Minkowskian signature  $\epsilon(1, 1, -1, -1)$ .
2. String world sheets and perhaps also partonic 2-surfaces might be understood as singularities of the effective metric. What happens that the effective metric with Euclidian signature  $\epsilon \times (1, 1, 1, 1)$  transforms to the signature  $\epsilon(1, 1, -1, -1)$  (say) at string world sheet so that one would have the degenerate signature  $\epsilon \times (1, 1, 0, 0)$  at the string world sheet.

What is amazing is that this works also number theoretically. It came as a total surprise to me that the notion of hyper-quaternions as a closed algebraic structure indeed exists. The hyper-quaternionic units would be given

by  $(1, I, iJ, iK)$ , where  $i$  is a commuting imaginary unit satisfying  $i^2 = -1$ . Hyper-quaternionic numbers defined as combinations of these units with real coefficients do form a closed algebraic structure which however fails to be a number field just like hyper-complex numbers do. Note that the hyper-quaternions obtained with real coefficients from the basis  $(1, iI, iJ, iK)$  fail to form an algebra since the product is not hyper-quaternion in this sense but belongs to the algebra of complexified quaternions. The same problem is encountered in the case of hyper-octonions defined in this manner. This has been a stone in my shoe since I feel strong disrelish towards Wick rotation as a trick for moving between different signatures.

3. Could also partonic 2-surfaces correspond to this kind of singular 2-surfaces? In principle, 2-D surfaces of 4-D space intersect at discrete points just as string world sheets and partonic 2-surfaces do so that this might make sense. By complex structure the situation is algebraically equivalent to the analog of plane with non-flat metric allowing all possible signatures  $(\epsilon_1, \epsilon_2)$  in various regions. At light-like curve either  $\epsilon_1$  or  $\epsilon_2$  changes sign and light-like curves for these two kinds of changes can intersect as one can easily verify by drawing what happens. At the intersection point the metric is completely degenerate and simply vanishes.
4. Replacing real 2-dimensionality with complex 2-dimensionality, one obtains by the universality of algebraic dimension the same result for partonic 2-surfaces and string world sheets. The braid ends at partonic 2-surfaces representing the intersection points of 2-surfaces of this kind would have completely degenerate effective metric so that the Kähler-Dirac gamma matrices would vanish implying that energy momentum tensor vanishes as does also the induced Kähler field.
5. The effective metric suffers a local conformal scaling in the critical deformations identified in the proposed manner. Since ordinary conformal group acts on Minkowski space and leaves the boundary of light-cone invariant, one has two conformal groups. It is not however clear whether the  $M^4$  conformal transformations can act as symmetries in TGD, where the presence of the induced metric in Kähler action breaks  $M^4$  conformal symmetry. As found, also in TGD framework the Kac-Moody currents assigned to the braid strands generate Yangian: this is expected to be true also for the Kac-Moody counterparts of the conformal algebra associated with quantum criticality. On the other hand, in twistor program one encounters also two conformal groups and the space in which the second conformal group acts remains somewhat mysterious object. The Lie algebras for the two conformal groups generate the conformal Yangian and the integrands of the scattering amplitudes are Yangian invariants. Twistor approach should apply in TGD if zero energy ontology is right. Does this mean a deep connection?

What is also intriguing that twistor approach in principle works in strict mathematical sense only at signatures  $\epsilon \times (1, 1, -1-1)$  and the scattering amplitudes in Minkowski signature are obtained by analytic continuation. Could the effec-

tive metric give rise to the desired signature? Note that the notion of massless particle does not make sense in the signature  $\epsilon \times (1, 1, 1, 1)$ .

These arguments provide genuine a support for the notion of quaternionicity and suggest a connection with the twistor approach.

## 7 Could One Define Dynamical Homotopy Groups In WCW?

Agostino Prastaro - working as professor at the University of Rome - has done highly interesting work with partial differential equations, also those assignable to geometric variational principles such as Kähler action in TGD [A6, A7]. I do not understand the mathematical details but the key idea is a simple and elegant generalization of Thom's cobordism theory, and it is difficult to avoid the idea that the application of Prastaro's idea might provide insights about the preferred extremals, whose identification is now on rather firm basis [K27].

One could also consider a definition of what one might call dynamical homotopy groups as a genuine characteristics of WCW topology. The first prediction is that the values of conserved classical Noether charges correspond to disjoint components of WCW. Could this mean that the natural topology in the parameter space of Noether charges zero modes of WCW metric) is p-adic? An analogous conjecture was made on basis of spin glass analogy long time ago. Second surprise is that the only the six lowest dynamical homotopy groups of WCW would be non-trivial. The finite number of these groups dictate by the dimension of embedding space suggests also an interpretation as analogs of homology groups.

In the following the notion of cobordism is briefly discussed and the idea of Prastaro about assigning cobordism with partial differential equations is discussed.

### 7.1 About Cobordism As A Concept

To get some background consider first the notion of cobordism (<http://tinyurl.com/y7wdhtmv>).

1. Thom's cobordism theory [A12] is inspired by the question "When an  $n$ -manifold can be represented as a boundary of  $n + 1$ -manifold". One can also pose additional conditions such as continuity, smoothness, orientability, one can add bundles structures and require that they are induced to  $n$ -manifold from that of  $n + 1$ -manifold. One can also consider sub-manifolds of some higher-dimensional manifold.

One can also fix  $n$ -manifold  $M$  and ask "What is the set of  $n$ -manifolds  $N$  with the property that there exists  $n + 1$ -manifold  $W$  having union of  $M \cup N$  as its boundary". One can also allow  $M$  to have boundary and pose the same question by allowing also the boundary of connecting  $n+1$ -manifold  $W$  contain also the orbits of boundaries of  $M$  and  $N$ .

The cobordism class of  $M$  can be defined as the set of manifolds  $N$  cobordant with  $M$  - that is connectable in this manner. They have same cobordism class

since cobordism is equivalence relation. The classes form also a group with respect to disjoint union. Cobordism is much rougher equivalence relation than diffeomorphy or homeomorphy since topology changes are possible. For instance, every 3-D closed un-oriented manifold is a boundary of a 4-manifold! Same is true for orientable cobordisms. Cobordism defines a category: objects are (say closed) manifolds and morphisms are cobordisms.

2. The basic result of Morse, Thom, and Milnor is that cobordism as topology changes can be engineered from elementary cobordisms. One take manifold  $M \times I$  and imbeds to its other n-dimensional end the manifold  $S^p \times D^q$ ,  $n = p + q$ , removes its interior and glues back  $D^{p+1} \times S^{q-1}$  along its boundary to the boundary of the resulting hole. This gives n-manifold with different topology, call it  $N$ . The outcome is a cobordism connecting  $M$  and  $N$  unless there are some obstructions.

There is a connection with Morse theory (<http://tinyurl.com/yeh4chg9>) in which cobordism can be seen as a mapping of  $W$  to a unit interval such that the inverse images define a slicing of  $W$  and the inverse images at ends correspond to  $M$  and  $N$ .

3. One can generalize the abstract cobordism to that for  $n$ -sub-manifolds of a given embedding space. This generalization is natural in TGD framework. This might give less trivial results since not all connecting manifolds are imbeddable into a given embedding space. If connecting 4-manifolds connecting 3-manifolds with Euclidian signature (of induced metric) are assumed to have a Minkowskian signature, one obtains additional conditions, which might be too strong (the classical result of Geroch [A13] implies that non-trivial cobordism implies closed time loops - impossible in TGD).

From TGD point of view this is too strong a condition and in TGD framework space-time surfaces with both Euclidian and Minkowskian signature of the induced metric are allowed. Also cobordisms singular as 4-surfaces are analogous to 3-vertices of Feynman diagrams are allowed.

## 7.2 Prastaro's Generalization Of Cobordism Concept To The Level Of Partial Differential Equations

I am not enough mathematician in technical sense of the word to develop overall view about what Prastaro has done and I have caught only the basic idea. I have tried to understand the articles [A6, A7] with title "Geometry of PDE's. I/II: Variational PDE's and integral bordism groups" (<http://tinyurl.com/yb9wey8c> and <http://tinyurl.com/y9x55qmk>), which seem to correspond to my needs. The key idea is to generalize the cobordism concept also to partial differential equations with cobordism replaced with the time evolution defined by partial differential equation. In particular, to geometric variational principles defining as their extremals the counterparts of cobordisms.

Quite generally, and especially so in the case of the conservation of Noether charges give rise to strong selection rules since two  $n$ -surfaces with different classical charges cannot be connected by extremals of the variational principle. Note however

that the values of the conserved charges depend on the normal derivatives of the embedding space coordinates at the  $n$ -dimensional ends of cobordism. If one poses additional conditions fixing these normal derivatives, the selection rules become even stronger. In TGD framework Bohr orbit property central for the notion of WCW geometry and holography allows to hope that conserved charges depend on 3-surfaces only.

What is so beautiful in this approach that it promises to generalize the notion of cobordism and perhaps also the notions of homotopy/homology groups so that they would apply to partial differential equations quite generally, and especially so in the case of geometric variational principles giving rise to  $n + 1$ -surfaces connecting  $n$ -surfaces characterizing the initial and final states classically. TGD with  $n = 3$  seems to be an ideal applications for these ideas.

Prastaro also proposes a generalization of cobordism theory to super-manifolds and quantum super-manifolds. The generalization in the case of quantum theory utilizing path integral does not pose conditions on classical connecting field configurations. In TGD framework these generalizations are not needed since fermion number is geometrized in terms of embedding space gamma matrices and super(-symplectic) symmetry is realized differently.

### 7.3 Why Prastaro's Idea Resonates So Strongly With TGD

Before continuing I want to make clear why Prastaro's idea resonates so strongly with TGD.

1. One of the first ideas as I started to develop TGD was that there might be selection rules analogous to those of quantum theory telling which 3-surfaces can be connected by a space-time surface. At that time I still believed in path integral formalism assuming that two 3-surfaces at different time slices with different values of Minkowski time can be connected by any space-time surface for which embedding space coordinates have first derivatives.

I soon learned about Thom's theory but was greatly disappointed since no selection rules were involved in the category of abstract 3-manifolds. I thought that possible selection rules should result from the imbeddability of the connecting four-manifold to  $H = M^4 \times CP_2$  but my gut feeling was that these rules are more or less trivial since so many connecting 4-manifolds exist and some of them are very probably imbeddable.

One possible source of selection rules could have been the condition that the induced metric has Minkowskian signature - one could justify it in terms of classical causality. This restricts strongly topology change in general relativity (<http://tinyurl.com/y6vuopgj>). Geroch's classical result [A13] states that non-trivial smooth Lorentz cobordism between compact 3-surfaces implies the existence of closed time loop - not possible in TGD framework. Second non-encouraging result is that scalar field propagating in trouser topology leads to an occurrence of infinite energy burst (<http://tinyurl.com/ybbuwfj>).

In the recent formulation of TGD however also Euclidian signature of the induced metric is allowed. For space-time counterparts of 3-particle vertices

three space-time surfaces are glued along their smooth 3-D ends whereas space-time surface fails to be everywhere smooth manifold. This picture fits nicely with the idea that one can engineer space-time surfaces by gluing them together along their ends.

2. At that time (before 1980) the discovery of the geometry of the “World of Classical Worlds” (WCW) as a possible solution to the failures of canonical quantization and path integral formalism was still at distance of ten years in future. Around 1985 I discovered the notion of WCW. I made some unsuccessful trials to construct its geometry, and around 1990 finally realized that 4-D general coordinate invariance is needed although basic objects are 3-D surfaces.

This is realized if classical physics is an exact part of quantum theory - not only something resulting in a stationary phase approximation. Classical variational principle should assign to a 3-surface a physically unique space-time surface - the analog of Bohr orbit - and the action for this surface would define Kähler function defining the Kähler geometry of WCW using standard formula.

This led to a notion of preferred extremal: absolute minimum of Kähler action was the first guess and might indeed make sense in the space-time regions with Euclidian signature of induced metric but not in Minkowskian regions, which give to the vacuum functional and exponential of Minkowskian Kähler action multiplied by imaginary unit coming from  $\sqrt{g}$  - just as in quantum field theories. Euclidian regions give the analog of the free energy exponential of thermodynamics and transform path integral to mathematically well-defined functional integral.

3. After having discovered the notion of preferred extremal, I should have also realized that an interesting generalization of cobordism theory might make sense after all, and could even give rise to the classical counterparts of the selection rules! For instance, conservation of isometry charges defines equivalence classes of 3-surfaces endowed with tangent space data. Bohr orbit property could fix the tangent space data (normal derivatives of embedding space coordinates) so that conserved classical charges would characterize 3-surfaces alone and thus cobordism equivalence classes and become analogous to topological invariants. This would be in spirit with the attribute “Topological” in TGD!

## 7.4 What Preferred Extremals Are?

The topology of WCW has remained mystery hitherto - partly due to my very limited technical skills and partly by the lack of any real physical idea. The fact, that p-adic topology seems to be natural at least as an effective topology for the maxima of Kähler function of WCW gave a hint but this was not enough.

I hope that the above summary has made clear why the idea about dynamical cobordism and even dynamical homotopy theory is so attractive in TGD framework. One could even hope that dynamics determines not only Kähler geometry but also the topology of WCW to some extent at least! To get some idea what might be

involved one must however first tell about the recent situation concerning the notion of preferred extremal.

1. The recent formulation for the notion of preferred extremal relies on strong form of General Coordinate Invariance (SGCI). SGCI states that two kinds of 3-surfaces can be identified as fundamental objects. Either the light-like 3-D orbits of partonic 2-surfaces defining boundaries between Minkowskian and Euclidian space-time regions or the space-like 3-D ends of space-time surfaces at boundaries of CD. Since both choices are equally good, partonic 2-surfaces and their tangent space-data at the ends of space-time should be the most economic choice.

This eventually led to the realization that partonic 2-surfaces and string world sheets should be enough for the formulation of quantum TGD. Classical fields in the interior of space-time surface would be needed only in quantum measurement theory, which demands classical physics in order to interpret the experiments.

2. The outcome is strong form of holography (SH) stating that quantum physics should be coded by string world sheets and partonic 2-surfaces inside given causal diamond (CD). SH is very much analogous to the AdS/CFT correspondence but is much simpler: the simplicity is made possible by much larger group of conformal symmetries.

If these 2-surfaces satisfy some consistency conditions one can continue them to 4-D space-time surface inside CD such that string world sheets are surfaces inside them satisfying the condition that charged (possibly all) weak gauge potentials identified as components of the induced spinor connection vanish at the string world sheets and also that energy momentum currents flow along these surfaces. String world sheets carry second quantized free induced spinor fields and fermionic oscillator operator basis is used to construct WCW gamma matrices.

3. The 3-surfaces at the ends of WCW must satisfy strong conditions to guarantee effective 2-dimensionality. Quantum criticality suggests the identification of these conditions. All Noether charges assignable to a sub-algebra of super-symplectic algebra isomorphic to it and having conformal weights which are  $n$ -multiples of those of entire algebra vanish/annihilate quantum states. One has infinite fractal hierarchy of broken super-conformal symmetries with the property that the sub-algebra is isomorphic with the entire algebra. This like a ball at the top of ball at the top of ....

The speculative vision is that super-symplectic subalgebra with weights coming as  $n$ -ples of those for the entire algebra acts as an analog of conformal gauge symmetries on light-like orbits of partonic 2-surfaces, and gives rise to a pure gauge degeneracy whereas other elements of super-symplectic algebra act as dynamical symmetries. The hierarchy of quantum criticalities defines hierarchies of symmetry breakings characterized by hierarchies of sub-algebras for which one  $n_{i+1}$  is divisible by  $n_i$ . The proposal is that conformal gauge invariance means that the analogs of Bohr orbits are determined only apart



from conformal gauge transformations forming to  $n_i$  conformal equivalence classes so that effectively one has  $n_i$  discrete degrees of freedom assignable to light-like partonic orbits.

4. In this framework manifolds  $M$  and  $N$  would correspond the 3-surfaces at the boundaries of CD and containing a collection strings carrying induced spinor fields. The connecting 4-surface  $W$  would contain string world sheets and the light-like orbits of partonic 2-surfaces as simultaneous boundaries for Minkowskian and Euclidian regions.

Propagator line has several meanings depending on whether one considers particles as strings, as single fermion states localizable at the ends of strings, or as Euclidian space-time regions or their light-like boundaries with singular induced metric having vanishing determinant. Vertices appear as generalizations of the stringy vertices and as generalization of the vertices of Feynman diagrams in which the incoming 4-surfaces meet along their ends.

1. Propagator line has several meanings depending on whether one considers particles as strings, as single fermion states localizable at the ends of strings, or as Euclidian space-time regions or their light-like boundaries with degenerate induced metric with vanishing determinant. Vertices appear as generalizations of the stringy vertices and as generalization of the vertices of Feynman diagrams in which the incoming 4-surfaces meet along their ends.

- (a) The lines of generalized Feynman graphs defined in topological sense are identified as slightly deformed pieces of  $CP_2$  defining wormhole contacts connecting two Minkowskian regions and having wormhole throats identified as light-like parton orbits as boundaries. Since there is a magnetic monopole flux through the wormhole contacts they must appear as pairs (also larger number is possible) in order that magnetic field lines can close. Elementary particles correspond to pairs of wormhole contacts. At both space-time sheets the throats are connected by magnetic flux tubes carrying monopole flux so that a closed flux tube results having a shape of an extremely flattened square and having wormhole contacts at its ends. It is a matter of taste, whether to call the light-like wormhole throats or their interiors as lines of the generalized Feynman/twistor diagrams.

The light-like orbits of partonic 2-surfaces bring strongly in mind the light-like 3-surfaces along which radiation fields can be restricted - kind of shockwaves at which the signature of the induced space-time metric changes its signature.

- (b) String world sheets as orbits of strings are also in an essential role and could be seen as particle like objects. String world sheets could as kind of singular solutions of field equations analogous to characteristics of hyperbolic differential equations. The isometry currents of Kähler action flow along string world sheets and field equations restricted to them are satisfied. As if one would have 2-dimensional solution.  $\sqrt{g_4}$  would of

course vanishes for genuinely 2-D solution but this one can argue that this is not a problem since  $\sqrt{g_4}$  can be eliminated from field equations. String world sheets could serve as 2-D analogs for a solution of hyperbolic field equations defining expanding wave front localized at 3-D light-like surface.

- (c) Propagation in the third sense of word is assignable to the ends of string world sheets at the light-like orbits of partonic 2-surfaces and possibly carrying fermion number. One could say that in TGD one has both fundamental fermions serving as building bricks of elementary particles and strings characterizing interactions between particles. Fermion lines are massless in 8-D sense. By strong form of holography this quantum description has 4-D description space-time description as a classical dual.

2. The topological description of interaction vertices brings in the most important deviation from the standard picture behind cobordism: space-time surfaces are not smooth in TGD framework. One allows topological analogs of 3-vertices of Feynman diagrams realized by connecting three 4-surfaces along their smooth 3-D ends. 3-vertex is also an analog (actually much more!) for the replication in biology. This vertex is *not* the analog of stringy trouser vertex for which space-time surface is continuous whereas 3-surface at the vertex is singular (also trouser vertex could appear in TGD).

The analog of trouser vertex for string world sheets means splitting of string and fermionic field modes decompose into superposition of modes propagating along the two branches. For instance, the propagation of photon along two paths could correspond to its geometric decay at trouser vertex not identifiable as “decay” to two separate particles.

For the analog of 3-vertex of Feynman diagram the 3-surface at the vertex is non-singular but space-time surface is singular. The gluing along ends corresponds to genuine 3-particle vertex.

The view about solution of PDEs generalizes dramatically but the general idea about cobordism might make sense also in the generalized context.

## 7.5 Could Dynamical Homotopy/Homology Groups Characterize WCW Topology?

The challenge is to at least formulate (with my technical background one cannot dream of much more) the analog of cobordism theory in this framework. One can actually hope even the analog of homotopy/homology theory.

1. To a given 3-surface one can assign its cobordism class as the set of 3-surfaces at the opposite boundary of CD connected by a preferred extremal. The 3-surfaces in the same cobordism class are characterized by same conserved classical Noether charges, which become analogs of topological invariants.

One can also consider generalization of cobordisms as analogs to homotopies by allowing return from the opposite boundary of CD. This would give rise

to first homotopy groupoid. One can even go back and forth several times. These dynamical cobordisms allow to divide 3-surfaces at given boundary of CD in equivalence classes characterized among other things by same values of conserved charges. One can also return to the original 3-surface. This could give rise to the analog of the first homotopy group  $\Pi_1$ .

2. If one takes the homotopy interpretation literally one must conclude that the 3-surfaces with different conserved Noether charges cannot be connected by any path in WCW - they belong to disjoint components of the WCW! The zeroth dynamical homotopy group  $\Pi_0$  of WCW would be non-trivial and its elements would be labelled by the conserved Noether charges defining topological invariants!

The values of the classical Noether charges would label disjoint components of WCW. The topology for the space of these parameters would be totally disconnected - no two points cannot be connected by a continuous path. p-Adic topologies are indeed totally disconnected. Could it be that p-adic topology is natural for the conserved classical Noether charges and the sectors of WCW are characterized by p-adic number fields and their algebraic extensions?

Long time ago I noticed that the 4-D spin glass degeneracy induced by the huge vacuum degeneracy of Kähler action implies analogy between the space of maxima of Kähler function and the energy landscape of spin glass systems [K15]. Ultrametricity (<http://tinyurl.com/y6vswdoh>) is the basic property of the topology of the spin glass energy landscape. p-Adic topology is ultrametric and the proposal was that the effective topology for the space of maxima could be p-adic.

3. Isometry charges are the most important Noether charges. These Noether charges are very probably not the only conserved charges. Also the generators in the complement of the gauge sub-algebra of symplectic algebra acting as gauge conformal symmetries could be conserved. All these conserved Noether charges would define a parameter space with a natural p-adic topology.

Since integration is problematic p-adically, one can ask whether only discrete quantum superpositions of 3-surfaces with different classical charges are allowed or whether one should even assume fixed values for the total classical Noether charges appearing in the scattering amplitudes.

I have proposed this kind of approach for the zero modes of WCW geometry not contributing to the Kähler metric except as parameters. The integration for zero modes is also problematic because there is no metric, which would define the integration measure. Since classical charges do not correspond to quantum fluctuating degrees of freedom they should correspond to zero modes. Hence these arguments are equivalent.

The above argument led to the identification of the analogs of the homotopy group  $\Pi_0$  and led to the idea about homotopy groupoid/group  $\Pi_1$ . The elements of  $\Pi_1$  would correspond to space-time surfaces, which run arbitrary number of times fourth and back and return to the initial 3-surface at the boundary of CD. If the

two preferred extremals connecting same pair of 3-surfaces can be deformed to each other, one can say that they are equivalent as dynamical homotopies (or cobordisms). What could be the allowed deformations? Are they cobordisms of cobordisms? What this could mean? Could they define the analog of homotopy groupoid  $\Pi_2$  as foliations of preferred extremals connecting the same 3-surfaces?

1. The number theoretic vision about generalized Feynman diagrams suggests a possible approach. Number theoretic ideas combined with the generalization of twistor approach [K27, K23] led to the vision that generalized Feynman graphs can be identified as sequences or webs of algebraic operations in the co-algebra defined by the Yangian assignable to super-symplectic algebra [A5] [?, ?, ?] and acting as symmetries of TGD. Generalized Feynman graphs would represent algebraic computations. Computations can be done in very many different ways and each of them corresponds to a generalized Feynman diagram. These computations transform give same final collection of “numbers” when the initial collection of “numbers” is given. Does this mean that the corresponding scattering amplitudes must be identical?

If so, a huge generalization of the duality symmetry of the hadronic string models would suggest itself. All computations can be reduced to minimal computations. Accordingly, generalized Feynman diagrams can be reduced to trees by eliminating loops by moving the ends of the loops to same point and snipping the resulting tadpole out! The snipped of tadpole would give a mere multiplicative factor to the amplitude contributing nothing to the scattering rate - just like vacuum bubbles contribute nothing in the case of ordinary Feynman diagrams.

2. How this symmetry could be realized? Could one just assume that only the minimal generalized Feynman diagrams contribute? - not a very attractive option. Or could one hope that only tree diagrams are allowed by the classical dynamics: this was roughly the original vision? The huge vacuum degeneracy of Kähler action implying non-determinism does not encourage this option. The most attractive and most predictive realization conforming with the idea about generalized Feynman diagrammatics as arithmetics would be that all the diagrams differing by these moves give the same result. An analogous symmetry has been discovered for twistor diagrams.
3. Suppose one takes seriously the snipping of a tadpole away from diagram as a move, which does not affect the scattering amplitude. Could this move correspond to an allowed elementary cobordism of preferred extremal? If so, scattering amplitudes would have purely topological meaning as representations of the elements of cobordism classes! TGD would indeed be what it was proposed to be but in much deeper sense than I thought originally. This could also conform with the interpretation of classical charges as topological invariants, realize adelic physics at the level of WCW, and conform with the idea about TGD as almost topological QFT and perhaps generalizing it to topological QFT in generalized sense.

4. One can imagine several interpretations for the snipping operation at space-time level. TGD allows a huge classical vacuum degeneracy: all space-time surfaces having Lagrangian manifold of  $CP_2$  as their  $CP_2$  projection are vacuum extremals of Kähler action. Also all  $CP_2$  type extremals having 1-D light-like curve as  $M^4$  projection are vacuum extremals but have non-vanishing Kähler action. This would not matter if one does not have superpositions since multiplicative factors are eliminated in scattering amplitudes. Could the tadpoles correspond to  $CP_2$  type vacuum extremals at space-time level?

There is also an alternative interpretation. In ZEO causal diamonds (CD) form a hierarchy and one can imagine that the sub-CDs of given CD correspond to quantum fluctuations. Could tadpoles be assigned to sub-CDs of CD be considered+

5. In this manner one could perhaps define elements of homotopy groupoid  $\Pi_2$  as foliations preferred extremals with same ends - these would be 5-D surfaces. If one has two such 5-D foliations with the same 4-D ends, one can form the reverse of the other and form a closed surface. This would be analogous to a map of  $S^2$  to WCW. If the two 5-D foliations cannot be transformed to each other, one would have something, which might be regarded as a non-trivial element of dynamical homotopy group  $\Pi_2$ .

One can ask whether one could define also the analogs of higher homology or homotopy groupoids and groupoids  $\Pi_3$  up to  $\Pi_5$  - the upper bound  $n = 5 = 8 - 3$  comes from the fact that foliations of foliations.. can have maximum dimension  $D = 8$  and from the dimension of  $D = 3$  of basic objects.

1. One could form a foliation of the foliations of preferred extremals as the element of the homotopy groupoid  $\Pi_3$ . Could allowed moves reduce to the snipping operation for generalized Feynman diagrams but performed along direction characterized by a new foliation parameter.
2. The topology of the zero mode sector of WCW parameterized by fixed values of conserved Noether charges as element of  $\Pi_0$  could be characterized by dynamical homotopy groups  $\Pi_n$ ,  $n = 1, \dots, 5$  - at least partially. These degrees of freedom could correspond to quantum fluctuating degrees of freedom. The Kähler structure of WCW and finite-D analogy suggests that all odd dynamical homotopy groups vanish so that  $\Pi_0$ ,  $\Pi_2$  and  $\Pi_4$  would be the only non-trivial dynamical homotopy groups. The vanishing of  $\Pi_1$  would imply that there is only single minimal generalized Feynman diagram contributing to the scattering amplitude. This also true if Feynman diagrams correspond to arithmetic operations.
3. Whether one should call these groups homotopy groups or homology groups is not obvious. The construction means that the foliations of foliations of ... can be seen as images of spheres suggesting "homotopy". The number of these groups is determined by the dimension of embedding space, which suggests "homology".

4. Clearly, the surfaces defining the dynamical homotopy groups/groupoids would be analogs of branes of M-theory but would be obtained constructing paths of paths of paths... by starting from preferred extremals. The construction of so called  $n$ -groups (<http://tinyurl.com/yckcjc1n>) brings strongly in mind this construction.

## **7.6 Appendix: About Field Equations Of TGD In Jet Bundle Formulation**

Prastaro utilizes jet bundle (<http://tinyurl.com/yb2575bm>) formulation of partial differential equations (PDEs). This notion allows a very terse formulation of general PDEs as compared to the old-fashioned but much more concrete formulation that I have used. The formulation is rather formula rich and reader might lose easily his/her patience since one must do hard work to learn which formulas follow trivially from the basic definitions.

I will describe this formulation in TGD framework briefly but without explicit field equations, which can be found at [K3]. To my view a representation by using a concrete example is always more reader friendly than the general formulas derived in some reference. I explain my view about the general ideas behind jet bundle formulation with minimal number amount of formulas. The reader can find explicit formulas from the Wikipedia link above.

The basic goal is to have a geometric description of PDE. In TGD framework the geometric picture is of course present from beginning: field patterns as 4-surfaces in field space - somewhat formal geometric objects - are replaced with genuine 4-surfaces in  $M^4 \times CP_2$ .

### **7.6.1 Field equations as conservation laws, Frobenius integrability conditions, and a connection with quaternion analyticity**

The following represents qualitative picture of field equations of TGD trying to emphasize the physical aspects. Also the possibility that Frobenius integrability conditions are satisfied and correspond to quaternion analyticity is discussed.

1. Kähler action is Maxwell action for induced Kähler form and metric expressible in terms of embedding space coordinates and their gradients. Field equations reduce to those for embedding space coordinates defining the primary dynamical variables. By GCI only four of them are independent dynamical variables analogous to classical fields.
2. The solution of field equations can be interpreted as a section in fiber bundle. In TGD the fiber bundle is just the Cartesian product  $X^4 \times CD \times CP_2$  of space-time surface  $X^4$  and causal diamond  $CD \times CP_2$ .  $CD$  is the intersection of future and past directed light-cones having two light-like boundaries, which are cone-like pieces of light-boundary  $\delta M_{\pm}^4 \times CP_2$ . Space-time surface serves as base space and  $CD \times CP_2$  as fiber. Bundle projection  $\Pi$  is the projection to the factor  $X^4$ . Section corresponds to the map  $x \rightarrow h^k(x)$  giving embedding space coordinates as functions of space-time coordinates. Bundle structure is now trivial and rather formal.

By GCI one could also take suitably chosen 4 coordinates of  $CD \times CP_2$  as space-time coordinates, and identify  $CD \times CP_2$  as the fiber bundle. The choice of the base space depends on the character of space-time surface. For instance  $CD$ ,  $CP_2$  or  $M^2 \times S^2$  ( $S^2$  a geodesic sphere of  $CP_2$ ), could define the base space. The bundle projection would be projection from  $CD \times CP_2$  to the base space. Now the fiber bundle structure can be non-trivial and make sense only in some space-time region with same base space.

3. The field equations derived from Kähler action must be satisfied. Even more: one must have a *preferred* extremal of Kähler action. One poses boundary conditions at the 3-D ends of space-time surfaces and at the light-like boundaries of  $CD \times CP_2$ .

One can fix the values of conserved Noether charges at the ends of  $CD$  (total charges are same) and require that the Noether charges associated with a sub-algebra of super-symplectic algebra isomorphic to it and having conformal weights coming as  $n$ -ples of those for the entire algebra, vanish. This would realize the effective 2-dimensionality required by SH. One must pose boundary conditions also at the light-like partonic orbits. So called weak form of electric-magnetic duality is at least part of these boundary conditions.

It seems that one must restrict the conformal weights of the entire algebra to be non-negative  $r \geq 0$  and those of subalgebra to be positive:  $mn > 0$ . The condition that also the commutators of sub-algebra generators with those of the entire algebra give rise to vanishing Noether charges implies that all algebra generators with conformal weight  $m \geq n$  vanish so the dynamical algebra becomes effectively finite-dimensional. This condition generalizes to the action of super-symplectic algebra generators to physical states.

$M^4$  time coordinate cannot have vanishing time derivative  $dm^0/dt$  so that four-momentum is non-vanishing for non-vacuum extremals. For  $CP_2$  coordinates time derivatives  $ds^k/dt$  can vanish and for space-like Minkowski coordinates  $dm^i/dt$  can be assumed to be non-vanishing if  $M^4$  projection is 4-dimensional. For  $CP_2$  coordinates  $ds^k/dt = 0$  implies the vanishing of electric parts of induced gauge fields. The non-vacuum extremals with the largest conformal gauge symmetry (very small  $n$ ) would correspond to cosmic string solutions for which induced gauge fields have only magnetic parts. As  $n$  increases, also electric parts are generated. Situation becomes increasingly dynamical as conformal gauge symmetry is reduced and dynamical conformal symmetry increases.

4. The field equations involve besides embedding space coordinates  $h^k$  also their partial derivatives up to second order. Induced Kähler form and metric involve first partial derivatives  $\partial_\alpha h^k$  and second fundamental form appearing in field equations involves second order partial derivatives  $\partial_\alpha \partial_\beta h^k$ .

Field equations are hydrodynamical, in other worlds represent conservation laws for the Noether currents associated with the isometries of  $M^4 \times CP_2$ . By GCI there are only 4 independent dynamical variables so that the conservation of  $m \leq 4$  isometry currents is enough if chosen to be independent. The

dimension  $m$  of the tangent space spanned by the conserved currents can be smaller than 4. For vacuum extremals one has  $m = 0$  and for massless extremals (MEs)  $m = 1$ ! The conservation of these currents can be also interpreted as an existence of  $m \leq 4$  closed 3-forms defined by the duals of these currents.

5. The hydrodynamical picture suggests that in some situations it might be possible to assign to the conserved currents flow lines of currents even globally. They would define  $m \leq 4$  global coordinates for some subset of conserved currents (4+8 for four-momentum and color quantum numbers). Without additional conditions the individual flow lines are well-defined but do not organize to a coherent hydrodynamic flow but are more like orbits of randomly moving gas particles. To achieve global flow the flow lines must satisfy the condition  $d\phi^A/dx^\mu = k_B^A J_\mu^B$  or  $d\phi^A = k_B^A J^B$  so that one can special of 3-D family of flow lines parallel to  $k_B^A J^B$  at each point - I have considered this kind of possibly in [K3] at detail but the treatment is not so general as in the recent case.

Frobenius integrability conditions (<http://tinyurl.com/yc6apam2>) follow from the condition  $d^2\phi^A = 0 = dk_B^A \wedge J^B + k_B^A dJ^B = 0$  and implies that  $dJ^B$  is in the ideal of exterior algebra generated by the  $J^A$  appearing in  $k_B^A J^B$ . If Frobenius conditions are satisfied, the field equations can define coordinates for which the coordinate lines are along the basis elements for a sub-space of at most 4-D space defined by conserved currents. Of course, the possibility that for preferred extremals there exists  $m \leq 4$  conserved currents satisfying integrability conditions is only a conjecture.

It is quite possible to have  $m < 4$ . For instance for vacuum extremals the currents vanish identically For MEs various currents are parallel and light-like so that only single light-like coordinate can be defined globally as flow lines. For cosmic strings (cartesian products of minimal surfaces  $X^2$  in  $M^4$  and geodesic spheres  $S^2$  in  $CP_2$  4 independent currents exist). This is expected to be true also for the deformations of cosmic strings defining magnetic flux tubes.

6. Cauchy-Riemann conditions in 2-D situation represent a special case of Frobenius conditions. Now the gradients of real and imaginary parts of complex function  $w = w(z) = u + iv$  define two conserved currents by Laplace equations. In TGD isometry currents would be gradients apart from scalar function multipliers and one would have generalization of C-R conditions. In citeallbprefextremals,twistorstory I have considered the possibility that the generalization of Cauchy-Riemann-Fueter conditions [A14, A11] (<http://tinyurl.com/yb8134b5>) could define quaternion analyticity - having many non-equivalent variants - as a defining property of preferred extremals. The integrability conditions for the isometry currents would be the natural physical formulation of CRF conditions. Different variants of CRF conditions would correspond to varying number of independent conserved isometry currents.
7. The problem caused by GCI is that there is infinite number of coordinate choices. How to pick a physically preferred coordinate system? One possible



manner to do this is to use coordinates for the projection of space-time surface to some preferred sub-space of embedding - geodesic manifold is an excellent choice. Only  $M^1 \times X^3$  geodesic manifolds are not possible but these correspond to vacuum extremals.

One could also consider a philosophical principle behind integrability. The variational principle itself could give rise to at least some preferred space-time coordinates in the same manner as TGD based quantum physics would realize finite measurement resolution in terms of inclusions of HFFs in terms of hierarchy of quantum criticalities and fermionic strings connecting partonic 2-surfaces. Frobenius integrability of the isometry currents would define some preferred coordinates. Their number need not be the maximal four however.

For instance, for massless extremals only light-like coordinate corresponding to the light-like momentum is obtained. To this one can however assign another local light-like coordinate uniquely to obtain integrable distribution of planes  $M^2$ . The solution ansatz however defines directly an integrable choice of two pairs of coordinates at embedding space level usable also as space-time coordinates - light-like local direction defining local plane  $M^2$  and polarization direction defining a local plane  $E^2$ . These choices define integrable distributions of orthogonal planes and local hypercomplex and complex coordinates. Pair of analogs of C-R equations is the outcome. I have called these coordinates Hamilton-Jacobi coordinates for  $M^4$ .

8. This picture allows to consider a generalization of the notion of solution of field equation to that of integral manifold (<http://tinyurl.com/yajn7cuz>. If the number of independent isometry currents is smaller than 4 (possibly locally) and the integrability conditions hold true, lower-dimensional sub-manifolds of space-time surface define integral manifolds as kind of lower-dimensional effective solutions. Genuinely lower-dimensional solutions would of course have vanishing  $\sqrt{g_4}$  and vanishing Kähler action.

String world sheets can be regarded as 2-D integral surfaces. Charged (possibly all) weak boson gauge fields vanish at them since otherwise the electromagnetic charge for spinors would not be well-defined. These conditions force string world sheets to be 2-D in the generic case. In special case 4-D space-time region as a whole can satisfy these conditions. Well-definedness of Kähler-Dirac equation [K29, K18] demands that the isometry currents of Kähler action flow along these string world sheets so that one has integral manifold. The integrability conditions would allow  $2 < m \leq n$  integrable flows outside the string world sheets, and at string world sheets one or two isometry currents would vanish so that the flows would give rise 2-D independent sub-flow.

9. The method of characteristics (<http://tinyurl.com/y9dcdydt>) is used to solve hyperbolic partial differential equations by reducing them to ordinary differential equations. The (say 4-D) surface representing the solution in the field space has a foliation using 1-D characteristics. The method is especially simple for linear equations but can work also in the non-linear case. For instance, the expansion of wave front can be described in terms of characteristics representing light rays. It can happen that two characteristics intersect and

a singularity results. This gives rise to physical phenomena like caustics and shock waves.

In TGD framework the flow lines for a given isometry current in the case of an integrable flow would be analogous to characteristics, and one could also have purely geometric counterparts of shockwaves and caustics. The light-like orbits of partonic 2-surface at which the signature of the induced metric changes from Minkowskian to Euclidian might be seen as an example about the analog of wave front in induced geometry. These surfaces serve as carriers of fermion lines in generalized Feynman diagrams. Could one see the particle vertices at which the 4-D space-time surfaces intersect along their ends as analogs of intersections of characteristics - kind of caustics? At these 3-surfaces the isometry currents should be continuous although the space-time surface has "edge".

10. The analogy with ordinary analyticity suggests that it might be possible to interpret string world sheets and partonic 2-surfaces appearing in strong form of holography (SH) as co-dimension 2 surfaces analogous to poles of analytic function in complex plane. Light-like 3-surfaces might be seen as analogs of cuts. The coding of analytic function by its singularities could be seen as analog of SH.

### 7.6.2 Jet bundle formalism

Jet bundle formalism (<http://tinyurl.com/yb2575bm>) is a modern manner to formulate PDEs in a coordinate independent manner emphasizing the local algebraic character of field equations. In TGD framework GCI of course guarantees this automatically. Beside this integrability conditions formulated in terms of Cartan's contact forms are needed.

1. The basic idea is to take the partial derivatives of embedding space coordinates as functions of space-time coordinates as independent variables. This increases the number of independent variables. Their number depends on the degree of the jet defined and for partial differential equation of order  $r$ , for  $n$  dependent variables, and for  $N$  independent variables the number of new degrees of freedom is determined by  $r$ ,  $n$ , and  $N$  - just by counting the total number of various partial derivatives from  $k = 0$  to  $r$ . For  $r = 1$  (first order PDE) it is  $N \times (1 + n)$ .
2. Jet at given space-time point is defined as a Taylor polynomial of the embedding space coordinates as functions of space-time coordinates and is characterized by the partial derivatives at various points treated as independent coordinates analogous to embedding space coordinate. Jet degree  $r$  is characterized by the degree of the Taylor polynomial. One can sum and multiply jets just like Taylor polynomials. Jet bundle assigns to the fiber bundle associated with the solutions of PDE corresponding jet bundle with fiber at each point consisting of jets for the independent variables ( $CD \times CP_2$  coordinates) as functions of the dependent variables (space-time coordinates).

3. The field equations from the variation of Kähler action are second order partial differential equations and in terms of jet coefficients they reduce to local algebraic equations plus integrability conditions. Since TGD is very non-linear one obtains polynomial equations at each point - one for each embedding space coordinate. Their number reduces to four by GCI. The minimum degree of jet bundle is  $r = 2$  if one wants algebraic equations since field equations are second order PDEs.
4. The local algebraic conditions are not enough. One must have also conditions stating that the new independent variables associated with partial derivatives of various order reduces to appropriate multiple partial derivatives of embedding space coordinates. These conditions can be formulated in terms of Cartan's contact forms, whose vanishing states these conditions. For instance, if  $dh^k$  is replaced by independent variable  $u^k$ , the condition  $dh^k - u^k = 0$  is true for the solution surfaces.
5. In TGD framework there are good motivations to break the non-orthodoxy and use 1-jets so that algebraic equations replaced by first order PDEs plus conditions requiring vanishing of contact forms. These equations state the conservation of isometry currents implying that the 3-forms defined by the duals of isometry currents are closed. As found, this formulation reveals in TGD framework the hydrodynamic picture and suggests conditions making the system integrable in Frobenius sense.

## 8 Twistor lift of TGD and WCW geometry

In the following a view about WCW geometry forced by twistor lift of TGD [K23, K2, K19, K24] is summarized. Twistor lift brings to the action a volume term but without breaking conformal invariance and without introducing cosmological constant as a fundamental dimensional dynamical coupling. The proposed construction of the gamma matrices of WCW giving rise to Kähler metric as anti-commutators is now in terms of the Noether super charges associated with the super-symplectic algebra. This I dare to regard as a very important step of progress.

### 8.1 Possible weak points of the earlier vision

To make progress it is wise to try to identify the possible weak points of the earlier vision.

1. The huge vacuum degeneracy of Kähler action [K10] defining the Kähler function of WCW Kähler metric is analogous to gauge degeneracy of Maxwell action and coded by symplectic transformations of  $CP_2$ . It implies that the degeneracy of the metric increases as one approaches vacuum extremals and is maximal for the space-time surfaces representing canonical embeddings of Minkowski space: Kähler action vanishes up to fourth order in deformation. The original interpretation was in terms of 4-D spin glass degeneracy assumed to be induced by quantum degeneracy.

One could however argue that classical non-determinism of Kähler action is not acceptable and that a small term removing the vacuum degeneracy is needed to make the situation mathematically acceptable. There is an obvious candidate: a volume term having an interpretation in terms of cosmological constant. This term however seems to mean the presence of length scale as a fundamental constant and is in conflict with the basic lesson learned from gauge theories teaching that only dimensionless couplings can be allowed.

2. The construction of WCW Kähler metric relies on the hypothesis that the basic result from the construction of loop space geometries [A8] generalizes: the Kähler metric should be essentially unique from the condition that the isometry group is maximal - this guarantees the existence of Riemann connection. For  $D = 3$  this condition is expected to be even stronger than for  $D = 1$ .

The hypothesis is that in zero energy ontology (ZEO) the symplectic group acting at the light-like boundaries of causal diamond (CD) (one has  $CD = cd \times CP_2$ , where  $cd$  is the intersection of future and past directed light-cones) acts as the isometries of the Kähler metric.

It would be enough to identify complexified WCW gamma matrices and define WCW metric in terms of their anti-commutators. The natural proposal is that gamma matrices are expressible as linear combinations of fermionic oscillator operators for second quantized induced spinor fields at space-time surface. One could even ask whether fermionic super charges and conserved fermionic Noether charges are involved with the construction.

The explicit construction of gamma matrices [K29, K18] has however been based on somewhat ad hoc formulas, and what I call effective 2-dimensionality argued to follow from quantum criticality is somewhat questionable as exact notion.

## 8.2 Twistor lift of TGD and ZEO

Twistor lift of TGD and ZEO meant a revolution in the view about WCW geometry and spinor structure.

1. The basic idea is to replace 4-D Kähler action with dimensionally reduced 6-D Kähler for the analog of twistor space of space-time surface. The induction procedure for the spinors would be generalized so that it applies to twistor structure [?]. The twistor structure of the embedding space is identified as the product of twistor spaces  $M^4 \times S^2$  of  $M^4$  and  $SU(3)/U(1) \times U(1)$  of  $CP_2$ . In momentum degrees of freedom the twistor space of  $M^4$  would be the usual  $CP_3$ .

Remarkably,  $M^4$  and  $CP_2$  are the only spaces allowing twistor space with Kähler structure [A9]. In the case of  $M^4$  the Kähler structure is a generalization of that for  $E^4$ . TGD would be unique from the existence of twistor lift. This predicts CP breaking at fundamental level possibly responsible for CP breaking and matter-antimatter asymmetry.

2. One would still have Kähler coupling strength  $\alpha_K$  as the only single dimensionless coupling strength, whose spectrum is dictated by quantum criticality meaning that it is analogous to critical temperature. All coupling constant like parameters would be determined by quantum criticality. Cosmological constant would not be fundamental constant and this makes itself visible also in the concrete expressions for conserved Noether currents. The breaking of the scale invariance removing vacuum degeneracy of 4-D Kähler action would be analogous to spontaneous symmetry breaking and would remove vacuum degeneracy and classical non-determinism.

The volume term would emerge from dimensional reduction required to give for the 6-surface the structure of  $S^2$  bundle having space-time surfaces as base space. Cosmological constant would be determined by dynamics and depend on p-adic length scale depending in the average on length scale of space-time sheet proportional to the cosmic time sense like  $1/a^2$ ,  $a$  cosmic time. This would solve the problem of large cosmological constant and predict extremely small cosmological constant in cosmic scales in the recent cosmology. This suggests that in long length scales one still has spin glass degeneracy realized in terms of many-sheeted space-time.

3. In ZEO 3-surface correspond to a union of 3-surfaces at the ends of space-time surfaces at boundaries of CD. There are many characterizations of quantum criticality.
  - (a) Preferred extremal property and quantum criticality would mean that one has simultaneously an extremal of both 4-D Kähler action and volume term except at singular 2-surfaces identified as string world sheets and their boundaries. In accordance with the universality of quantum critical dynamics, one would have outside singularities local dynamics without dependence on Kähler coupling strength. The interpretation would be as geometric generalization of massless fields also characterizing criticality.
  - (b) Another characterization of preferred extremal is as a space-time surfaces using sub-algebra  $S_m$  of symplectic algebra  $S$  for which generators have conformal weights coming as  $m$ -tuples of those for the full symplectic algebra. Both  $S_m$  and  $[S, S_m]$  would have vanishing Noether charges. For the induced spinor fields analogous condition would hold true. Effectively the infinite number of radial conformal weights of the symplectic algebra associated with the light-like radial coordinate of  $\delta M_{\pm}^4$  would reduce to a finite number.
  - (c) A further characterization would be in terms of  $M^8 - H$  duality [?]. Preferred extremals in  $H$  would be images of of space-time surfaces in  $M^8$  under  $M^8 - H$  duality. The latter would correspond to roots of octonionic polynomials with coefficients in an extension of rationals. Therefore space-time surfaces in  $H$  satisfying field equations plus preferred extremal conditions would correspond to surfaces described by algebraic

equations in  $M^8$ . Algebraic dynamics would be dual to differential dynamics.

- (d) In adelic physics [?, ?] the hierarchy of Planck constants  $h_{eff}/h_0 = n$  with  $n$  having an interpretation as dimensions of Galois group of extension of rationals would define further correlate of quantum criticality. The scaled up Compton lengths proportional to  $h_{eff}$  would characterize the long range fluctuations associated with quantum criticality.

### 8.3 The revised view about WCW metric and spinor structure

In this framework one can take a fresh approach to the construction of the spinor structure and Kähler metric of WCW. The basic vision is rather conservative. Rather than inducing ad hoc formulas for WCW gamma matrices one tries to identify Noether the elements super-algebra as Noether charges containing also the gamma matrices as Noether super charges.

1. The simplest guess is that the algebra generated by fermionic Noether charges  $Q^A$  for symplectic transformations  $h^k \rightarrow h^k + \epsilon j^{Ak}$  assumed to induce isometries of WCW and Noether supercharges  $Q_n$  and their conjugates for the shifts  $\Psi \rightarrow \Psi + \epsilon u_n$ , where  $u_n$  is a solution of the modified Dirac equation, and  $\epsilon$  is Grassmann number are enough to generate algebra containing the gamma matrix algebra.
2. The commutators  $\Gamma_n^A = [Q^A, Q_n]$  are super-charges labelled by  $(A, n)$ . One would like to identify them as gamma matrices of WCW. The problem is that they are labelled by  $(A, n)$  whereas isometry generators are labelled by  $A$  only just as symplectic Noether charges. Do all supercharges  $\Gamma_n^A$  except  $\Gamma_0^A$  corresponding to  $u_0 = constant$  annihilate the physical states so that one would have 1-1 correspondence? This would be analogous to what happens quite generally in super-conformal algebras.
3. The anti-commutators of  $\Gamma_0^A$  would give the components of the Kähler metric. The allowance of singular surfaces having 2-D string world sheets as singularities would give to the metric also stringy component besides 3-D component and possible 0-D components at the ends of string. Metric 2-D property would not be exact as assumed originally.

This construction can be blamed for the lack of explicitness. The general tendency in the development of TGD has been replacement of explicit but somewhat ad hoc formulas with principles. Maybe this reflects to my own ageing and increasing laziness but my own view is that principles are what matter and get abstracted only very slowly. The less formulas, the better!

## 9 Does 4-D action generate lower-dimensional terms dynamically?

The original proposal was that the action defining the preferred extremals is 4-D Kähler action. Later it became obvious that there must be also 2-D string world sheet term present and probably also 1-D term associated with string boundaries at partonic 2-surfaces. The question has been whether these lower-D terms in the action are primary or generated dynamically. By super-conformal symmetry the same question applies to the fermionic part of the action. The recent formulation based on the twistor lift of TGD contains also volume term but the question remains the same.

During years several motivations for the proposal that preferred extremals of action principle including also volume term for twistor lift of Kähler action are minimal surfaces which are singular at 2-D string world sheets and perhaps also at their boundaries.

In particular, quantum criticality would be realized as a minimal surface property realized by holomorphy in suitably generalized sense [?, ?]. The reason is that the holomorphic solutions of minimal surface equations involve no coupling parameters as the universality of the dynamics at quantum criticality demands.

Minimal surface equation would be true apart from possible singular surfaces having dimension  $D = 2, 1, 0$ .  $D = 2$  corresponds to string world sheets and partonic 2-surfaces. If there are 0-D singularities they would be associated with the ends of orbits of partonic 2-surfaces at boundaries of causal diamond (CD). Minimal surfaces are solutions of non-linear variant of massless d'Alembertian having as effective sources the singular surfaces at which d'Alembertian equation fails. The analogy with gauge theories is highly suggestive: singular surfaces would act as sources of massless field.

Strings world sheets seem to be necessary. The basic question is whether the singular surfaces are postulated from the beginning and there is action associated with them or whether they emerge dynamical from 4-D action. One can consider two extreme options.

**Option I:** There is an explicit assignment of action to the singular surfaces from the beginning. A transfer of Noether charges between space-time interior and string world sheets is possible. This kind of transfer process can take place also between string world sheets and their light-like boundaries and happens if the normal derivatives of embedding space coordinates are discontinuous at the singular surface.

**Option II:** No separate action is assigned with the singular surfaces. There could be a transfer of Noether charges between 4-D Kähler and volume degrees of freedom at the singular surfaces causing the failure of minimal surface property in 4-D sense. But could singular surfaces carry Noether currents as 2-D delta function like densities?

This is possible if the discontinuity of the normal derivatives generates a 2-D singular term to the action. Conservation laws require that at string world sheets energy momentum tensor should degenerate to a 2-D tensor parallel to and concentrated at string world sheet. Only 4-D action would be needed - this was actually

the original proposal. Strings and particles would be essentially edges of space-time - this is not possible in GRT. Same could happen also at its boundaries giving rise to point like particles. Super-conformal symmetry would make this possible also in the fermionic sector.

For both options the singular surfaces would provide a concrete topological picture about the scattering process at the level of single space-time surface and telling what happens to the initial state. The question is whether Option I actually reduces to Option II. If the 2-D term is generated to 4-D action dynamically, there is no need to postulate primary 2-D action.

## 9.1 Can Option II generate separate 2-D action dynamically?

The following argument shows that Option II with 4-D primary action can generate dynamically 2-D term into the action so that no primary action need to be assigned with string world sheets.

### 9.1.1 Dimensional hierarchy of surfaces and strong form of holography

String world sheets having light-like boundaries at the light-like orbits of partonic 2-surfaces are certainly needed to realize strong form of holography [K29]. Partonic 2-surfaces emerge automatically as the ends of the orbits of wormhole contacts.

1. There could (but need not) be a separate terms in the primary action corresponding to string world sheets and their boundaries. This hierarchy bringing in mind branes would correspond to the hierarchy of classical number fields formed by reals, complex numbers, quaternions (space-time surface), and octonions (embedding space in  $M^8$ -side of  $M^8$  duality). The tangent - or normal spaces of these surfaces would inherit real, complex, and quaternionic structures as induced structure. The number theoretic interpretation would allow to see these surfaces as images of those surfaces in  $M^8$  mapped to  $H$  by  $M^8 - H$  duality. Therefore it would be natural to assign action to these surfaces.
2. This makes in principle possible the transfer of classical and quantum charges between space-time interior and string world sheets and between from string world sheets to their light-like boundaries. TGD variant of twistor Grassmannian approach [K19, K24] relies on the assumption that the boundaries of string world sheets at partonic orbits carry quantum numbers. Quantum criticality realized in terms of minimal surface property realized holomorphically is central for TGD and one can ask whether it could play a role in the definition of S-matrix and identification of particles as geometric objects.
3. For preferred extremals string world sheets (partonic 2-surfaces) would be complex (co-complex) manifolds in octonionic sense. Minimal surface equations would hold true outside string world sheets. Conservation of various charges would require that the divergences of canonical momentum currents at string world sheet would be equal to the discontinuities of the normal components of the canonical momentum currents in interior. These discontinuities



would correspond to discontinuities of normal derivatives of embedding space coordinates and are acceptable. Similar conditions would hold true at the light-like boundaries of string world sheets at light-like boundaries of parton orbits. String world sheets would not be minimal surfaces and minimal surface property for space-time surface would fail at these surfaces.

Quantum criticality for string world sheets would also correspond to minimal surface property. If this is realized in terms of holomorphy, the field equations for Kähler and volume parts at string world sheets would be satisfied separately and the discontinuities of normal components for the canonical momentum currents in the interior would vanish at string world sheets.

4. The idea about asymptotic states as free particles would suggest that normal components of canonical momentum currents are continuous near the boundaries of CD as boundary conditions at least. The same must be true at the light-like boundaries of string world sheets. Minimal surface property would reduce to the property of being light-like geodesics at light-like partonic 2-surface. If this is not assumed, the orbit is space-like. Even if these conditions are realized, one can imagine the possibility that at string world sheets 4-D minimal surface equation fails and there is transfer of charges between Kähler and volume degrees of freedom (Option II) and therefore breaking of quantum criticality.

If the exchange of Noether charges vanishes everywhere at string world sheets and boundaries, one could argue that they represent independent degrees of freedom and that TGD reduces to string model. The proposed equation for coupling constant evolution however contains a coefficients depending on the total action so that this would not be the case.

5. Assigning action to the lower-D objects requires additional coupling parameters. One should be able to express these parameters in terms of the parameters appearing in 4-D action ( $\alpha_K$  and cosmological constant). For string sheets the action containing cosmological term is 4-D and Kähler action for  $X^2 \times S^2$ , where  $S^2$  is non-dynamical twistor sphere is a good guess. Kähler action gets contributions from  $X^2$  and  $S^2$ . If the 2-D action is generated dynamically as a singular term of 4-D action its coupling parameters are those of 4-D action.
6. There is a temptation to interpret this picture as a realization of strong form of holography (SH) in the sense that one can deduce the space-time surfaces by using data at string world sheets and partonic 2-surfaces and their light-like orbits. The vanishing of normal components of canonical momentum currents would fix the boundary conditions.

If double holography  $D = 4 \rightarrow D = 2 \rightarrow D = 1$  were satisfied it might be even possible to reduce the construction of S-matrix to the proposed variant of twistor Grassmann approach. This need not be the case: p-adic mass calculations rely on p-adic thermodynamics for the excitations of massless particles having  $CP_2$  mass scale and it would seem that the double holography can makes sense for massless states only.

In  $M^8$ -picture [?] the information about space-time surface is coded by a polynomial defined at real line having coefficients in an extension of rationals. This real line for octonions corresponds to the time axis in the rest system rather than light-like orbit as light-like boundary of string world sheet.

### 9.1.2 Stringy quantum criticality?

The original intuition [?] was that there are canonical momentum currents between Kähler and volume degrees of freedom at singular surfaces but no transfer of canonical momenta between interior and string world sheets nor string world sheets and their boundaries. Also string world sheets would be minimal surfaces as also the intuition from string models suggests. Could also the stringy quantum criticality be realized?

1. Some embedding space coordinates  $h^k$  must have discontinuous partial derivatives in directions normal to the string world sheet so that 3-surface has 1-D edge along fermionic string connecting light-like curves at partonic 2-surfaces in both Minkowskian and Euclidian regions. A closed highly flattened rectangle with long and short edges would be associated with closed monopole flux tube in the case of wormhole contact pairs assigned with elementary particles. 3-surfaces would be “edgy” entities and space-time surfaces would have 2-D and 1-D edges. In condensed matter physics these edges would be regarded as defects.
2. Quantum criticality demands that the dynamics of string world sheets and of interior effectively decouple. Same must take place for the dynamics of string world sheets and their boundaries. Decoupling allows also string world sheets to be minimal surfaces as analogs of complex surfaces whereas string world sheet boundaries would be light-like (their deformations are always space-like) so that one obtains both particles and string like objects.
3. By field equations the sums for the divergences of stringy canonical momentum currents and the corresponding singular parts of these currents in the interior must vanish. By quantum criticality in interior the divergences of Kähler and volume terms vanish separately. Same must happen for the sums in case of string world sheets and their boundaries. The discontinuity of normal derivatives implies that the contribution from the normal directions to the divergence reduces to the sum of discontinuities in two normal directions multiplied by 2-D delta function. This contribution is in the general case equal to the divergence of corresponding stringy canonical momentum current but must vanish if one has quantum criticality also at string world sheets and their boundaries.

The separate continuity of Kähler and volume parts of canonical momentum currents would guarantee this but very probably implies the continuity of the induced metric and Kähler form and therefore of normal derivatives so that there would be no singularity. However, the condition that total canonical momentum currents are continuous makes sense, and indeed implies a transfer of various conserved charges between Kähler action and volume degrees of

freedom at string world sheets and their boundaries in normal directions as was conjectured in [?].

4. What about the situation in fermionic degrees of freedom? The action for string world sheet  $X^2$  would be essentially of Kähler action for  $X^2 \times S^2$ , where  $S^2$  is twistor sphere. Since the modified gamma matrices appearing in the modified Dirac equation are determined in terms of canonical momentum densities assignable to the modified Dirac action, there could be similar transfer of charges involved with the fermionic sector and the divergences of Noether charges and super-charges assignable to the volume action are non-vanishing at the singular surfaces. The above mechanism would force decoupling between interior spinors and string world sheets spinors also for the modified Dirac equation since modified gamma matrices are determined by the bosonic action.

**Remark:** There is a delicacy involved with the definition of modified gamma matrices, which for volume term are proportional to the induced gamma matrices (projections of the embedding space gamma matrices to space-time surface). Modified gamma matrices are proportional to the contractions  $T_k^\alpha \Gamma^k$  of canonical momentum densities  $T^{\alpha k} = \partial L / \partial (\partial_\alpha h^k)$  with embedding space gamma matrices  $\Gamma^k$ . To get dimension correctly in the case of volume action one must divide away the factor  $\Lambda / 8\pi G$ . Therefore fermionic super-symplectic currents do not involve this factor as required.

It remains an open question whether the string quantum criticality is realized everywhere or only near the ends of string world sheets near boundaries of CD.

### 9.1.3 String world sheet singularities as infinitely sharp edges and dynamical generation of string world sheet action

The condition that the singularities are 2-D string world sheets forces 1-D edges of 3-surfaces to be infinitely sharp.

Consider an edge at 3-surface. The divergence from the discontinuity contains contributions from two normal coordinates proportional to a delta function for the normal coordinate and coming from the discontinuity. The discontinuity must be however localized to the string rather than 2-surface. There must be present also a delta function for the second normal coordinate. Hence the value of also discontinuity must be infinite. One would have infinitely sharp edge. A concrete example is provided by function  $y = |x|^\alpha$   $\alpha < 1$ . This kind of situation is encountered in Thom's catastrophe theory for the projection of the catastrophe: in this case one has  $\alpha = 1/2$ . This argument generalizes to 3-D case but visualization is possible only as a motion of infinitely sharp edge of 3-surface.

Kähler form and metric are second degree monomials of partial derivatives so that an attractive assumption is that  $g_{\alpha\beta}$ ,  $J_{\alpha\beta}$  and therefore also the components of volume and Kähler energy momentum tensor are continuous. This would allow  $\partial_{n_i} h^k$  to become infinite and change sign at the discontinuity as the idea about infinitely sharp edge suggests. This would reduce the continuity conditions for canonical momentum currents to rather simple form

$$T^{n_i n_j} \Delta \partial_{n_j} h^k = 0 . \quad (9.1)$$

which in turn would give

$$T^{n_i n_j} = 0 \quad (9.2)$$

stating that canonical momentum is conserved but transferred between Kähler and volume degrees of freedom. One would have a condition for a continuous quantity conforming with the intuitive view about boundary conditions due to conservation laws. The condition would state that energy momentum tensor reduces to that for string world sheet at the singularity so that the system becomes effectively 2-D. I have already earlier proposed this condition.

The reduction of 4-D locally to effectively 2-D system raises the question whether any separate action is needed for string world sheets (and their boundaries)? The generated 2-D action would be similar to the proposed 2-D action. By super-conformal symmetry similar generation of 2-D action would take place also in the fermionic degrees of freedom. I have proposed also this option already earlier. This would mean that Option II is enough.

The following gives a more explicit analysis of the singularities. The vanishing on the discontinuity for the sum of normal derivative gives terms with varying degree of divergence. Denote by  $n_i$  resp.  $t_i$  the coordinate indices in the normal resp. tangent space. Suppose that some derivative  $\partial_{n_i} h^k$  become infinite at string. One can introduce degree  $n_D$  of divergence for a quantity appearing as part of canonical momentum current as the degree of the highest monomial consisting of the diverging derivatives  $\partial_{n_i} h^k$  appearing in quantity in question. For the leading term in continuity conditions for canonical momentum currents of total action one should have  $n_D = 2$  to give the required 2-D delta function singularity.

- $\partial_{n_i} h^k$  has  $n_D \leq 1$ . If it is also discontinuous - say changes sign - one has  $n_D = 2$  for  $\Delta \partial_{n_i} h^k$  in direction  $n_i$ .
- One has  $n_D(g_{t_i t_j}) = 0$ ,  $n_D(g_{t_i n_j}) = 1$ ,  $n_D(g_{n_i n_i}) = 2$  and  $n_D(g_{n_i n_j}) = 1$  or 2 for  $i \neq j$ . One has  $n_D(g) = 4$  ( $g = \det(g_{\alpha\beta})$ ). For contravariant metric one has  $n_D(g^{t_i t_j}) = 0$  and  $n_D(g^{n_i n_j}) = n_D(g^{n_i n_j}) = -2$  as is easy to see from the formula for  $g^{\alpha\beta}$  in terms of cofactors.
- Both Kähler and volume terms in canonical momentum current are proportional to  $\sqrt{g}$  with  $n_D(\sqrt{g}) = 2$  having leading term proportional to 2-determinant  $\sqrt{\det(g_{n_i n_j})}$ . In Kähler action the leading term comes from tangent space part  $J_{ij}$  and has  $n_D = -1$  coming from the partial derivative. The remaining parts involving  $J_{t_i n_j}$  or  $J_{n_i n_j}$  have  $n_D < 0$ .
- Consider the behavior of the contribution of volume term to the canonical momentum currents. For  $g^{n_i t_j} \partial_{t_j} h^k \sqrt{g}$  one has  $n_D = 0$  so that this term is finite. For  $g^{n_i n_j} \partial_{n_j} h^k \sqrt{g}$  one has  $n_D \leq 1$  and this term can be infinite as also its discontinuity coming solely from the change of sign for  $\partial_{n_j} h^k$ . If  $\partial_{n_j} h^k$

is infinite and changes sign, one can have  $n_D = 2$  as required by 2-D delta function singularity.

The continuity condition for the canonical momentum current would state the vanishing of  $n_D = 2$  discontinuity but would not imply separate vanishing of discontinuity for Kähler and volume parts of canonical momentum currents - this in accordance with the idea about canonical momentum transfer. If the sign of partial derivative only changes the coefficient of the partial derivative must vanish so that the condition reduces to the condition  $T^{n_i n_j} = 0$  already given for the components of the total energy momentum tensor, which would be continuous by the above assumption.

## 9.2 Kähler calibrations: an idea before its time?

While updating book introductions I was surprised to find that I had talked about so called calibrations of sub-manifolds as something potentially important for TGD and later forgotten the whole idea! A closer examination however demonstrated that I had ended up with the analog of this notion completely independently later as the idea that preferred extremals are minimal surfaces apart from 2-D singular surfaces, where there would be exchange of Noether charges between Kähler and volume degrees of freedom.

1. The original idea that I forgot too soon was that the notion of calibration (see <http://tinyurl.com/y31yead3>) generalizes and could be relevant for TGD. A calibration in Riemann manifold  $M$  means the existence of a  $k$ -form  $\phi$  in  $M$  such that for any orientable  $k$ -D sub-manifold the integral of  $\phi$  over  $M$  equals to its  $k$ -volume in the induced metric. One can say that metric  $k$ -volume reduces to homological  $k$ -volume.

Calibrated  $k$ -manifolds are minimal surfaces in their homology class, in other words their volume is minimal. Kähler calibration is induced by the  $k^{\text{th}}$  power of Kähler form and defines calibrated sub-manifold of real dimension  $2k$ . Calibrated sub-manifolds are in this case precisely the complex sub-manifolds. In the case of  $CP_2$  they would be complex curves (2-surfaces) as has become clear.

2. By the Minkowskian signature of  $M^4$  metric, the generalization of calibrated sub-manifold so that it would apply in  $M^4 \times CP_2$  is non-trivial. Twistor lift of TGD however forces to introduce the generalization of Kähler form in  $M^4$  (responsible for CP breaking and matter antimatter asymmetry) and calibrated manifolds in this case would be naturally analogs of string world sheets and partonic 2-surfaces as minimal surfaces. Cosmic strings are Cartesian products of string world sheets and complex curves of  $CP_2$ . Calibrated manifolds, which do not reduce to Cartesian products of string world sheets and complex surfaces of  $CP_2$  should also exist and are minimal surfaces.

One can also have 2-D calibrated surfaces and they could correspond to string world sheets and partonic 2-surfaces which also play key role in TGD. Even discrete points assignable to partonic 2-surfaces and representing fundamental fermions play a key role and would trivially correspond to calibrated surfaces.

3. Much later I ended up with the identification of preferred extremals as minimal surfaces by totally different route without realizing the possible connection with the generalized calibrations. Twistor lift and the notion of quantum criticality led to the proposal that preferred extremals for the twistor lift of Kähler action containing also volume term are minimal surfaces. Preferred extremals would be separately minimal surfaces and extrema of Kähler action and generalization of complex structure to what I called Hamilton-Jacobi structure would be an essential element. Quantum criticality outside singular surfaces would be realized as decoupling of the two parts of the action. May be all preferred extremals be regarded as calibrated in generalized sense.

If so, the dynamics of preferred extremals would define a homology theory in the sense that each homology class would contain single preferred extremal. TGD would define a generalized topological quantum field theory with conserved Noether charges (in particular rest energy) serving as generalized topological invariants having extremum in the set of topologically equivalent 3-surfaces.

It is interesting to recall that the original proposal for the preferred extremals as absolute minima of Kähler action has transformed during years to a proposal that they are absolute minima of volume action within given homology class and having fixed ends at the boundaries of CD.

4. The experience with  $CP_2$  would suggest that the Kähler structure of  $M^4$  defining the counterpart of form  $\phi$  is unique. There is however infinite number of different closed self-dual Kähler forms of  $M^4$  defining what I have called Hamilton-Jacobi structures. These forms can have subgroups of Poincare group as symmetries. For instance, magnetic flux tubes correspond to given cylindrically symmetry Kähler form. The problem disappears as one realizes that Kähler structures characterize families of preferred extremals rather than  $M^4$  itself.

If the notion of calibration indeed generalizes, one ends up with the same outcome - preferred extremals as minimal surfaces with 2-D string world sheets and partonic 2-surfaces as singularities - from many different directions.

1. Quantum criticality requires that dynamics does not depend on coupling parameters so that extremals must be separately extremals of both volume term and Kähler action and therefore minimal surfaces for which these degrees of freedom decouple except at singular 2-surfaces, where the necessary transfer of Noether charges between two degrees of freedom takes place at these. One ends up with string picture but strings alone are of course not enough. For instance, the dynamical string tension is determined by the dynamics for the twistor lift.
2. Almost topological QFT picture implies the same outcome: topological QFT property fails only at the string world sheets.
3. Discrete coupling constant evolution, vanishing of loop corrections, and number theoretical condition that scattering amplitudes make sense also in p-adic

number fields, requires a representation of scattering amplitudes as sum over resonances realized in terms of string world sheets.

4. In the standard QFT picture about scattering incoming states are solutions of free massless field equations and interaction regions the fields have currents as sources. This picture is realized by the twistor lift of TGD in which the volume action corresponds to geodesic length and Kähler action to Maxwell action and coupling corresponds to a transfer of Noether charges between volume and Kähler degrees of freedom. Massless modes are represented by minimal surfaces arriving inside causal diamond (CD) and minimal surface property fails in the scattering region consisting of string world sheets.
5. Twistor lift forces  $M^4$  to have generalize Kähler form and this in turn strongly suggests a generalization of the notion of calibration. At physics side the implication is the understanding of CP breaking and matter anti-matter asymmetry.
6.  $M^8 - H$  duality requires that the dynamics of space-time surfaces in  $H$  is equivalent with the algebraic dynamics in  $M^8$ . The effective reduction to almost topological dynamics implied by the minimal surface property implies this. String world sheets (partonic 2-surfaces) in  $H$  would be images of complex (co-complex sub-manifolds) of  $X^4 \subset M^8$  in  $H$ . This should allow to understand why the partial derivatives of embedding space coordinates can be discontinuous at these edges/folds but there is no flow between interior and singular surface implying that string world sheets are minimal surfaces (so that one has conformal invariance).

The analogy with foams in 3-D space deserves to be noticed.

1. Foams can be modelled as 2-D minimal surfaces with edges meeting at vertices. TGD space-time could be seen as a dynamically generated foam in 4-D many-sheeted space-time consisting of 2-D minimal surfaces such that also the 4-D complement is a minimal surface. The counterparts for vertices would be light-like curves at light like orbits of partonic 2-surfaces from which several string world sheets can emanate.
2. Can one imagine something more analogous to the usual 3-D foam? Could the light-like orbits of partonic 2-surfaces define an analog of ordinary foam? Could also partonic 2-surfaces have edges consisting of 2-D minimal surfaces joined along edges representing strings connecting fermions inside partonic 2-surface?

For years ago I proposed what I called as symplectic QFT (SQFT) as an analog of conformal QFT and as part of quantum TGD [K4]. SQFT would have symplectic transformations as symmetries, and provide a description for the symplectic dynamics of partonic 2-surfaces. SQFT involves an analog of triangulation at partonic 2-surfaces and Kähler magnetic fluxes associated with them serve as observables. The problem was how to fix this kind of network. Partonic foam could serve as a concrete physical realization for the symplectic network and have fundamental fermions at vertices. The edges at partonic

2-surfaces would be space-like geodesics. The outcome would be a calibration involving objects of all dimensions  $0 \leq D \leq 4$  - a physical analog of homology theory.

## 10 Could metaplectic group have some role in TGD framework?

Metaplectic group appears as a covering group of linear symplectic group  $Sp(2n, F)$  for any number field and its representations can be regarded as analog of spinor representations of the rotation group. Since infinite-D symplectic group of  $\delta M_+^4 \times CP_2$ , where  $\delta M_+^4$  is light-cone boundary, appears as an excellent candidate for the isometries of "world of classical worlds" in zero energy ontology (ZEO), one can ask whether and how the notion of metaplectic group generalizes to TGD framework [K10, ?, K18, K13, K6, K5].

The condition for the existence of metaplectic structure is same as those for the spinor structure and not met in the case of  $CP_2$ . One however expects that also in the case of metaplectic structure the modified metaplectic structure exists is one couples spinors to an odd integer multiple of Kähler gauge potential. For triality 1 representation assignable to quarks one has  $n = 1$ . The fact that the center of  $SU(3)$  is  $Z_3$  suggests that metaplectic group for  $CP_2$  is 3- or 6-fold covering of symplectic group instead of 2-fold covering.

Besides the ordinary representations of  $SL(2, C)$  also the possibly existing analogs of metaplectic representations of  $SL(2, C) = Sp(2, C)$  acting on wave functions at hyperbolic space  $H_3$  at  $a^2 = t^2 - r^2$  hyperboloid of  $M_+^4$  are cosmologically interesting since the many-sheeted space-time in number theoretic vision allows quantum coherence in even cosmological scales and there are indications for periodic red-shift suggests tessellations of  $H_3$  analogous to lattices in  $E^3$  and defined by discrete subgroup of  $Sl(2, C)$ .

### 10.1 Heisenberg group, symplectic group, and metaplectic group

The following gives a brief summary of basics related to Heisenberg group, symplectic group, and metaplectic group.

#### 10.1.1 Heisenberg group

1. The matrix representation of the simplest Heisenberg group <http://tinyurl.com/y2fomegs> is given by matrices

$$\begin{pmatrix} 1 & a & c \\ 0 & 1 & b \\ 0 & 0 & 1 \end{pmatrix} \tag{10.1}$$

A 3-D Lie group is in question. The multiplication for group elements  $(a_1, b_1, c_1)$  and  $(a_2, b_2, c_2)$  is given by  $(a_1, b_1, c_1) \circ (a_2, b_2, c_2) = (a_1 + a_2, b_1 + b_2, c_1 + c_2 - a_1 b_2)$ .



The coefficients (a,b,c) can belong to any ring since the inverse can be expressed using only product and sum as  $(-a, -b, ab - c)$ . In particular, discrete variants of Heisenberg group such as those associated with extensions of rationals, exist. For odd primes one can define Heisenberg group modulo  $p$  as group of order  $p^3$  in finite field  $F_p$ .

$2n + 1$ -D Heisenberg group consists of upper triangular with unit matrix at diagonal.

2. Continuous Heisenberg group is a nilpotent Lie group of dimension  $d = 3$ . Nilpotency means that its Lie algebra elements are nilpotent. The Lie algebra is generated by upper-diagonal matrices and the commutation relations for the Lie algebra basis are  $[X, Y] = Z, [X, Z] = 0, [Y, Z] = 0$ . The coordinate  $X = q$  and differential operator  $Y = p = \hbar \partial_q, Z = i\hbar 1$  satisfying  $[p, q] = i\hbar Id$ , define a concrete representation of the Lie algebra of the simplest 3-D Heisenberg group in the space of functions  $f(q)$ . By introducing  $n$  pairs of coordinates commuting to unit matrix one obtains  $2n + 1$ -D Heisenberg group.

### 10.1.2 Symplectic group

Symplectic group acts as automorphisms of Heisenberg group. Symplectic group leaves acts in function algebra of function  $H(p, q)$  leaving invariant Poisson bracket  $\{H_1, H_2\} = \partial_q H_1 \partial_p H_2 - \partial_q H_2 \partial_p H_1$ . The Poisson bracket  $\{p, q\} = 1$  giving the element of  $J_{p,q} = 1$  symplectic form remaining invariant under symplectic transformations. Exponentiation of any Hamiltonian  $H(p, q)$  acting as Hamiltonian generates symplectic flows. Symplectic group is infinite-D.

3-D linear symplectic group  $Sp(2, F)$  is obtained as a special case. In continuous case Hamiltonians are linear functions of  $p$  and  $q$  so that the action by Poisson bracket is linear. General linear symplectic group  $Sp(2n, F)$  acts in  $2n$ -D space spanned by the analogs of  $(q_p, p_i)$ . When symplectic form is accompanied by complex structure and Kähler form symplectic isometries define a finite-D subgroup of symplectic group. For instance, in case of  $CP_2$  symplectic isometries define group  $SU(3)$ .

### 10.1.3 Metaplectic group

Metaplectic group  $Mp_m(2n, F)$  (see <http://tinyurl.com/y5mpswy8> and <http://tinyurl.com/y4kjys3e>) is an  $m$ -fold covering of the linear symplectic group  $Sp(2n, F)$ . Metaplectic group like also linear symplectic group metaplectic group is defined for all number fields, in particular  $p$ -adic number fields and even adèles. All representations of the metaplectic group are infinite-D (non-compactness is not the only reason: even finite-D non-unitary matrix representations fail to exist).

$Sp(2, R)$  co-incides with a covering group the special linear group  $SL(2, R)$  acting as real Möbius transformations in upper half-plane. Metaplectic group does

not allow finite-D matrix representations and all representations are infinite-dimensional. Metaplectic group can be regarded as  $m$ -fold cover of symplectic group and in Weil representation the cover can be chosen to be 2-fold cover.

The elements for the metaplectic group  $M_2(2, R)$  as 2-fold covering of  $Sp(2, R)$  have representation as pairs  $(g, \epsilon)$  with  $g$  a Möbius transformation represented by matrix  $(a, b; c, d)$  with unit determinant acting as  $z \rightarrow (az + b)/(cz + d)$  and with  $\epsilon(z)^2 = cz + d$ . The product of group elements is given by  $(g_1, \epsilon_1)(g_2, \epsilon_2) = (g_1 g_2, \epsilon)$ ,  $\epsilon(z) = \epsilon_1(g_2(z))\epsilon_2(z)$ . The entities transforming in this manner are not functions but analogous to spinors and one can speak of symplectic spinors.

3. One can generalize the notion of symplectic structure to that of metaplectic structure. The topological conditions (the second Stiefel-Whitney class vanishes) for the existence of metaplectic structure for given symplectic manifold are same as for the spinor structure.

Interestingly, in the case of  $CP_2$  this condition is not satisfied and the problem is circumvented by coupling  $CP_2$  spinors to an odd multiple of Kähler gauge potential giving rise to Kähler form: this is essential for obtain electroweak couplings correctly for the induced spinor structure at space-time surface. Since Kähler form relates so closely to symplectic structure, it is reasonable to expect that also in case of  $CP_2$  ( $CP_{2n}$ ) symplectic spinors exist.

The center of isometry group  $SU(3)$  of  $CP_2$  is  $Z_3$  acting trivial on  $CP_2$  coordinates. The action is analogous to that of Möbius transformations being induced by linear action of  $SU(3)$  on projective coordinates  $(z_1, z_2, z_3)$  and by the projective map such as  $(z_1, z_2, z_3) \rightarrow (z_1/z_3, z_2/z_3, 1)$  in given coordinate patch defined by a choice of two complex coordinates  $(z_i, z_j)$  now  $(z_1/z_3, z_2/z_3)$ . Do symplectic spinors transform like  $CP_2$  spinors under metaplectic action of  $SU(3)$ ?

$CP_2$  spinors with unit coupling to Kähler gauge potential allow triality  $t = \pm 1$  partial impossible without the coupling making possible spinor structure and presumably also metaplectic structure. Does this mean that in the case of  $CP_2$  the metaplectic group must be identified as 3-fold or possibly 6-fold covering of symplectic group. The holonomy group is electroweak  $U(2)$  and acts like  $SU(2) \times U(1)$ . Does holonomy group acts as double covering of  $SO(3)$  and as 3-fold covering of  $U(1)$  giving 6-fold covering of tangent space group  $SO(4)$ ?

## 10.2 Symplectic group in TGD

In TGD the symplectic transformations of  $\delta M_+^4 \times CP_2$ , where  $\delta M_+^4$  is light-cone boundary, and generated by Hamiltonian algebra, are central and act in the "world of classical worlds" (WCW) [K10, ?, K18, K13, K6, K5].

1. WCW is formed by pairs of 3-surfaces with members at opposite boundaries of causal diamond  $CD = cd \times CP_2$  of embedding space  $H = M^4 \times CP_2$ .  $cd$  is causal diamond of  $M^4$  defined as intersection of future and past directed light-cones. The members of the pair are connected by preferred extremal of

action defined by twistor lift of TGD: it is sum of Kähler action and volume term. Preferred extremal is analogous to Bohr orbit.

2. The obvious question is whether also infinite-D symplectic group of  $\delta M_+^4 \times CP_2$  allows metaplectic variant. Second question is how symplectic spinors relate to ordinary spinors. Are ordinary spinors of  $H$  symplectic spinors as one might expect?
3. In TGD the spinors of "world of classical worlds" (WCW) [K10, ?, K18] should have interpretation as symplectic spinors. Spinors of WCW are fermionic Fock states created by quark oscillator operators replacing theta parameters in super-coordinates and in super-spinors of super variant of embedding space  $H$ . Their local composites appear as monomials with vanishing quark number in hermitian super-coordinates of super-variant of  $H$  and in super-quark-spinors of super- $H$  containing only monomials with odd quark number. These super-fields differ from those of standard SUSY since monomials of theta parameters are replaced with monomials of quark oscillator operators and Majorana spinors are not in question.

Infinite-D metaplectic group  $\delta M_+^4 \times CP_2$  should act on WCW spinor fields and the action should be induced from action in  $H$ .

### 10.3 Kac-Moody type approach to representations of symplectic/metaplectic group

Representations of the symplectic/metaplectic group. Kac-Moody type approach is strongly suggested physically. Kac-Moody group has Lie-algebra which is central extension of the Lie-algebra of local gauge transformation. Kac-Moody algebra elements are labelled by elements with conformal weight  $n \in Z$  but also the variant  $n \geq 0$  ("half-algebra" exists as sub-algebra is clear from the commutation relations.

1. Let  $r$  denote the radial light-like coordinate of light-cone boundary  $\delta M_+^4 \times CP_2$ .  $\delta M_+^4 = S^2 \times R^+$  is metrically 2-sphere  $S^2$  and this implies extension of usual conformal invariance for  $S^2$  to conformal invariance localized with respect to  $r$  and explains why 4-D Minkowski space is physically unique.

Radially local conformal transformations  $z \rightarrow f(r, z)$  of light-cone boundary with scaling  $r \rightarrow |df(r, z, zbar)/dz|^{-1} \times r$  in light-cone radial coordinate  $r$  compensating for the conformal scaling factor  $|df(r, z, zbar)/dz|^2$  as isometries of light-cone boundary as also color rotation local with respect to  $r$ . One has radially local  $S = SO(3) \times SU(3)$  as isometries of light-cone boundary. This would serve as the TGD variant of color gauge symmetry.

2. Effective localization of the symplectic algebra of  $S^2 \times CP_2$  with respect to the radial light-like coordinate  $r$ . Denote the radial conformal weight  $h$ .

**Option 1:** Radial waves of form  $r^h$ ,  $h = -1/2 + iy$  (something to do with zeros of zeta) behave like plane waves with wave vector  $y$  for in inner product defined by integration measure  $dr$ . Orthogonal plane-wave basis effectively.

Restriction to causal diamond CD defined as intersection of future and past directed light-cones implies  $r \leq r_{max}$  defining the size of CD and periodic boundary conditions for a discrete basis  $r^h$ . If  $h = -1/2 + iy$  corresponds to a zero of zeta, the size of CD determined by  $r_{max}$  is quantized. For instance,  $\sin(y \ln(r_{max})) = 0$  would imply  $\ln(r_{max}) = n \times \pi/y$ . Also  $\cos(y \ln(r_{max})) = 0$  can be considered.

**Option 2:** One can include the real part of  $h$  to the integration measure of inner product defined as  $d\mu = dr/r$ . This is dimensionless and very natural by scaling invariance. For this choice one has  $h = iy$  and the connection with Riemann zeta is not anymore natural.  $r_{max} = \exp(n \times \pi/y)$  would give periodic boundary conditions.

For  $y = k\pi$  one would have  $r_{max} = \exp(1/k)$ ,  $k$  integer. This conforms with the adelic picture since the infinite-D extension of rationals generated by  $e^{1/k}$  induces finite-D extension of p-adic numbers since  $e^p$  is ordinary p-adic number.

$y = k\pi/\log(p)$  gives  $r_{max} = p^{n/k}$  and one can construct finite-D extensions of rationals allowing roots of  $p$ .

3. Super-symplectic algebra is assumed to have fractal structure . There is a hierarchy of isomorphic super-symplectic sub-algebras  $SSA_n$ ,  $n = 1, 2, \dots$ , for which conformal weights n-multiples of the weights for the entire algebra.

**Option 1:** One would have also conformal weights  $n(-1/2 + iy)$  for these radial waves however inner product using  $d\mu = dr$  as integration measure does reduce to inner product for plane waves but to  $\int r^{-n+1} \exp(in(y_1 - y_2)) du$ ,  $u = \log(r/r_0)$ . This leads out from the original state space. The modification of the integration measure to  $d\mu = r^{(n-1)} dr$  does not seem plausible.

**Option 2:** Identify the conformal weight as  $h = iy$  and include the real part  $-1/2$  to the dimensionless integration measure  $d\mu = dr/r$ . This allows fractal hierarchy  $h = niy$ . This seems to be the only elegant option so that the connection with Riemann zeta seems artificial

This picture leads to some conjectures and questions.

1. Sub-algebra  $SSA_n$  and its commutator with entire algebra SSA represented trivially for physical states. Also classical Noether charges vanish: this gives strong conditions on preferred extremals and makes them analogs of Bohr orbits: only preferred pairs of 3-surfaces at opposite boundaries of CD are connected by preferred extremal. Hierarchy of state spaces is the outcome.

This would be generalization of Super Virasoro conditions for which only the entire algebra would act trivially apart from the scaling generator  $L_0$ .

2. Could the hierarchies of extensions of rationals with dimensions  $n_1|n_2|...$  ( $|$  is for "divides") correspond to hierarchies of inclusions of hyper-finite factors.
3. Could the hierarchies of  $SSA_n$  with  $n_1|n_2|...$  correspond to hierarchies of extensions of extensions of.... of rationals with dimensions  $n_1|n_2|...$

$\delta M_+^4 \times CP_2$  is metrically  $S^2 \times CP_2$  and this leads to some questions.

1. Could one have Kac-Moody type representation of the symplectic algebra of  $S^2 \times CP_2$ , which is radially local and involves central extension? This is physically suggestive.
2. Symplectic isometries of  $S^2 \times CP_2$  local with respect to  $r$  would define a sub-representation.

Hamiltonians products of  $\delta M_+^4 \times CP_2$  Hamiltonians for  $\delta M_+^4$  and  $CP_2$  labelled by angular momentum  $j$  and by the 2 Casimirs of triality  $t = 0$  color representations.

Isometry algebras  $SO(3)$  and  $SU(3)$  are sub-algebras of symplectic algebra determined by Hamiltonians at light-cone boundary in given representation to themselves. There are no higher-D sub-algebras so that one cannot consider hierarchy analogous to the hierarchy of sub-algebras labelled by radial conformal weights as n-multiples of weights of the entire algebra.

This in turn leads to a series of questions concerning what happens if one takes gauge symmetry and Kac-Moody symmetry as its analog as a physical guideline.

1. The metaplectic group of  $SL(2, R)$  has only infinite-D representations but no matrix representations. Can this be true also for the metaplectic representation of infinite-D for  $SO(3) \times SU(3)$  which is compact and allow finite-D unitary ordinary representations.  $SO(3)$  must be lifted to  $SU(2)$  and this is natural for quark spinors.  $SU(3)$  allows only triality  $t = 0$  partial waves.

Since  $SU(3)$  has  $Z_3$  as center one expects that the notion of metaplectic representation in this case generalizes so that one has 3-fold covering of function space instead of 2-fold one. Quark spinors indeed allow  $CP_2$  partial waves which are in  $t = 1$  representations. As already noticed  $CP_2$  allows does not allow metaplectic structure in standard sense but the coupling to the Kähler gauge potential probably makes this possible since the condition for the existence of generalized metaplectic structure is same as for the existence of modified spinor structure.

2. Should one treat all  $S^2$  Hamiltonians with  $l > 1$  as gauge degrees of freedom? A possible interpretation would be in terms of finite measurement resolution and analog of Kac-Moody symmetry acting very much like gauge symmetry representing the finite measurement resolution. Symplectic group would effectively reduce to  $SO(3) \times SU(3)$ . If so, one would have  $SO(3) \times SU(3)$  gauge theory with  $l = 1$  states and spin 1/2 states with color as particles.
3. Only quark triplets and singlets of fermions and color octets of gluons are observed. Without any additional conditions TGD predicts infinite number of spinor harmonics. For  $CP_2$  spinor harmonics there is a correlation for the color quantum numbers and electroweak quantum numbers of spinor harmonic. In QCD the color representation of quark does not however depend on electroweak quantum numbers. Also the masses of spinor harmonics depend on electroweak quantum numbers and are typically very large.

*Remark:* One could of course ask whether quarks could move in different color partial waves but having  $t = 1$ . This however seems rather implausible.

The proposal is that Kac-Moody type generators can be used to build massless states with have correct correlation between color represented as angular momentum like quantum number and electroweak quantum numbers. Could the experimental absence of higher color partial waves be due to the fact the gauge nature of higher excitations of symplectic algebra making higher color partial waves of quarks and leptons gauge degrees of freedom?

4. What about  $l = 1$  states assignable to  $SO(3)$ ? Twistor lift of TGD predicts that also  $M^4$  has analog of Kähler form and induced  $U(1)$  gauge field analogous to induced Kähler form. The physical effects are weak and would be responsible for CP breaking and matter antimatter asymmetry. Could the  $l = 1$  triplet correspond to this  $U(1)$  gauge boson somewhat like  $SU(3)$  octet corresponds to gluon (gluon is identified as pair of quark and antiquark at different positions)?
5. How does this relate to the analog of metaplectic group for  $SO(3) \times SU(3)$ ? What about the central extension of  $SO(3) \times SU(3)$  assignable to spinor representations with weight  $n = 1/2$ . If one adds to the Hamilton associated with rotation generator  $L_z$  around z-axis in  $SO(3)$  and to hyper-charge generator  $Y$  of  $SU(3)$  a constant, one obtains what looks like central extension at the level of Poisson brackets since right hand side of brackets receives an additive constant. In  $SU(3)$  degrees of freedom one can have only  $t = 0$  color partial waves for scalars but for spinors one obtains the  $t = 1$  waves and can say that color partial waves possess and anomalous hyper-charge  $Y$ .

The spectra of  $L_z$  and  $Y$  are shifted but Killing vector fields are not affected. The couplings of isometry generators are changed since there is coupling proportional to Hamiltonian. This does not seem to have have interpretation as a mere gauge transformation since it makes  $t = 1$  color partial waves possible for quarks.

## 10.4 Relationship to modular functions

The metaplectic representations involve in basic form  $Sp(2n, F)$ ,  $F$  any number field.

1.  $n = 1$  is physically special: one has  $Sp(2, C) = SL(2, C)$ , which is double covering of Lorentz group. The so called modular representations giving rise to basic functions appearing in number theory are related to the representations of  $SL(2, C)$  with the condition that  $SL(2, Z)$  or its discrete subgroup (there are infinite number of them) is represented either trivially or mere projective factor. In the representation realizing  $SL(2, C)$  as Möbius transformations  $z \rightarrow (Az + B)/Cz + D$  or upper half-plane one has  $f(z) \rightarrow (Cz + D)^k f(z)$  when  $(A, B; C, D)$  represents element of  $SL(2, Z)$  or its subgroup  $G$ .  $k$  is integer or half integer. One has modular invariance apart from the projective factor.

Although these nodularity conditions apply only to a discrete subgroup of  $SL(2, R)$  they they imply projective invariance of the analytic functions involved so that projectively their support of the function reduces to  $G \setminus H$ ,  $H$  upper complex plane analogous to unit cell. Could this kind of conditions correspond to the proposed analogs of Kac-Moody type gauge conditions proposed for symplectic symmetries of  $\delta M_{\pm}^4 \times CP_2$ ?

2.  $SO(3, 1)$  acts as isometries of the hyperbolic space  $H_3$  identifiable as the hyperboloid  $H_3$  as  $a^2 = t^2 - r^2 = \text{constant}$  surface of future light-cone  $M_+^4$ :  $a$  defines in TGD Lorentz invariant cosmic time and is natural embedding space coordinate in ZEO. Since  $SL(2, Z)$  has infinite number of discrete subgroups, one has infinite number of tessellations of  $H_3$  analogous to lattices in 3-D Euclidian space.

In TGD quantum coherence is possible in even cosmological scales since TGD predicts hierarchy of effective values of Planck constants. Could one have quantum coherent structures represented as tessellations of the hyperboloid? The prediction would be quantization of redshift as reflection of quantization of distances from given point of tessellations to other points. Evidence for this kind of quantization has been observed.

3. Finite measurement resolution suggests consideration of tessellations as discretization of  $H_3$  and assignable to extensions of rationals and also to subgroups of  $SL(2, Z)$ . This would mean discretized wave functions in the tessellation. This would be like wave function for particle in discrete lattice in  $E^3$ . On the other hand, modular functions with projective modular invariance would be analogs for wave functions of particles periodic symmetry implied by lattice but represented projectively.

Could one decompose the representation to products of modular forms as projective representations in coset space  $SL(2, C)/\Gamma$ ,  $\Gamma$  a discrete subgroup of  $SL(2, C)$  and of representations of discrete subgroup corresponding to finite measurement resolution. This would be like representation of wave functions as products of discrete lattice wave function and wave functions in the space of momenta modulo lattice momenta: Fermi sphere would be replaced by the coset space  $SU(2)/G$ .

4. The projective factor  $\epsilon^2(Z) = (Cz + D)^k$  is essential for the projective representation of  $Sp(2, C)$ . Is it possible to generalize this factor acting on upper complex plane to the case of  $H_3$ ? If subgroup  $SO(3)$  is represented projectively, then one can use for  $H_3$  coordinates  $(r, \theta, \phi)$ , such that  $r$  as radius of sphere  $S^2$  remains invariant under  $r$  and  $SO(3)$  acts on the complex coordinate of  $S^2$  transforming linearly under  $SO(1)$  as  $z \rightarrow (Az + B)/(Cz + D)$  so that the projective factor can be identified. These representations would be analogous to modular representations: the discrete subgroup of  $SL(2, C)$  would be replaced with  $SU(2)$ .

It would seem that it must be replaced with  $SU(2)$  as subgroup. Could one generalize the notion of modular form invariant under discrete subgroup of

$SL(2, C)$  so that the discrete subgroup would become discrete subgroup of  $SO(3)$  ( $SU(2)$ ).

Platonic solids are lattices at  $S^2$  and their isometries and finite subgroups  $D(2n)$  appear in McKay correspondence relating discrete subgroups of  $SU(2)$  and ADE Lie groups. Finite measurement resolution as dual interpretation. What about infinite discrete subgroups. Does invariance mean projective  $SU(2)$  invariance (the case when  $n = 0$ )

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