

# Introduction to "Physics as Generalized Number Theory"

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### Abstract

There are two basic approaches to the construction of quantum TGD. The first approach relies on the vision of quantum physics as infinite-dimensional Kähler geometry for the “world of classical worlds” identified as the space of 3-surfaces in certain 8-dimensional space. Essentially a generalization of the Einstein’s geometrization of physics program is in question.

The second vision identifies physics as a generalized number theory and involves three threads: various p-adic physics and their fusion together with real number based physics to a larger structure, the attempt to understand basic physics in terms of classical number fields (in particular, identifying associativity condition as the basic dynamical principle), and infinite primes whose construction is formally analogous to a repeated second quantization of an arithmetic quantum field theory.

#### 1. *p-Adic physics and their fusion with real physics*

The basic technical problems of the fusion of real physics and various p-adic physics to single coherent whole relate to the notion of definite integral both at space-time level, embedding space level and the level of WCW (the “world of classical worlds”). The expressibility of WCW as a union of symmetric spaces leads to a proposal that harmonic analysis of symmetric spaces can be used to define various integrals as sums over Fourier components. This leads to the proposal the p-adic variant of symmetric space is obtained by an algebraic continuation through a common intersection of these spaces, which basically reduces to an algebraic variant of coset space involving algebraic extension of rationals by roots of unity. This brings in the notion of angle measurement resolution coming as  $\Delta\phi = 2\pi/p^n$  for given p-adic prime  $p$ . Also a proposal how one can complete the discrete version of symmetric space to a continuous p-adic version emerges and means that each point is effectively replaced with the p-adic variant of the symmetric space identifiable as a p-adic counterpart of the real discretization volume so that a fractal p-adic variant of symmetric space results.

If the Kähler geometry of WCW is expressible in terms of rational or algebraic functions, it can in principle be continued the p-adic context. One can however consider the possibility that the integrals over partonic 2-surfaces defining flux Hamiltonians exist p-adically as Riemann sums. This requires that the geometries of the partonic 2-surfaces effectively reduce to finite sub-manifold geometries in the discretized version of  $\delta M_+^4 \times CP_2$ . If Kähler action is required to exist p-adically same kind of condition applies to the space-time surfaces themselves. These strong conditions might make sense in the intersection of the real and p-adic worlds assumed to characterize living matter.

#### 2. *TGD and classical number fields*

The basic vision is that the geometry of the infinite-dimensional WCW (“world of classical worlds”) is unique from its mere existence. This leads to its identification as union of symmetric spaces whose Kähler geometries are fixed by generalized conformal symmetries. This fixes space-time dimension and the decomposition  $M^4 \times S$  and the idea is that the symmetries of the Kähler manifold  $S$  make it somehow unique. The motivating observations are that the dimensions of classical number fields are the dimensions of partonic 2-surfaces, space-time surfaces, and embedding space and  $M^8$  can be identified as hyper-octonions- a sub-space of complexified octonions obtained by adding a commuting imaginary unit. This stimulates some questions.

Could one understand  $S = CP_2$  number theoretically in the sense that  $M^8$  and  $H = M^4 \times CP_2$  be in some deep sense equivalent (“number theoretical compactification” or  $M^8 - H$  duality)? Could associativity define the fundamental dynamical principle so that space-time surfaces could be regarded as associative or co-associative (defined properly) sub-manifolds of  $M^8$  or equivalently of  $H$ .

One can indeed define the associative (co-associative) 4-surfaces using octonionic representation of gamma matrices of 8-D spaces as surfaces for which the Kähler-Dirac gamma matrices span an associate (co-associative) sub-space at each point of space-time surface. In fact, only octonionic structure is needed. Also  $M^8 - H$  duality holds true if one assumes that this associative sub-space at each point contains preferred plane of  $M^8$  identifiable as a preferred commutative or co-commutative plane (this condition generalizes to an integral distribution of commutative planes in  $M^8$ ). These planes are parametrized by  $CP_2$  and this leads to  $M^8 - H$  duality.

WCW itself can be identified as the space of 4-D local sub-algebras of the local Clifford algebra of  $M^8$  or  $H$  which are associative or co-associative. An open conjecture is that this characterization of the space-time surfaces is equivalent with the preferred extremal property

of Kähler action with preferred extremal identified as a critical extremal allowing infinite-dimensional algebra of vanishing second variations.

### 3. Infinite primes

The construction of infinite primes is formally analogous to a repeated second quantization of an arithmetic quantum field theory by taking the many particle states of previous level elementary particles at the new level. Besides free many particle states also the analogs of bound states appear. In the representation in terms of polynomials the free states correspond to products of first order polynomials with rational zeros. Bound states correspond to  $n^{\text{th}}$  order polynomials with non-rational but algebraic zeros at the lowest level. At higher levels polynomials depend on several variables.

The construction might allow a generalization to algebraic extensions of rational numbers, and also to classical number fields and their complexifications obtained by adding a commuting imaginary unit. Special class corresponds to hyper-octonionic primes for which the imaginary part of ordinary octonion is multiplied by the commuting imaginary unit so that one obtains a sub-space  $M^8$  with Minkowski signature of metric. Also in this case the basic construction reduces to that for rational or complex rational primes and more complex primes are obtained by acting using elements of the octonionic automorphism group which preserve the complex octonionic integer property.

Can one map infinite primes/integers/rationals to quantum states? Do they have space-time surfaces as correlates? Quantum classical correspondence suggests that if infinite rationals can be mapped to quantum states then the mapping of quantum states to space-time surfaces automatically gives the map to space-time surfaces. The question is therefore whether the mapping to quantum states defined by WCW spinor fields is possible. A natural hypothesis is that number theoretic fermions can be mapped to real fermions and number theoretic bosons to WCW (“world of classical worlds”) Hamiltonians.

The crucial observation is that one can construct infinite hierarchy of rational units by forming ratios of infinite integers such that their ratio equals to one in real sense: the integers have interpretation as positive and negative energy parts of zero energy states. One can generalize the construction to quaternionic and octonionic units. One can construct also sums of these units with complex coefficients using commuting imaginary unit and these sums can be normalized to unity and have interpretation as states in Hilbert space. These units can be assumed to possess well defined standard model quantum numbers. It is possible to map the quantum number combinations of WCW spinor fields to these states. Hence the points of  $M^8$  can be said to have infinitely complex number theoretic anatomy so that quantum states of the universe can be mapped to this anatomy. One could talk about algebraic holography or number theoretic Brahman=Atman identity.

Also the question how infinite primes might relate to the p-adicization program and to the hierarchy of Planck constants is discussed.

## 1 Introduction

Physics as a generalized number theory program involves three threads: various p-adic physics and their fusion together with real number based physics to a larger structure [K27]. the attempt to understand basic physics in terms of classical number fields [K28]. and infinite primes [K26] whose construction is formally analogous to a repeated second quantization of an arithmetic quantum field theory. A common denominator of these approaches is a precise mathematical formulation for the notion of finite measurement resolution, which could be taken as one of the basic guiding principles of quantum TGD and is at quantum level realized in terms of inclusions of hyper-finite factors about which configuration space spinor fields provide an example [K36]. In the following these threads are described briefly. More detailed summaries will be given in separate articles.

### 1.1 P-Adic Physics And Unification Of Real And P-Adic Physics

p-Adic numbers [A31, A22, A23] became a part of TGD through the successes of p-adic thermodynamics in the description of elementary particle massivation [K18]. The p-adicization program attempts to construct physics in various number fields as an algebraic continuation of physics in the field of rationals (or appropriate extension of rationals). The program involves in an essential

way the generalization of number concept obtained by fusing reals and p-adic number fields to a larger structure by gluing them together along common rationals.

The program involves in an essential way the generalization of number concept obtained by fusing reals and p-adic number fields to a larger structure by gluing them together along common rationals or their algebraic extension. The resulting structure is a generalization of adeles by fusing reals and various p-adic number fields to a book-like structure with pages defined by the number fields glued together along rationals or their algebraic extension in which case the extension induces the extension of p-adic number fields. This structure in turn induces similar structure for embedding spaces, space-time surfaces, and even WCW.

### 1.1.1 Real and p-adic regions of the space-time as geometric correlates of matter and mind

One could end up with p-adic space-time sheets via field equations. The solutions of the equations determining space-time surfaces are restricted by the requirement that the coordinates are real. When this is not the case, one might apply instead of a real completion with some p-adic completion. It however seems that p-adicity is present at deeper level and automatically present via the generalization of the number concept obtained by fusing reals and p-adics along rationals and common algebraics.

p-Adic non-determinism due to the presence of non-constant functions with a vanishing derivative implies extreme flexibility and therefore suggests the identification of the p-adic regions as seats of cognitive representations. Unlike the completion of reals to complex numbers, the completions of p-adic numbers preserve the information about the algebraic extension of rationals and algebraic coding of quantum numbers must be associated with “mind like” regions of space-time. p-Adics and reals would be in the same relationship as map and territory.

The implications are far-reaching and consistent with TGD inspired theory of consciousness: p-adic regions are present even at elementary particle level and provide some kind of model of “self” and of external world. In fact, p-adic physics would model the p-adic cognitive regions representing real elementary particle regions rather than elementary particles themselves! p-Adic mass calculations would be a model of a model!

### 1.1.2 The generalization of the notion of number

The unification of real physics of material work and p-adic physics of cognition and intentionality leads to the generalization of the notion of number field. Reals and various p-adic number fields are glued along their common rationals (and common algebraic numbers too) to form a fractal book like structure. Allowing all possible finite-dimensional extensions of p-adic numbers brings additional pages to this “Big Book”.

At space-time level the book like structure corresponds to the decomposition of space-time surface to real and p-adic space-time sheets glued together along the common back. What this back means is however not what comes first in mind: a subset of space-time points for which preferred embedding space coordinates in an algebraic extension or rationals. This would lead to serious problems with GCI. One must define the intersection of realities and p-adicities at the level of WCW, and demand that the intersection corresponds to space-time surfaces with parameters (WCW coordinats) in the algebraic extension of rationals. The strong form of holography allows to construct space-time surface from string world sheets and partonic 2-surfaces serving as “space-time genes”, and the parameters correspond by conformal invariance to general coordinate invariant conformal moduli for these 2-surfaces. The adelization of TGD reduces to an algebraic continuation of the moduli and various quantum numbers to various number fields.

This has deep implications for the view about cognition. For instance, two points infinitesimally near p-adically are infinitely distant in real sense so that cognition becomes a cosmic phenomenon.

### 1.1.3 Zero energy ontology, cognition, and intentionality

One could argue that conservation laws forbid p-adic-real phase transitions in practice so that cognitions (intentions) realized as real-to-padic (p-adic-to-real) transitions would not be possible. The situation changes if one accepts zero energy ontology [K7, K6].

### 1. Zero energy ontology classically

In TGD inspired cosmology [K25] the imbeddings of Robertson-Walker cosmologies are vacuum extremals. Same applies to the imbeddings of Reissner-Nordström solution [K32] and in practice to all solutions of Einstein's equations imbeddable as extremals of Kähler action. Since four-momentum currents define a collection of vector fields rather than a tensor in TGD, both positive and negative signs for energy corresponding to two possible assignments of the arrow of the geometric time to a given space-time surface are possible. This leads to the view that all physical states have vanishing net energy classically and that physically acceptable universes are creatable from vacuum.

The result is highly desirable since one can avoid unpleasant questions such as “What are the net values of conserved quantities like rest mass, baryon number, lepton number, and electric charge for the entire universe?”, “What were the initial conditions in the big bang?”, “If only single solution of field equations is selected, isn't the notion of physical theory meaningless since in principle it is not possible to compare solutions of the theory?”. This picture fits also nicely with the view that entire universe understood as quantum counterpart 4-D space-time is recreated in each quantum jump and allows to understand evolution as a process of continual re-creation.

### 2. Zero energy ontology at quantum level

The construction of S-matrix [K17, K6] leads to the conclusion that all physical states identified as zero energy states in ZEO possess vanishing conserved quantum numbers but that for a given zero energy state one can identify opposite quantum numbers to the opposite boundaries of causal diamond (CD). Note that ZEO also superposition of states with different conserved quantum numbers at given boundary: this would allow a more natural understanding of Bose-Einstein condensate of Cooper pairs.

Furthermore, the entanglement coefficients between positive and negative energy components of the state have interpretation as  $M$ -matrix identifiable as a “complex square root” of density matrix expressible as a product of positive diagonal square root of the density matrix and of a unitary S-matrix. S-matrix thus becomes a property of the zero energy state and physical states code by their structure what is usually identified as quantum dynamics.

The collection of  $M$ -matrices defines an orthonormal state basis for zero energy states and together they define unitary  $U$ -matrix characterizing transition amplitudes between zero energy states. This matrix would not be however the counterpart of the usual S-matrix. Rather the unitary matrix phase of a given  $M$ -matrix would define the S-matrix measured in laboratory.

At space-time level this would mean that positive energy component and negative energy component are at a temporal distance characterized by the time scale of the causal diamond (CD) and the rational (perhaps integer) characterizing the value of Planck constant for the state in question. The interpretation in terms of a mini bang followed by a mini crunch suggests itself also. CDs are indeed important also in TGD inspired cosmology [K25].

### 3. Hyper-finite factors of type $II_1$ and new view about S-matrix

The representation of S-matrix as unitary entanglement coefficients would not make sense in ordinary quantum theory but in TGD the von Neumann algebra in question is not a type I factor as for quantum mechanics or a type III factor as for quantum field theories, but what is called hyper-finite factor of type  $II_1$  [K36]. This algebra is an infinite-dimensional algebra with the almost defining, and at the first look very strange, property that the infinite-dimensional unit matrix has unit trace. The infinite dimensional Clifford algebra spanned by the configuration space gamma matrices (configuration space understood as the space of 3-surfaces, the “of classical worlds”) is indeed very naturally algebra of this kind since infinite-dimensional Clifford algebras provide a canonical representations for hyper-finite factors of type  $II_1$ .

It has turned out that the fractal structure of HFFs implying hierarchies of Jones inclusions has the hierarchy of quantum criticalities and associated hierarchy of Planck constants  $h_{eff} = n \times h$  as counterparts. Also the hierarchy of algebraic extensions of rationals partially labelled by the integer  $n$  defined by the product of the ramified primes of the extension seems to be closely related to these hierarchies.

### 4. The new view about quantum measurement theory

This mathematical framework leads to a new kind of quantum measurement theory. The basic assumption is that only a finite number of degrees of freedom can be quantum measured in a given measurement and the rest remain untouched. What is known as Jones inclusions  $\mathcal{N} \subset \mathcal{M}$  of von Neumann algebras allow to realize mathematically this idea [K36].  $\mathcal{N}$  characterizes measurement resolution and quantum measurement reduces the entanglement in the non-commutative quantum space  $\mathcal{M}/\mathcal{N}$ . The outcome of the quantum measurement is still represented by a unitary S-matrix but in the space characterized by  $\mathcal{N}$ . It is not possible to end up with a pure state with a finite sequence of quantum measurements.

The obvious objection is that the replacement of a universal S-matrix coding entire physics with a state dependent unitary entanglement matrix is too heavy a price to be paid for the resolution of the above mentioned paradoxes. Situation could be saved if the S-matrices have fractal structure. The quantum criticality of TGD Universe indeed implies fractality. The possibility of an infinite sequence of Jones inclusions for hyperfinite type  $II_1$  factors isomorphic as von Neumann algebras expresses this fractal character algebraically. Thus one can hope that the S-matrix appearing as entanglement coefficients is more or less universal in the same way as Mandelbrot fractal looks more or less the same in all length scales and for all resolutions. Whether this kind of universality must be posed as an additional condition on entanglement coefficients or is an automatic consequence of unitarity in type  $II_1$  sense is an open question.

#### 1.1.4 What number theoretical universality might mean?

Number theoretic universality has been one of the basic guide lines in the construction of quantum TGD. There are two forms of the principle.

1. The strong form of number theoretical universality states that physics for any system should effectively reduce to a physics in algebraic extension of rational numbers at the level of  $M$ -matrix so that an interpretation in both real and p-adic sense (allowing a suitable algebraic extension of p-adics) is possible. One can however worry whether this principle only means that physics is algebraic so that there would be no need to talk about real and p-adic physics at the level of  $M$ -matrix elements. It is not possible to get rid of real and p-adic numbers at the level of classical physics since calculus is a prerequisite for the basic variational principles used to formulate the theory. For this option the possibility of completion is what poses conditions on  $M$ -matrix.
2. The weak form of principle requires only that both real and p-adic variants of physics make sense and that the intersection of these physics consist of physics associated with various algebraic extensions of rational numbers. In this rational physics would be like rational numbers allowing infinite number of algebraic extensions and real numbers and p-adic number fields as its completions. Real and p-adic physics would be completions of rational physics. In this framework criticality with respect to phase transitions changing number field becomes a viable concept. This form of principle allows also purely p-adic phenomena such as p-adic pseudo non-determinism assigned to imagination and cognition. Genuinely p-adic physics does not however allow definition of notions like conserved quantities since the notion of definite integral is lacking and only the purely local form of real physics allows p-adic counterpart.

Strong form of holography suggests a rather elegant and concrete realization of this vision based on string world sheets and partonic 2-surfaces as “space-time genes” and having conformal moduli in an algebraic extension of rationals.

Experience has taught that it is better to avoid too strong statements and perhaps the weak form of the principle is enough. It is however clear that number theoretical criticality could provide important insights to quantum TGD. p-Adic thermodynamics [K18] is an excellent example of this. Needless to say, zero energy ontology is absolutely essential: otherwise this kind of transitions would not make sense.

#### 1.1.5 p-Adicization by algebraic continuation

The basic challenges of the p-adicization program are following.

1. The first problem -the conceptual one- is the identification of preferred coordinates in which functions are algebraic and for which algebraic values of coordinates are in preferred position. This problem is encountered both at the level of space-time, embedding space, and configuration space. Here the group theoretical considerations play decisive role and the selection of preferred coordinates relates closely to the selection of quantization axes. This selection has direct physical correlates at the level of embedding space and the hierarchy of Planck constants has interpretation as a correlate for the selection of quantization axes [K10].

Algebraization does not necessarily mean discretization at space-time level: for instance, the coordinates characterizing partonic 2-surface can be algebraic so that algebraic point of the configuration space results and surface is not discretized. If this kind of function spaces are finite-dimensional, it is possible to fix  $X^2$  completely data for a finite number of points only.

2. Local physics generalizes as such to p-adic context (field equations, etc...). The basic stumbling block of this program is integration already at space-time (Kähler action, flux Hamiltonians, etc...). The problem becomes really horrible looking at configuration space level (functional integral). Algebraic continuation could allow to circumvent this difficulty. Needless to say, the requirement that the continuation exists must pose immensely tight constraints on the physics. Also the existence of the Kähler geometry does this and the solution to the constraint is that WCW is a union of symmetric spaces.

In the case of symmetric spaces Fourier analysis generalizes to harmonics analysis and one can reduce integration to summation for functions allowing Fourier decomposition. In p-adic context the existence of plane waves requires an algebraic extension allowing roots of unity characterizing the measurement accuracy for angle like variables. This leads in the case of symmetric spaces to a general p-adicization recipe. One starts from a discrete variant of the symmetric space for which points correspond to roots of unity and replaces each discrete point with its p-adic completion representing the p-adic variant of the symmetric space so that kind of fractal variant of the symmetric space is obtained. There is an infinite hierarchy of p-adicizations corresponding to measurement resolutions and to the choice of preferred coordinates and the interpretation is in terms of cognitive representations. This requires a refined view about General Coordinate Invariance taking into account the fact that cognition is also part of the quantum state.

One general idea which results as an outcome of the generalized notion of number is the idea of a universal function continuable from a function mapping rationals to rationals or to a finite extension of rationals to a function in any number field. This algebraic continuation is analogous to the analytical continuation of a real analytic function to the complex plane.

1. Rational functions with rational coefficients are obviously functions satisfying this constraint. Algebraic functions with rational coefficients satisfy this requirement if appropriate finite-dimensional algebraic extensions of p-adic numbers are allowed. Exponent function is such a function.
2. For instance, residue calculus essential in the construction of N-point functions of conformal field theory might be generalized so that the value of an integral along the real axis could be calculated by continuing it instead of the complex plane to any number field via its values in the subset of rational numbers forming the rim of the book like structure having number fields as its pages. If the poles of the continued function in the finitely extended number field allow interpretation as real numbers it might be possible to generalize the residue formula. One can also imagine of extending residue calculus to any algebraic extension. An interesting situation arises when the poles correspond to extended p-adic rationals common to different pages of the “great book”. Could this mean that the integral could be calculated at any page having the pole common. In particular, could a p-adic residue integral be calculated in the ordinary complex plane by utilizing the fact that in this case numerical approach makes sense.
3. Algebraic continuation is the basic tool of p-adicization program. Entire physics of the TGD Universe should be algebraically continuable to various number fields. Real number based



physics would define the physics of matter and p-adic physics would describe correlates of cognition.

4. For instance, the idea that number theoretically critical partonic 2-surfaces are expressible in terms of rational functions with rational or algebraic coefficients so that also p-adic variants of these surfaces make sense, is very attractive.
5. Finite sums and products respect algebraic number property and the condition of finiteness is coded naturally by the notion of finite measurement resolution in terms of the notion of (number theoretic) braid. This simplifies dramatically the algebraic continuation since configuration space reduces to a finite-dimensional space and the space of configuration space spinor fields reduces to finite-dimensional function space.

The real configuration space can well contain sectors for which p-adicization does not make sense. For instance, if the exponent of Kähler function and Kähler are not expressible in terms of algebraic functions with rational or at most algebraic functions or more general functions making sense p-adically, the continuation is not possible. p-Adic non-determinism in p-adic sectors makes also impossible the continuation to real sector. All this is consistent with vision about rational and algebraic physics as as analog of rational and algebraic numbers allowing completion to various continuous number fields.

Due to the fact that real and p-adic topologies are fundamentally different, ultraviolet and infrared cutoffs in the set of rationals are unavoidable notions and correspond to a hierarchy of different physical phases on one hand and different levels of cognition on the other hand. For instance, most points p-adic space-time sheets reside at infinity in real sense and p-adically infinitesimal is infinite in real sense. Two types of cutoffs are predicted p-adic length scale cutoff and a cutoff due to phase resolution related to the hierarchy of Planck constants. Zero energy ontology provides natural realization for the p-adic length scale cutoff. The latter cutoff seems to correspond naturally to the hierarchy of algebraic extensions of p-adic numbers and quantum phases  $\exp(i2\pi/n)$ ,  $n \geq 3$ , coming as roots of unity and defining extensions of rationals and p-adics allowing to define p-adically sensible trigonometric functions. These phases relate closely to the hierarchy of quantum groups, braid groups, and  $II_1$  factors of von Neumann algebra.

## 1.2 TGD And Classical Number Fields

This chapter is second one in a multi-chapter devoted to the vision about TGD as a generalized number theory. The basic theme is the role of classical number fields in quantum TGD. A central notion is  $M^8 - H$  duality which might be also called number theoretic compactification. This duality allows to identify embedding space equivalently either as  $M^8$  or  $M^4 \times CP_2$  and explains the symmetries of standard model number theoretically. These number theoretical symmetries induce also the symmetries dictating the geometry of the “world of classical worlds” (WCW) as a union of symmetric spaces. This infinite-dimensional Kähler geometry is expected to be highly unique from the mere requirement of its existence requiring infinite-dimensional symmetries provided by the generalized conformal symmetries of the light-cone boundary  $\delta M^4_+ \times S$  and of light-like 3-surfaces and the answer to the question what makes 8-D embedding space and  $S = CP_2$  so unique would be the reduction of these symmetries to number theory.

Zero energy ontology has become the corner stone of both quantum TGD and number theoretical vision. In zero energy ontology either light-like or space-like 3-surfaces can be identified as the fundamental dynamical objects, and the extension of general coordinate invariance leads to effective 2-dimensionality (strong form of holography) in the sense that the data associated with partonic 2-surfaces and the distribution of 4-D tangent spaces at them located at the light-like boundaries of causal diamonds (CDs) defined as intersections of future and past directed light-cones code for quantum physics and the geometry of WCW.

The basic number theoretical structures are complex numbers, quaternions and octonions, and their complexifications obtained by introducing additional commuting imaginary unit  $\sqrt{-1}$ . Hyper-octonionic (-quaternionic,-complex) sub-spaces for which octonionic imaginary units are multiplied by commuting  $\sqrt{-1}$  have naturally Minkowskian signature of metric. The question is whether and how the hyper-structures could allow to understand quantum TGD in terms of

classical number fields. The answer which looks the most convincing one relies on the existence of octonionic representation of 8-D gamma matrix algebra.

1. The first guess is that associativity condition for the sub-algebras of the local Clifford algebra defined in this way could select 4-D surfaces as associative (hyper-quaternionic) sub-spaces of this algebra and define WCW purely number theoretically. The associative sub-spaces in question would be spanned by the modified gamma matrices defined by the Kähler-Dirac action fixed by the variational principle (Kähler action) selecting space-time surfaces as preferred extremals.
2. This condition is quite not enough: one must strengthen it with the condition that a preferred commutative and thus hyper-complex sub-space is contained in the tangent space of the space-time surface. This condition actually generalizes somewhat since one can introduce a family of so called Hamilton-Jacobi coordinates for  $M^4$  allowing an integrable distribution of decompositions of tangent space to the space of non-physical and physical polarizations [K3]. The physical interpretation is as a number theoretic realization of gauge invariance selecting a preferred local commutative plane of non-physical polarizations.
3. Even this is not yet the whole story: one can define also the notions of co-associativity and co-commutativity applying in the regions of space-time surface with Euclidian signature of the induced metric. The basic unproven conjecture is that the decomposition of space-time surfaces to associative and co-associative regions containing preferred commutative *resp.* co-commutative 2-plane in the 4-D tangent plane is equivalent with the preferred extremal property of Kähler action and the hypothesis that space-time surface allows a slicing by string world sheets and by partonic 2-surfaces.

### 1.2.1 Hyper-octonions and hyper-quaternions

The discussions for years ago with Tony Smith [A39] stimulated very general ideas about space-time surface as an associative, quaternionic sub-manifold of octonionic 8-space (for octonions see [A10]). Also the observation that quaternionic and octonionic primes have norm squared equal to prime in complete accordance with p-adic length scale hypothesis, led to suspect that the notion of primeness for quaternions, and perhaps even for octonions, might be fundamental for the formulation of quantum TGD. The original idea was that space-time surfaces could be regarded as four-surfaces in 8-D embedding space with the property that the tangent spaces of these spaces can be locally regarded as 4- *resp.* 8-dimensional quaternions and octonions.

It took some years to realize that the difficulties related to the realization of Lorentz invariance might be overcome by replacing quaternions and octonions with hyper-quaternions and hyper-octonions. Hyper-quaternions *resp.* -octonions is obtained from the algebra of ordinary quaternions and octonions by multiplying the imaginary part with  $\sqrt{-1}$  and can be regarded as a sub-space of complexified quaternions *resp.* octonions. The transition is the number theoretical counterpart of the transition from Riemannian to pseudo-Riemannian geometry performed already in Special Relativity. The loss of number field and even sub-algebra property is not fatal and has a clear physical meaning. The notion of primeness is inherited from that for complexified quaternions *resp.* octonions.

Note that hyper-variants of number fields make also sense p-adically unlike the notions of number fields themselves unless restricted to be algebraic extensions of rational variants of number fields. What deserves separate emphasis is that the basic structure of the standard model would reduce to number theory.

### 1.2.2 Number theoretical compactification and $M^8 - H$ duality

The notion of hyper-quaternionic and octonionic manifold makes sense but it not plausible that  $H = M^4 \times CP_2$  could be endowed with a hyper-octonionic manifold structure. Situation changes if  $H$  is replaced with hyper-octonionic  $M^8$ . Suppose that  $X^4 \subset M^8$  consists of hyper-quaternionic and co-hyper-quaternionic regions. The basic observation is that the hyper-quaternionic sub-spaces of  $M^8$  with a fixed hyper-complex structure (containing in their tangent space a fixed hyper-complex subspace  $M^2$  or at least one of the light-like lines of  $M^2$ ) are labeled by points of  $CP_2$ .

Hence each hyper-quaternionic and co-hyper-quaternionic four-surface of  $M^8$  defines a 4-surface of  $M^4 \times CP_2$ . One can loosely say that the number-theoretic analog of spontaneous compactification occurs: this of course has nothing to do with dynamics.

This picture was still too naive and it became clear that not all known extremals of Kähler action contain fixed  $M^2 \subset M^4$  or light-like line of  $M^2$  in their tangent space.

1. The first option represents the minimal form of number theoretical compactification.  $M^8$  is interpreted as the tangent space of  $H$ . Only the 4-D tangent spaces of light-like 3-surfaces  $X_l^3$  (wormhole throats or boundaries) are assumed to be hyper-quaternionic or co-hyper-quaternionic and contain fixed  $M^2$  or its light-like line in their tangent space. Hyper-quaternionic regions would naturally correspond to space-time regions with Minkowskian signature of the induced metric and their co-counterparts to the regions for which the signature is Euclidian. What is of special importance is that this assumption solves the problem of identifying the boundary conditions fixing the preferred extremals of Kähler action since in the generic case the intersection of  $M^2$  with the 3-D tangent space of  $X_l^3$  is 1-dimensional. The surfaces  $X^4(X_l^3) \subset M^8$  would be hyper-quaternionic or co-hyper-quaternionic but would not allow a local mapping between the 4-surfaces of  $M^8$  and  $H$ .
2. One can also consider a more local map of  $X^4(X_l^3) \subset H$  to  $X^4(X_l^3) \subset M^8$ . The idea is to allow  $M^2 \subset M^4 \subset M^8$  to vary from point to point so that  $S^2 = SO(3)/SO(2)$  characterizes the local choice of  $M^2$  in the interior of  $X^4$ . This leads to a quite nice view about strong geometric form of  $M^8 - H$  duality in which  $M^8$  is interpreted as tangent space of  $H$  and  $X^4(X_l^3) \subset M^8$  has interpretation as tangent for a curve defined by light-like 3-surfaces at  $X_l^3$  and represented by  $X^4(X_l^3) \subset H$ . Space-time surfaces  $X^4(X_l^3) \subset M^8$  consisting of hyper-quaternionic and co-hyper-quaternionic regions would naturally represent a preferred extremal of  $E^4$  Kähler action. The value of the action would be same as  $CP_2$  Kähler action.  $M^8 - H$  duality would apply also at the induced spinor field and at the level of configuration space.
3. Strong form of  $M^8 - H$  duality satisfies all the needed constraints if it represents Kähler isometry between  $X^4(X_l^3) \subset M^8$  and  $X^4(X_l^3) \subset H$ . This implies that light-like 3-surface is mapped to light-like 3-surface and induced metrics and Kähler forms are identical so that also Kähler action and field equations are identical. The only differences appear at the level of induced spinor fields at the light-like boundaries since due to the fact that gauge potentials are not identical.
4. The map of  $X_l^3 \subset H \rightarrow X_l^3 \subset M^8$  would be crucial for the realization of the number theoretical universality.  $M^8 = M^4 \times E^4$  allows linear coordinates as those preferred coordinates in which the points of embedding space are rational/algebraic. Thus the point of  $X^4 \subset H$  is algebraic if it is mapped to algebraic point of  $M^8$  in number theoretic compactification. This of course restricts the symmetry groups to their rational/algebraic variants but this does not have practical meaning. Number theoretical compactification could thus be motivated by the number theoretical universality.
5. The possibility to use either  $M^8$  or  $H$  picture might be extremely useful for calculational purposes. In particular,  $M^8$  picture based on  $SO(4)$  gluons rather than  $SU(3)$  gluons could perturbative description of low energy hadron physics. The strong  $SO(4)$  symmetry of low energy hadron physics can be indeed seen direct experimental support for the  $M^8 - H$  duality.

### 1.3 Infinite Primes

The notion of prime seems to capture something very essential about what it is to be elementary building block of matter and has become a fundamental conceptual element of TGD. The notion of prime gains its generality from its reducibility to the notion of prime ideal of an algebra. Thus the notion of prime is well-defined, not only in case of quaternions and octonions, but also in the case of hyper-quaternions and -octonions, which are especially natural physically and for which numbers having zero norm correspond physically to light-like 8-vectors. Many interpretations for infinite primes have been competing for survival but it seems that the recent state of TGD allows to exclude some of them from consideration.

### 1.3.1 The notion of infinite prime

Simple arguments show that the  $p$ -adic prime characterizing the 3-surface representing the entire universe increases in a statistical sense in the sequence of quantum jumps: the reason is simply that the size of primes is bounded below. This leads to a peculiar paradox: if the number of quantum jumps already occurred is infinite, this prime is most naturally infinite. On the other hand, if one assumes that only finite number of quantum jumps have occurred, one encounters the problem of understanding why the initial quantum history was what it was. Furthermore, since the size of the 3-surface representing the entire Universe is infinite,  $p$ -adic length scale hypothesis suggest also that the  $p$ -adic prime associated with the entire universe is infinite.

These arguments motivate the attempt to construct a theory of infinite primes and to extend quantum TGD so that also infinite primes are possible. Rather surprisingly, one can construct infinite primes by repeating a procedure analogous to a quantization of a super symmetric arithmetic quantum field theory. At given level of hierarchy one can identify the decomposition of space-time surface to  $p$ -adic regions with the corresponding decomposition of the infinite prime to primes at lower level of infinity: at the basic level are finite primes for which one cannot find any formula.

This and other observations suggest that the Universe of quantum TGD might basically provide a physical representation of number theory allowing also infinite primes. The proposal suggests also a possible generalization of real numbers to a number system akin to hyper-reals introduced by Robinson in his non-standard calculus [A20] providing rigorous mathematical basis for calculus. In fact, some rather natural requirements lead to a unique generalization for the concepts of integer, rational and real. Somewhat surprisingly, infinite integers and reals can be regarded as infinite-dimensional vector spaces with integer and real valued coefficients respectively and this raises the question whether the tangent space for the configuration space of 3-surfaces could be regarded as the space of generalized 8-D hyper-octonionic numbers.

### 1.3.2 Infinite primes and physics in TGD Universe

Several different views about how infinite primes, integers, and rationals might be relevant in TGD Universe have emerged.

#### 1. *Infinite primes, cognition, and intentionality*

The correlation of infinite primes with cognition is the first fascinating possibility and this possibility has stimulated several ideas.

1. The hierarchy of infinite primes associated with algebraic extensions of rationals leading gradually towards algebraic closure of rationals would in turn define cognitive hierarchy corresponding to algebraic extensions of  $p$ -adic numbers.
2. Infinite primes form an infinite hierarchy so that the points of space-time and embedding space can be seen as infinitely structured and able to represent all imaginable algebraic structures. Certainly counter-intuitively, single space-time point -or more generally wave functions in the space of the units associated with the point- might be even capable of representing the quantum state of the entire physical Universe in its structure. For instance, in the real sense surfaces in the space of units correspond to the same real number 1, and single point, which is structure-less in the real sense could represent arbitrarily high-dimensional spaces as unions of real units. For real physics this structure is completely invisible and is relevant only for the physics of cognition. One can say that Universe is an algebraic hologram, and there is an obvious connection both with Brahman=Atman identity of Eastern philosophies and Leibniz's notion of monad.
3. One can assign to infinite primes at  $n^{th}$  level of hierarchy rational functions of  $n$  rational arguments which form a natural hierarchical structure in that highest level corresponds to a polynomial with coefficients which are rational functions of the arguments at the lower level. One can solve one of the arguments in terms of lower ones to get a hierarchy of algebraic extensions. At the lowest level algebraic extensions of rationals emerge, at the next level algebraic extensions of space of rational functions of single variable, etc... This would suggest that infinite primes code for the correlation between quantum states and the algebraic extensions appearing in their their physical description and characterizing their

cognitive correlates. The hierarchy of infinite primes would also correlate with a hierarchy of logics of various orders (hierarchy of statements about statements about...).

### 2. *Infinite primes and super-symmetric quantum field theory*

Consider next the physical interpretation.

1. The discovery of infinite primes suggested strongly the possibility to reduce physics to number theory. The construction of infinite primes can be regarded as a repeated second quantization of a super-symmetric arithmetic quantum field theory. This suggests that configuration space spinor fields or at least the ground states of associated super-conformal representations could be mapped to infinite primes in both bosonic and fermionic degrees of freedom. The process might generalize so that it applies in the case of quaternionic and octonionic primes and their hyper counterparts. This hierarchy of second quantizations means enormous generalization of physics to what might be regarded a physical counterpart for a hierarchy of abstractions about abstractions about.... The ordinary second quantized quantum physics corresponds only to the lowest level infinite primes.
2. The ordinary primes appearing as building blocks of infinite primes at the first level of the hierarchy could be identified as coding for p-adic primes assignable to fermionic and bosonic partons identified as 2-surfaces of a given space-time sheet. The hierarchy of infinite primes would correspond to hierarchy of space-time sheets defined by the topological condensate. This leads also to a precise identification of p-adic and real variants of bosonic partonic 2-surfaces as correlates of intention and action and pairs of p-adic and real fermionic partons as correlates for cognitive representations.
3. The idea that infinite primes characterize quantum states of the entire Universe, perhaps ground states of super-conformal representations, if not all states, could be taken further. It turns out that this idea makes sense when one considers discrete wave functions in the space of infinite primes and that one can indeed represent standard model quantum numbers in this way.
4. The number theoretical supersymmetry suggests also space-time supersymmetry TGD framework. Space-time super-symmetry in its standard form is not possible in TGD Universe and this cheated me to believe that this supersymmetry is completely absent in TGD Universe. The progress in the understanding of the properties of the modified Dirac action however led to a generalization of the space-time super-symmetry as a dynamical and broken symmetry of quantum TGD [?].

Here however emerges the idea about the number theoretic analog of color confinement. Rational (infinite) primes allow not only a decomposition to (infinite) primes of algebraic extensions of rationals but also to algebraic extensions of quaternionic and octonionic (infinite) primes. The physical analog is the decomposition of a particle to its more elementary constituents. This fits nicely with the idea about number theoretic resolution represented as a hierarchy of Galois groups defined by the extensions of rationals and realized at the level of physics in terms of Jones inclusions [K36] defined by these groups having a natural action on space-time surfaces, induced spinor fields, and on configuration space spinor fields representing physical states [K7].

### 3. *Infinite primes and physics as number theory*

The hierarchy of algebraic extensions of rationals implying corresponding extensions of p-adic numbers suggests that Galois groups, which are the basic symmetry groups of number theory, should have concrete physical representations using induced spinor fields and configuration space spinor fields and also infinite primes and real units formed as infinite rationals. These groups permute zeros of polynomials and thus have a concrete physical interpretation both at the level of partonic 2-surfaces dictated by algebraic equations and at the level of braid hierarchy. The vision about the role of hyperfinite factors of  $II_1$  and of Jones inclusions as descriptions of quantum measurements with finite measurement resolution leads to concrete ideas about how these groups are realized.

$G_2$  acts as automorphisms of hyper-octonions and  $SU(3)$  as its subgroup respecting the choice of a preferred imaginary unit. The discrete subgroups of  $SU(3)$  permuting to each other hyper-octonionic primes are analogous to Galois group and turned out to play a crucial role in the understanding of the correspondence between infinite hyper-octonionic primes and quantum states predicted by quantum TGD.

#### 4. *The notion of finite measurement resolution as the key concept*

TGD predicts several hierarchies: the hierarchy of space-time sheets, the hierarchy of infinite primes, the hierarchy of Jones inclusions identifiable in terms of finite measurement resolution [K36]. the dark matter hierarchy characterized by increasing values of  $\hbar$  [K10]. the hierarchy of extensions of a given p-adic number field. TGD inspired theory of consciousness predicts the hierarchy of selves and quantum jumps with increasing duration with respect to geometric time. These hierarchies should be closely related.

The notion of finite measurement resolution turns out to be the key concept: the p-adic norm of the rational defined by the infinite prime characterizes the angle measurement resolution for given p-adic prime  $p$ . It is essential that one has what might be called a state function reduction selecting a fixed p-adic prime which could be also infinite. This gives direct connections with cognition and with the p-adicization program relying also on angle measurement resolution. Also the value the integers characterizing the singular coverings of  $CD$  and  $CP_2$  defining as their product Planck constant characterize the measurement resolution for a given p-adic prime in  $CD$  and  $CP_2$  degrees of freedom. This conforms with the fact that elementary particles are characterized by two infinite primes. Hence finite measurement resolution ties tightly together the three threads of the number theoretic vision. Finite measurement resolution relates also closely to the inclusions of hyper-finite factors central for TGD inspired quantum measurement theory so that the characterization of the finite measurement resolution, which has been the ugly duckling of theoretical physics, transforms to a beautiful swan.

#### 5. *Space-time correlates of infinite primes*

Infinite primes code naturally for Fock states in a hierarchy of super-symmetric arithmetic quantum field theories. Quantum classical correspondence leads to ask whether infinite primes could also code for the space-time surfaces serving as symbolic representations of quantum states. This would a generalization of algebraic geometry would emerge and could reduce the dynamics of Kähler action to algebraic geometry and organize 4-surfaces to a physical hierarchy according to their algebraic complexity. Note that this conjecture should be consistent with two other conjectures about the dynamics of space-time surfaces (space-time surfaces as preferred extrema of Kähler action and space-time surfaces as quaternionic or co-quaternionic (as associative or co-associative) 4-surfaces of hyper-octonion space  $M^8$ ).

The representation of space-time surfaces as algebraic surfaces in  $M^8$  is however too naive idea and the attempt to map hyper-octonionic infinite primes to algebraic surfaces has not led to any concrete progress.

The solution came from quantum classical correspondence, which requires the map of the quantum numbers of configuration space spinor fields to space-time geometry. The Kähler-Dirac equation with measurement interaction term realizes this requirement. Therefore, if one wants to map infinite rationals to space-time geometry it is enough to map infinite primes to quantum numbers. This map can be indeed achieved thanks to the detailed picture about the interpretation of the symmetries of infinite primes in terms of standard model symmetries.

### 1.3.3 Generalization of ordinary number fields: infinite primes and cognition

Both fermions and p-adic space-time sheets are identified as correlates of cognition in TGD Universe. The attempt to relate these two identifications leads to a rather concrete model for how bosonic generators of super-algebras correspond to either real or p-adic space-time sheets (actions and intentions) and fermionic generators to pairs of real space-time sheets and their p-adic variants obtained by algebraic continuation (note the analogy with fermion hole pairs).

The introduction of infinite primes, integers, and rationals leads also to a generalization of classical number fields since an infinite algebra of real (complex, etc...) units defined by finite ratios of infinite rationals multiplied by ordinary rationals which are their inverses becomes possible. These

units are not units in the p-adic sense and have a finite p-adic norm which can differ from one. This construction generalizes also to the case of hyper- quaternions and -octonions although non-commutativity and in case of octonions also non-associativity pose technical problems. Obviously this approach differs from the standard introduction of infinitesimals in the sense that sum is replaced by multiplication meaning that the set of real and also more general units becomes infinitely degenerate.

Infinite primes form an infinite hierarchy so that the points of space-time and imbedding space can be seen as infinitely structured and able to represent all imaginable algebraic structures. Certainly counter-intuitively, single space-time point is even capable of representing the quantum state of the entire physical Universe in its structure. For instance, in the real sense surfaces in the space of units correspond to the same real number 1, and single point, which is structure-less in the real sense could represent arbitrarily high-dimensional spaces as unions of real units.

One might argue that for the real physics this structure is invisible and is relevant only for the physics of cognition. On the other hand, one can consider the possibility of mapping the configuration space and configuration space spinor fields to the number theoretical anatomies of a single point of embedding space so that the structure of this point would code for the world of classical worlds and for the quantum states of the Universe. Quantum jumps would induce changes of configuration space spinor fields interpreted as wave functions in the set of number theoretical anatomies of single point of embedding space in the ordinary sense of the word, and evolution would reduce to the evolution of the structure of a typical space-time point in the system. Physics would reduce to space-time level but in a generalized sense. Universe would be an algebraic hologram, and there is an obvious connection both with Brahman=Atman identity of Eastern philosophies and Leibniz's notion of monad.

Infinite rationals are in one-one correspondence with quantum states and in zero energy ontology hyper-octonionic units identified as ratios of the infinite integers associated with the positive and negative energy parts of the zero energy state define a representation of WCW spinor fields. The action of subgroups of  $SU(3)$  and rotation group  $SU(2)$  preserving hyper-octonionic and hyper-quaternionic primeness and identification of momentum and electro-weak charges in terms of components of hyper-octonionic primes makes this representation unique. Hence Brahman-Atman identity has a completely concrete realization and fixes completely the quantum number spectrum including particle masses and correlations between various quantum numbers.

The appendix of the book gives a summary about basic concepts of TGD with illustrations. Pdf representation of same files serving as a kind of glossary can be found at <http://tgdtheory.fi/tgdglossary.pdf> [L2].

## **2 P-Adic Physics And The Fusion Of Real And P-Adic Physics To A Single Coherent Whole**

In this section basic facts about p-adic numbers [A31, A22, A23] and the question about their relation to real numbers are discussed. Also the basic technicalities related to the notion of p-adic physics are discussed. Also included is a section about the physics in the intersection of real and p-adic worlds relevant to living systems in TGD Universe.

### **2.1 Background**

It is good to start with a summary of the basic mathematical problems related to the p-adicization of physics and a rough formulation for how one might resolve these problems.

#### **2.1.1 Problems**

It is far from obvious what the p-adic counterpart of real physics could mean and how one could fuse together real and p-adic physics. Therefore it is good to list the basic problems and proposals for their solution.

The first problem concerns the correspondence between real and p-adic numbers.

1. The success of p-adic mass calculations involves the notions of p-adic probability, thermodynamics, and the mapping of p-adic probabilities to the real ones by a continuous correspondence  $x = \sum x_n p^n \rightarrow Id(x) = \sum x_n p^{-n}$  that I have christened canonical identification.

The naïve guess is that canonical identification in some form could relate also real and p-adic preferred extremals and define cognitive representations at space-time level. The problem is that  $I$  does not respect symmetries defined by isometries and also general coordinate invariance is possible only if one can identify preferred embedding space coordinates. The reason is that  $I$  does not commute with the basic arithmetic operations.  $I$  allows several variants and it is possible to have correspondence which respects symmetries in arbitrary accuracy in preferred coordinates. Thus  $I$  can play a role at space-time level only if one defines symmetries modulo measurement resolution.  $I$  would make sense only in the interval defining the measurement resolution for a given coordinate variable and the p-adic effective topology would make sense just because the finite measurement resolution does not allow to well-order the points.

2. The identification of real and p-adic numbers via rationals common to all number fields - or more generally along algebraic extension of rationals- respects symmetries and algebra but is not continuous. At the embedding space level preferred coordinates are required also now. The maximal symmetries of the embedding space allow identification of this kind of coordinates. They are not unique. For instance,  $M^4$  linear coordinates look very natural but for  $CP_2$  trigonometric functions of angle like coordinates look more suitable and Fourier analysis suggests strongly the introduction of algebraic extensions involving roots of unity. Partly the non-uniqueness has an interpretation as an embedding space correlate for the selection of the quantization axes. The symmetric space [A16] property of WCW gives hopes that general coordinate invariance in quantal sense can be realized. The existence of p-adic harmonic analysis suggests a discretization of the p-adic variant of embedding space and WCW based on roots of unity.
3. One can consider a compromise between the two correspondences. Discretization via common algebraic points can be completed to a p-adic continuum by assigning to each real discretization interval (say angle increment  $2\pi/N$ ) p-adic numbers with norm smaller than one.
4. It however turned out that more imaginative approach is needed [K35]. Strong form of holography allows to identify string world sheets and partonic 2-surfaces as space-time genes. One can transcend the discretization in an algebraic extension of rationals from space-time level to the level of WCW by demanding that the parameters characterizing these surfaces are in an algebraic extension of rationals. Also cutoffs can be introduced at this level. The outcome is general coordinate invariant (GCI) and problems with symmetries and GCI are avoided. Besides this answers to the basic questions of p-adicization emerge. One can assign to string world sheets purely number theoretically preferred primes and even generalize the p-adic length scale hypothesis using Negentropy Maximization Principle (NMP) [K15].

Second problem relates to integration and Fourier analysis. Both these procedures are fundamental for physics - be it classical or quantum. The p-adic variant of definite integral does not exist in the sense required by the action principles of physics although classical partial differential equations assigned to a particular variational principle make perfect sense. Fourier analysis is also possible only if one allows algebraic extension of p-adic numbers allowing a sufficient number of roots of unity correlating with the measurement resolution of angle. The finite number of them has interpretation in terms of finite angle resolution. Fourier analysis provides also an algebraic realization of definite integral when one integrates over the entire manifold as one indeed does in the case of WCW. If the space in question allows maximal symmetries as WCW and embedding space do, there are excellent hopes of having p-adic variants of both integration and harmonic analysis and the above described procedure allows a precise completion of the discretized variant of real manifold to its continuous p-adic variant.

The third problem relates to the definitions of the p-adic variants of Riemannian, symplectic [A26, A18, A17], and Kähler [A8] geometries. It is possible to generalize formally the notion of Riemann metric although non-local quantities like areas and total curvatures do not make sense if



defined in terms of integrals. If all relevant quantities assignable to the geometry (family of Hamiltonians defining isometries, Killing vector fields, components of metric and Kähler form, Kähler function, etc...) are expressible in terms of rational functions involving only rational numbers as coefficients of polynomials, they allow an algebraic continuation to the p-adic context and the p-adic variant of the geometry makes sense.

The fourth problem relates to the question what one means with p-adic quantum mechanics. In TGD framework p-adic quantum theory utilizes p-adic Hilbert space. The motivation is that the notions of p-adic probability and unitarity are well defined. From the beginning it was clear that the straightforward generalization of Schrödinger equation is not very interesting physically and gradually the conviction has developed that the most realistic approach must rely on the attempt to find the p-adic variant of the TGD inspired quantum physics in all its complexity. The recent approach starts from a rather concrete view about generalized Feynman diagrams defining the points of WCW and leads to a rather detailed view about what the p-adic variants of QM could be and how they could be fused with real QM to a larger structure. Even more, just the requirement that this p-adicization exists, gives very powerful constraints on the real variant of the quantum TGD. Very briefly, algebraic continuation of the scattering amplitudes expressible using data associated with string world sheets and partonic 2-surfaces to various number fields allows to achieve number theoretical universality.

The fifth problem relates to the notion of information in p-adic context. p-Adic thermodynamics leads naturally to the question what p-adic entropy might mean and this in turn leads to the realization that for rational or even algebraic probabilities p-adic variant of Shannon entropy can be negative and has minimum for a unique prime. One can say that the entanglement in the intersection of real and p-adic worlds is negentropic. This leads to rather fascinating vision about how negentropic entanglement (see **Fig. ??** in the Appendix) makes it possible for living systems to overcome the second law of thermodynamics. The formulation of quantum theory in the intersection of real and living worlds becomes the basic challenge.

The proposed solutions to the technical problems could be rephrased in terms of the notion of algebraic universality. Various p-adic physics are obtained as algebraic continuation of real physics through the common algebraic points of real and p-adic worlds and by performing completion in the sense that the interval corresponding to finite measurement resolution are replaced with their p-adic counterpart via canonical identification. This allows to have exact symmetries as their discrete variants and also a continuous correspondence if desired.

### 2.1.2 Program

These ideas lead to a reasonably well defined p-adicization program. Try to define precisely the concepts of the p-adic space-time and “world of classical worlds” (WCW), formulate the finite-p p-adic versions of quantum TGD. Try to fuse together real and various p-adic quantum TGDs are to form a full theory of physics and cognition.

The construction of the p-adic TGD necessitates the generalization of the basic tools of standard physics such as differential and integral calculus, the concept of Hilbert space, Riemannian geometry, group theory, action principles, and the notions of probability and unitarity to the p-adic context. Also new physical thinking and philosophy is needed. The notions of Zero Energy Ontology (ZEO), hierarchy of Planck constants reducible to a hierarchy of quantum criticalities, Negentropy Maximization Principle (NMP), strong form of holography, etc.. are essential but not discussed in detail in the following.

Quite recently it has become clear that strong holography implied by strong form of general coordinate invariance (GCI) is the crux of the construction. WCW has a book-like adelic structure. String world sheets and partonic 2-surfaces serve as number theoretically universal “space-time genes” and induced by algebraic extensions of rationals shared by reals and appropriate extensions of p-adic numbers. This core structure could be called intersection of reality and various p-adicities, the back of the Big Book. What can be said about quantum physics utilizes information about this structure continued algebraically to various real and p-adic sectors.

In the following I try to describe the most central problems and ideas of the p-adicization program. Page number of a readable article must be finite and this has forced to leave away a lot of topics. p-Adic mass calculations [K18], which form the corner stone of the entire approach would require entire article series. The vision about how to define generalized Feynman diagrams and their

p-adic variants by utilizing the assumption that WCW is symmetric space allowing algebraization of functional integral crucial for the entire approach is discussed [L1]. Here huge symmetries of WCW, which include super-symplectic symmetry and generalize the conformal symmetries of string models, are in key role [K11, K8]. Negentropy Maximization Principle [K15] relevant for understanding the profound implications of the negentropic entanglement, in particular how the preferred p-adic primes emerge [K35] is not discussed. The applications of p-adic length scale hypothesis to the physics of living matter [K29] and the model of cognition [K4, K20] would provide additional insights and motivations but have been also left out.

## 2.2 Summary Of The Basic Physical Ideas

In the following various ways to end up with p-adic physics and with the idea about p-adic physics as physics of cognition are discussed. There is also the idea about p-adic topology as an effective topology of real space-time surfaces in finite measurement resolution implying discretization but this idea is not so compelling.

### 2.2.1 p-Adic mass calculations briefly

p-Adic mass calculations based on p-adic thermodynamics with energy replaced with the generator  $L_0 = zd/dz$  of infinitesimal scaling are described in the first part of [K18].

1. p-Adic thermodynamics could be justified by the randomness of the motion of partonic 2-surfaces restricted only by the light-likeness of the orbit.
2. It is essential that the conformal symmetries associated with the light-like coordinates of parton and light-cone boundary are not gauge symmetries but dynamical symmetries. The point is that there are two kinds of super-conformal symmetries [A14, A15]: the super-symplectic conformal symmetries assignable to the light-like boundaries of  $CD \times CP_2$  and super Kac-Moody symmetries [A7] assignable to light-like 3-surfaces defining fundamental dynamical objects. In so called coset construction [A33] the differences of super-conformal generators of these algebras annihilate the physical states. This leads to a generalization of Equivalence Principle since one can assign four-momentum to the generators of both algebras identifiable as inertial *resp.* gravitational four-momentum. A second important consequence is that the generators of either algebra do not act like gauge transformations so that it makes sense to construct p-adic thermodynamics for them.
3. In p-adic thermodynamics scaling generator  $L_0$  having conformal weights as its eigen values replaces energy and Boltzmann weight  $\exp(H/T)$  is replaced by  $p^{L_0/T_p}$ . The quantization  $T_p = 1/n$  of conformal temperature and thus quantization of mass squared scale is implied by number theoretical existence of Boltzmann weights. p-Adic length scale hypothesis states that primes  $p \simeq 2^k$ ,  $k$  integer. A stronger hypothesis is that  $k$  is prime (in particular Mersenne prime or Gaussian Mersenne) makes the model very predictive and fine tuning is not possible.

Mersenne primes are very special number theoretically because bit as the unit of information unit corresponds to  $\log(2)$  and can be said to exist for  $M_n$ -adic topology. The reason is that  $\log(1+p)$  existing always p-adically corresponds for  $M_n = 2^n - 1$  to  $\log(2^n) \equiv n\log(2)$  so that one has  $\log(2 \equiv \log(1 + M_n)/n$ . Since the powers of 2 modulo  $p$  give all integers  $n \in \{1, p-1\}$  by Fermat's theorem, one can say that the logarithms of all integers modulo  $M_n$  exist in this sense and therefore the logarithms of all p-adic integers not divisible by  $p$  exist. For other primes one must introduce a transcendental extension containing  $\log(a)$  where  $a$  is so called primitive root. One could criticize the identification since  $\log(1 + M_n)$  corresponding in the real sense to  $n$  bits corresponds in p-adic sense to a very small information content since the p-adic norm of the p-adic bit is  $1/M_n$ .

The basic mystery number of elementary particle physics defined by the ratio of Planck mass and proton mass follows thus from number theory once  $CP_2$  radius is fixed to about  $10^4$  Planck lengths. Mass scale becomes additional discrete variable of particle physics so that there is not more need to force top quark and neutrinos with mass scales differing by 12 orders of magnitude to the same multiplet of gauge group. Electron, muon, and  $\tau$  correspond to Mersenne prime  $k = 127$

(the largest non-super-astrophysical Mersenne), and Mersenne primes  $k = 113, 107$ . Intermediate gauge bosons and photon correspond to Mersenne  $M_{89}$ , and graviton to  $M_{127}$ .

The value of  $k$  for quark can depend on hadronic environment [K19] and this would produce precise mass formulas for low energy hadrons. This kind of dependence conforms also with the indications that neutrino mass scale depends on environment [?]. Amazingly, the biologically most relevant length scale range between 10 nm and 4  $\mu\text{m}$  contains four Gaussian Mersennes  $(1+i)^n - 1$ ,  $n = 151, 157, 163, 167$  and scaled copies of standard model physics in cell length scale could be an essential aspect of macroscopic quantum coherence prevailing in cell length scale.

p-Adic mass thermodynamics is not quite enough: also Higgs boson is needed and wormhole contact carrying fermion and anti-fermion quantum numbers at the light-like wormhole throats is excellent candidate for Higgs [K13]. The coupling of Higgs to fermions can be small and induce only a small shift of fermion mass: this could explain why Higgs has not been observed. Also the Higgs contribution to mass squared can be understood thermodynamically if identified as absolute value for the thermal expectation value of the eigenvalues of the Kähler-Dirac operator having interpretation as complex square root of conformal weight.

The original belief was that only Higgs corresponds to wormhole contact. The assumption that fermion fields are free in the conformal field theory applying at parton level forces to identify all gauge bosons as wormhole contacts connecting positive and negative energy space-time sheets [K13]. Fermions correspond to topologically condensed  $CP_2$  type extremals with single light-like wormhole throat. Gravitons are identified as string like structures involving pair of fermions or gauge bosons connected by a flux tube. Partonic 2-surfaces are characterized by genus which explains family replication phenomenon and an explanation for why their number is three emerges [K5]. Gauge bosons are labeled by pairs  $(g_1, g_2)$  of handle numbers and can be arranged to octet and singlet representations of the resulting dynamical  $SU(3)$  symmetry. Ordinary gauge bosons are  $SU(3)$  singlets and the heaviness of octet bosons explains why higher boson families are effectively absent. The different character of bosons could also explain why the p-adic temperature for bosons is  $T_p = 1/n < 1$  so that Higgs contribution to the mass dominates.

The basis challenge is to understand why elementary particles seem to be characterized by preferred p-adic primes and why these primes seem to obey p-adic length scale hypothesis- that is be near but below powers of two.

### 2.2.2 p-Adic length scale hypothesis, ZEO, and hierarchy of Planck constants

ZEO and the hierarchy of Planck constants realized in terms of the generalization of the embedding space lead to a deeper understanding of the origin of the p-adic length scale hypothesis.

#### 1. ZEO

In ZEO one replaces positive energy states with zero energy states with positive and negative energy parts of the state at the light-like boundaries of CD. All conserved quantum numbers of the positive and negative energy states are of opposite sign so that these states can be created from vacuum. “Any physical state is creatable from vacuum” becomes thus a basic principle of quantum TGD and together with the notion of quantum jump resolves several philosophical problems (What was the initial state of universe?, What are the values of conserved quantities for Universe?, Is theory building completely useless if only single solution of field equations is realized?). At the level of elementary particle physics positive and negative energy parts of zero energy state are interpreted as initial and final states of a particle reaction so that quantum states become physical events.

At the level of WCW ZEO means that pairs of 3-surfaces residing at opposite boundaries of CD become basis objects or equivalent preferred extremals of Kähler acting [K2] having these 3-surfaces at ends replaced space-like 3-surfaces as basic objects. Preferred extremal property means that these space-time surfaces become archetypal spatiotemporal patterns: biologist would talk about behaviors, functions, or self-organization patterns [K14]. Self-organization is however understood in 4-D sense.

#### 2. Does the finiteness of measurement resolution dictate the laws of physics?

The hypothesis that the mere finiteness of measurement resolution could determine the laws of quantum physics [K6] completely belongs to the category of not at all obvious first principles.

The basic observation is that the Clifford algebra [A2] spanned by the gamma matrices of the “world of classical worlds” represents a von Neumann algebra [A35] known as hyperfinite factor of type II<sub>1</sub> (HFF) [K6, K36, K10]. HFF [A25, A32] is an algebraic fractal having infinite hierarchy of included sub-algebras isomorphic to the algebra itself [A1]. The structure of HFF is closely related to several notions of modern theoretical physics such as integrable statistical physical systems [A40], anyons [?], quantum groups and conformal field theories [A24], and knots and topological quantum field theories [A37, A29].

ZEO is second key element. In ZEO these inclusions allow an interpretation in terms of a finite measurement resolution: in the standard positive energy ontology this interpretation is not possible. Inclusion hierarchy defines in a natural manner the notion of coupling constant evolution and p-adic length scale hypothesis follows as a prediction. In this framework the extremely heavy machinery of renormalized quantum field theory involving the elimination of infinities is replaced by a precisely defined mathematical framework. More concretely, the included algebra creates states which are equivalent in the measurement resolution used. Zero energy state can be modified in a time scale shorter than the time scale of the zero energy state itself.

One can imagine two kinds of measurement resolutions. The element of the included algebra can leave the quantum numbers of the positive and negative energy parts of the state invariant, which means that the action of subalgebra leaves M-matrix invariant. The action of the included algebra can also modify the quantum numbers of the positive and negative energy parts of the state such that the zero energy property is respected. In this case the Hermitian operators subalgebra must commute with M-matrix.

The temporal distance between the tips of CD corresponds to the secondary p-adic time scale  $T_{p,2} = \sqrt{p}T_p$  by a simple argument based on the observation that light-like randomness of light-like 3-surface is analogous to Brownian motion. This gives the relationship  $T_p = L_p^2/Rc$ , where  $R$  is  $CP_2$  size. The action of the included algebra corresponds to an addition of zero energy parts to either positive or negative energy part of the state and is like addition of quantum fluctuation below the time scale of the measurement resolution. The natural hierarchy of time scales is obtained as  $T_n = 2^{-n}T$  since these insertions must belong to either upper or lower half of the causal diamond. This implies that preferred p-adic primes are near powers of 2. For electron the time scale in question is .1 seconds defining the fundamental biorhythm of 10 Hz.

M-matrix representing a generalization of S-matrix and expressible as a product of a positive square root of the density matrix and unitary S-matrix would define the dynamics of quantum theory [K6]. The notion of thermodynamical state would cease to be a theoretical fiction and in a well-defined sense quantum theory could be regarded as a square root of thermodynamics. Connes tensor product [A25] provides a mathematical description of the finite measurement resolution but does not fix the M-matrix as was the original hope. The remaining challenge is the calculation of M-matrix and the progress induced by ZEO during last years has led to rather concrete proposal for the construction of M-matrix.

It turns out however that the mathematical representation for the notion of finite resolution for angle measurement serves as a common denominator for all basic approaches to quantum TGD: the Kähler geometry [A8] of WCW identified as a union of infinite-dimensional symmetric spaces, inclusions of hyper finite factors as representation of finite measurement resolution, p-adicization program, the role of classical number fields [A10, A4, A13], and infinite primes so that it is fair to say that all approaches to TGD which originally seemed almost independent, converge to a coherent mathematical structure.

### 3. How do p-adic coupling constant evolution and p-adic length scale hypothesis emerge?

In zero energy ontology zero energy states have as embedding space correlates causal diamonds for which the distance between the tips of the intersecting future and past directed light-cones comes as integer multiples of a fundamental time scale:  $T_n = n \times T_0$ . p-Adic length scale hypothesis allows to consider a stronger hypothesis  $T_n = 2^n T_0$  and its generalization a slightly more general hypothesis  $T_n = p^n T_0$ ,  $p$  prime. It however seems that these scales are dynamically favored but that also other scales are possible.

Could the coupling constant evolution in powers of 2 implying time scale hierarchy  $T_n = 2^n T_0$  (or  $T_p = p T_0$ ) induce p-adic coupling constant evolution and explain why p-adic length scales correspond to  $L_p \propto \sqrt{p}R$ ,  $p \simeq 2^k$ ,  $R$   $CP_2$  length scale? This looks attractive but there is a

problem. p-Adic length scales come as powers of  $\sqrt{2}$  rather than 2 and the strongly favored values of  $k$  are primes and thus odd so that  $n = k/2$  would be half odd integer. This problem can be solved.

1. The observation that the distance traveled by a Brownian particle during time  $t$  satisfies  $r^2 = Dt$  suggests a solution to the problem. p-Adic thermodynamics applies because the partonic 3-surfaces  $X^2$  are as 2-D dynamical systems random apart from light-likeness of their orbit. For  $CP_2$  type vacuum extremals the situation reduces to that for a one-dimensional random light-like curve in  $M^4$ . The orbits of Brownian particle would now correspond to light-like geodesics  $\gamma_3$  at  $X^3$ . The projection of  $\gamma_3$  to a time=constant section  $X^2 \subset X^3$  would define the 2-D path  $\gamma_2$  of the Brownian particle. The  $M^4$  distance  $r$  between the end points of  $\gamma_2$  would be given  $r^2 = Dt$ . The favored values of  $t$  would correspond to  $T_n = 2^n T_0$  (the full light-like geodesic). p-Adic length scales would result as  $L^2(k) = DT(k) = D2^k T_0$  for  $D = R^2/T_0$ . Since only  $CP_2$  scale is available as a fundamental scale, one would have  $T_0 = R$  and  $D = R$  and  $L^2(k) = T(k)R$ .
2. p-Adic primes near powers of 2 would be in preferred position. p-Adic time scale would not relate to the p-adic length scale via  $T_p = L_p/c$  as assumed implicitly earlier but via  $T_p = L_p^2/R_0 = \sqrt{p}L_p$ , which corresponds to secondary p-adic length scale. For instance, in the case of electron with  $p = M_{127}$  one would have  $T_{127} = .1$  second which defines a fundamental biological rhythm. Neutrinos with mass around .1 eV would correspond to  $L(169) \simeq 5 \mu\text{m}$  (size of a small cell) and  $T(169) \simeq 1. \times 10^4$  years. A deep connection between elementary particle physics and biology becomes highly suggestive.
3. In the proposed picture the p-adic prime  $p \simeq 2^k$  would characterize the thermodynamics of the random motion of light-like geodesics of  $X^3$  so that p-adic prime  $p$  would indeed be an inherent property of  $X^3$ . For  $T_p = pT_0$  the above argument is not enough for p-adic length scale hypothesis and p-adic length scale hypothesis might be seen as an outcome of a process analogous to natural selection. Resonance like effect favoring octaves of a fundamental frequency might be in question. In this case,  $p$  would a property of CD and all light-like 3-surfaces inside it and also that corresponding sector of WCW.

The above proposal involves of course ad hoc elements and can be seen only as a first attempt to understand what is involved. Later a more refined approach will be discussed.

#### 4. Mersenne primes and Gaussian Mersennes

The generalization of the embedding space required by the postulated hierarchy of Planck constants [K10] means a book like structure for which the pages are products of singular coverings or factor spaces of CD (causal diamond defined as intersection of future and past directed light-cones) and of  $CP_2$  [K10]. This predicts that Planck constants are rationals and that a given value of Planck constant corresponds to an infinite number of different pages of the Big Book, which might be seen as a drawback. If only singular covering spaces are allowed the values of Planck constant are products of integers and given value of Planck constant corresponds to a finite number of pages given by the number of decompositions of the integer to two different integers. The definition of the book like structure assigns to a given CD preferred quantization axes and so that quantum measurement has direct correlate at the level of moduli space of CDs.

TGD inspired quantum biology and number theoretical considerations suggest preferred values  $h_{eff}/h = n$ ,  $n$  as integer. Ruler and compass integers defined by the products of distinct Fermat primes and power of two are number theoretically favored values for these integers because the phases  $\exp(i2\pi/n)$  in this case are number theoretically very simple and should have emerged first in the number theoretical evolution via algebraic extensions of p-adics and of rationals. p-Adic length scale hypothesis favors powers of two as values of  $n$ .

One can however ask whether a more precise characterization of preferred Mersennes could exist and whether there could exist a stronger correlation between hierarchies of p-adic length scales and Planck constants. Mersenne primes  $M_k = 2^k - 1$ ,  $k \in \{89, 107, 127\}$ , and Gaussian Mersennes  $M_{G,k} = (1 + i)k - 1$ ,  $k \in \{113, 151, 157, 163, 167, 239, 241.. \}$  are expected to be physically highly interesting and up to  $k = 127$  indeed correspond to elementary particles. The number theoretical miracle is that all the four p-adic length scales with  $k \in \{151, 157, 163, 167\}$  are in the biologically

highly interesting range 10 nm-2.5  $\mu\text{m}$ ). The question has been whether these define scaled up copies of electro-weak and QCD type physics with ordinary value of  $h_{eff}$ . The proposal that this is the case and that these physics are in a well-defined sense induced by the dark scaled up variants of corresponding lower level physics leads to a prediction for the preferred values of  $r = 2^{k_d}$ ,  $k_d = k_i - k_j$ .

Dark variant of exotic nuclear physics implies exotic physics with ordinary value of Planck constant in the new scale in a resonant manner: dark gauge bosons transform to their ordinary variants with the same Compton length. This transformation is natural since in length scales below the Compton length the gauge bosons behave as massless and free particles. As a consequence, lighter variants of weak bosons emerge and QCD confinement scale becomes longer.

This proposal will be referred to as Mersenne hypothesis. It leads to strong predictions about EEG [K9] since it predicts a spectrum of preferred Josephson frequencies for a given value of membrane potential and also assigns to a given value of  $h_{eff}$  a fixed size scale having interpretation as the size scale of the body part or magnetic body. Also a vision about evolution of life emerges. Mersenne hypothesis is especially interesting as far as new physics in condensed matter length scales is considered: this includes exotic scaled up variants of the ordinary nuclear physics and their dark variants. Even dark nucleons are possible and this gives justification for the model of dark nucleons predicting the counterparts of DNA, RNA, tRNA, and amino-acids as well as realization of vertebrate genetic code [K33] .

These exotic nuclear physics with ordinary value of Planck constant could correspond to ground states that are almost vacuum extremals corresponding to homologically trivial geodesic sphere of  $CP_2$  near criticality to a phase transition changing Planck constant. Ordinary nuclear physics would correspond to homologically non-trivial geodesic sphere and far from vacuum extremal property. For vacuum extremals of this kind classical  $Z^0$  field proportional to electromagnetic field is present and this modifies dramatically the view about cell membrane as Josephson junction. The model for cell membrane as almost vacuum extremal indeed led to a quantitative breakthrough in TGD inspired model of EEG and is therefore something to be taken seriously. The safest option concerning empirical facts is that the copies of electro-weak and color physics with ordinary value of Planck constant are possible only for almost vacuum extremals - that is at criticality against phase transition changing Planck constant.

### 2.2.3 The origin of the preferred p-adic length scales

This question was posed already for two decades ago but remained without a convincing answer. Quite recently however the number theoretical vision allowed to understand both the origin of preferred p-adic number fields and the emergence of p-adic length scale hypothesis in a generalized form. Preferred primes are near but below powers prime which can be also larger than  $p = 2$ .

The preferred primes correspond to so called ramified rational primes, which split in to products of the primes of the extension. If some prime appears as higher than first power, one has ramification. The number of ramified primes is finite.

In strong form of holography p-adic continuations of 2-surfaces to preferred extremals identifiable as imaginations would be easy due to the existence of p-adic pseudo-constants. The continuation could fail for most configurations of partonic 2-surfaces and string world sheets in the real sector: the interpretation would be that some space-time surfaces can be imagined but not realized [K20]. For certain extensions the number of realizable imaginations could be exceptionally large. These extensions would be winners in the number theoretic fight for survival and corresponding ramified primes would be preferred p-adic primes. Whether the preferred primes satisfy p-adic length scale hypothesis or its generalization from  $p = 2$  to small primes remains an open question.

The value of effective Planck constant  $h_{eff}/h = n$  corresponds to the number of sheets of some kind of covering space defined by the space-time surface. The discretization of the space-time surface identified as a monadic manifold [L3] with embedding space preferred coordinates in extension of rationals defining the adèle has Galois group of extension as a group of symmetries permuting the sheets of the covering group. Therefore  $n = h_{eff}/h$  would naturally correspond to the dimension of the extension dividing the order of its Galois group.

Weak form of NMP allows to understand the emergence of preferred p-adic length scales. NMP favors ramified primes, for which the integer  $n$  is power of single prime  $p$ . If  $n$  is a prime slightly

below  $n_{max} = p^n$  defining the dimension of the sub-space corresponding to maximal negentropy gain, weak form of NMP favors its selection since the p-adic topology is farthest from the discrete topology assignable to formal p-adic topology characterized by  $p = 1$  [K35].

#### 2.2.4 p-Adic physics and the notion of finite measurement resolution

Canonical identification mapping p-adic numbers to reals in a continuous manner plays a key role in some applications of TGD and together with the discretization necessary to define the p-adic variants of integration and harmonic analysis suggests that p-adic topology identified as an effective topology could provide an elegant manner to characterize finite measurement resolution.

1. Finite measurement resolution can be characterized as an interval of minimum length. Below this length scale one cannot distinguish points from each other. A natural definition for this inability could be as an inability to well-order the points. The real topology is too strong in the modelling in kind of situation since it brings in large amount of processing of pseudo information whereas p-adic topology which lacks the notion of well-ordering could be more appropriate as effective topology and together with a binary cutoff could allow to get rid of the irrelevant information.
2. This suggest that canonical identification applies only inside the intervals defining finite measurement resolution in a given discretization of the space considered by say small cubes. The canonical identification is unique only modulo diffeomorphism applied on both real and p-adic side but this is not a problem since this would only reflect the absence of the well-ordering lost by finite measurement resolution. Also the fact that the map makes sense only at positive real axis would be natural if one accepts this identification.

This interpretation would suggest that there is an infinite hierarchy of measurement resolutions characterized by the value of the p-adic prime. This would mean quite interesting refinement of the notion of finite measurement resolution. At the level of quantum theory it could be interpreted as a maximization of p-adic entanglement negentropy as a function of the p-adic prime. Perhaps one might say that there is a unique p-adic effective topology allowing to maximize the information content of the theory relying on finite measurement resolution.

#### 2.2.5 p-Adic numbers and the analogy of TGD with spin-glass

The vacuum degeneracy of the Kähler action leads to a precise spin glass analogy at the level of the WCW geometry and the generalization of the energy landscape concept to TGD context leads to the hypothesis about how p-adicity could be realized at the level of WCW. Also the concept of p-adic space-time surface emerges rather naturally.

##### 1. Spin glass briefly

The basic characteristic of the spin glass phase [?] is that the direction of the magnetization varies spatially, being constant inside a given spatial region, but does not depend on time. In the real context this usually leads to large surface energies on the surfaces at which the magnetization direction changes. Regions with different direction of magnetization clearly correspond non-vacuum regions separated by almost vacuum regions. Amusingly, if 3-space is effectively p-adic and if magnetization direction is p-adic pseudo constant, no surface energies are generated so that p-adics might be useful even in the context of the ordinary spin glasses.

Spin glass phase allows a great number of different ground states minimizing the free energy. For the ordinary spin glass, the partition function is the average over a probability distribution of the coupling constants for the partition function with Hamiltonian depending on the coupling constants. Free energy as a function of the coupling constants defines “energy landscape” and the set of free energy minima can be endowed with an ultra-metric distance function using a standard construction [A38] .

##### 2. Vacuum degeneracy of Kähler action

The Kähler action defining WCW geometry allows enormous vacuum degeneracy: any four-surface for which the induced Kähler form vanishes, is an extremal of the Kähler action. Induced

Kähler form vanishes if the  $CP_2$  projection of the space-time surface is Lagrangian manifold [A9] of  $CP_2$ : these manifolds are at most two-dimensional and any canonical transformation of  $CP_2$  creates a new Lagrangian sub-manifold [A9]. An explicit representation for Lagrangian sub-manifolds is obtained using some canonical coordinates  $P_i, Q_i$  for  $CP_2$ : by assuming

$$P_i = \partial_i f(Q_1, Q_2) \quad , \quad i = 1, 2 \quad ,$$

where  $f$  arbitrary function of its arguments. One obtains a 2-dimensional sub-manifold of  $CP_2$  for which the induced Kähler form proportional to  $dP_i \wedge dQ^i$  vanishes. The roles of  $P_i$  and  $Q_i$  can obviously be interchanged. A familiar example of Lagrange manifolds are  $p_i = \text{constant}$  surfaces of the ordinary  $(p_i, q_i)$  phase space.

Since vacuum degeneracy is removed only by the classical gravitational interaction there are good reasons to expect large ground state degeneracy, when the system corresponds to a small deformation of a vacuum extremal. This degeneracy is very much analogous to the ground state degeneracy of spin glass but is 4-dimensional.

### 3. Vacuum degeneracy of the Kähler action and physical spin glass analogy

Quite generally, the dynamical reason for the physical spin glass degeneracy is the fact that Kähler action has a huge vacuum degeneracy. Any 4-surface with  $CP_2$  projection, which is a Lagrangian sub-manifold (generically two-dimensional), is vacuum extremal. This implies that space-time decomposes into non-vacuum regions characterized by non-vanishing Kähler magnetic and electric fields such that the (presumably thin) regions between the non-vacuum regions are vacuum extremals. Therefore no surface energies are generated. Also the fact that various charges and momentum and energy can flow to larger space-time sheets via wormholes is an important factor making possible strong field gradients without introducing large surfaces energies. From a given preferred extremal of Kähler action one obtains a new one by adding arbitrary space-time surfaces which is vacuum extremal and deforming them.

The symplectic invariance of the Kähler action for vacuum extremals allows a further understanding of the vacuum degeneracy. The presence of the classical gravitational interaction spoils the canonical group  $Can(CP_2)$  as gauge symmetries of the action and transforms it to the isometry group of  $CH$ . As a consequence, the  $U(1)$  gauge degeneracy is transformed to a spin glass type degeneracy and several, perhaps even infinite number of maxima of Kähler function become possible. Given sheet has naturally as its boundary the 3-surfaces for which two maxima of the Kähler function coalesce or are created from single maximum by a cusp catastrophe [A30]. In catastrophe regions there are several sheets and the value of the maximum Kähler function determines which give a measure for the importance of various sheets. The quantum jumps selecting one of these sheets can be regarded as phase transitions.

In TGD framework classical non-determinism forces to generalize the notion of the 3-surface by replacing it with a sequence of space like 3-surfaces having time like separations such that the sequence characterizes uniquely one branch of multi-furcation. This characterization works when non-determinism has discrete nature. For  $CP_2$  type extremals which are bosonic vacua, basic objects are essentially four-dimensional since  $M_+^4$  projection of  $CP_2$  type extremal is random light like curve. This effective four-dimensionality of the basic objects makes it possible to topologize Feynman diagrammatics of quantum field theories by replacing the lines of Feynman diagrams with  $CP_2$  type extremals.

In TGD framework spin glass analogy holds true also in the time direction, which reflects the fact that the vacuum extremals are non-deterministic. For instance, by gluing vacuum extremals with a finite space-time extension (also in time direction!) to a non-vacuum extremal and deforming slightly, one obtains good candidates for the degenerate preferred extremals. This non-determinism is expected to make the preferred extremals of the Kähler action highly degenerate. The construction of S-matrix at the high energy limit suggests that since a localization selecting one degenerate maximum occurs, one must accept as a fact that each choice of the parameters corresponds to a particular S-matrix and one must average over these choices to get scattering rates. This averaging for scattering rates corresponds to the averaging over the thermodynamical partition functions for spin glass. A more general is that one allows final state wave functions to depend on the zero modes which affect S-matrix elements: in the limit that wave functions are completely localized, one ends up with the simpler scenario.



#### 4. *p*-Adic non-determinism and spin glass analogy

One must carefully distinguish between cognitive and physical spin-glass analogy. Cognitive spin-glass analogy is due to the *p*-adic non-determinism. *p*-Adic pseudo constants induce a non-determinism which essentially means that *p*-adic extrema depend on the *p*-adic pseudo constants which depend on a finite number of positive binary digits of their arguments only. Thus *p*-adic extremals are glued from pieces for which the values of the integration constants are genuine constants. Obviously, an optimal cognitive representation is achieved if pseudo constants reduce to ordinary constants.

More precisely, any function

$$\begin{aligned} f(x) &= f(x_N) , \\ x_N &= \sum_{k \leq N} x_k p^k , \end{aligned} \quad (2.1)$$

which does not depend on the binary digits  $x_n$ ,  $n > N$  has a vanishing *p*-adic derivative and is thus a pseudo constant. These functions are piecewise constant below some length scale, which in principle can be arbitrary small but finite. The result means that the constants appearing in the solutions the *p*-adic field equations are constants functions only below some length scale. For instance, for linear differential equations integration constants are arbitrary pseudo constants. In particular, the *p*-adic counterparts of the preferred extremals are highly degenerate because of the presence of the pseudo constants. This in turn means a characteristic randomness of the spin glass also in the time direction since the surfaces at which the pseudo constants change their values do not give rise to infinite surface energy densities as they would do in the real context.

The basic character of cognition would be spin glass like nature making possible “engineering” at the level of thoughts (planning) whereas classical non-determinism of the Kähler action would make possible “engineering” at the level of the real world.

#### **2.2.6 Life as islands of rational/algebraic numbers in the seas of real and *p*-adic continua?**

The possibility to define entropy differently for rational/algebraic entanglement and the fact that number theoretic entanglement entropy can be negative raises the question about which kind of systems can possess this kind of entanglement. I have considered several identifications but the most elegant interpretation is based on the idea that living matter resides in the intersection of real and *p*-adic worlds, somewhat like rational numbers live in the intersection of real and *p*-adic number fields. This intersection would be number theoretically universal in the sense that algebraic extension of rationals would be the number field but in rather abstract sense: for the parameters defining the WCW coordinates characterizing space-time surface rather than points of space-time surface.

The observation that Shannon entropy allows an infinite number of number theoretic variants for which the entropy can be negative in the case that probabilities are algebraic numbers leads to the idea that living matter in a well-defined sense corresponds to the intersection of real and *p*-adic worlds. This would mean that the mathematical expressions for the space-time surfaces (or at least 3-surfaces or partonic 2-surfaces and their 4-D tangent planes) make sense in both real and *p*-adic sense for some primes *p*. Same would apply to the expressions defining quantum states. In particular, entanglement probabilities would be rationals or algebraic numbers so that entanglement can be negentropic and the formation of bound states in the intersection of real and *p*-adic worlds generates information and is thus favored by NMP.

This picture has also a direct connection with consciousness.

1. The generation of non-rational (non-algebraic) bound state entanglement between the system and external world means that the system loses consciousness during the state function reduction process following the *U*-process generating the entanglement. What happens that the Universe corresponding to given CD decomposes to two un-entangled subsystems, which in turn decompose, and the process continues until all subsystems have only entropic bound state entanglement or negentropic algebraic entanglement with the external world.

2. If the sub-system generates entropic bound state entanglement in the process, it loses consciousness. Note that the entanglement entropy of the sub-system is a sum over entanglement entropies over all subsystems involved. This hierarchy of subsystems corresponds to the hierarchy of sub-CDs so that the survival without a loss of consciousness depends on what happens at all levels below the highest level for a given self. In more concrete terms, ability to stay conscious depends on what happens at cellular level too. The stable evolution of systems having algebraic entanglement is expected to be a process proceeding from short to long length scales as the evolution of life indeed is.
3.  $U$ -process generates a superposition of states in which any sub-system can have both real and algebraic entanglement with the external world. This would suggest that the choice of the type of entanglement is a volitional selection. A possible interpretation is as a choice between good and evil. The hedonistic complete freedom resulting as the entanglement entropy is reduced to zero on one hand, and the algebraic bound state entanglement implying correlations with the external world and meaning giving up the maximal freedom on the other hand. The hedonistic option is risky since it can lead to non-algebraic bound state entanglement implying a loss of consciousness. The second option means expansion of consciousness - a fusion to the ocean of consciousness as described by spiritual practices.
4. This formulation means a sharpening of the earlier statement “Everything is conscious and consciousness can be only lost” with the additional statement “This happens when non-algebraic bound state entanglement is generated and the system does not remain in the intersection of real and p-adic worlds anymore”. Clearly, the quantum criticality of TGD Universe seems has very many aspects and life as a critical phenomenon in the number theoretical sense is only one of them besides the criticality of the space-time dynamics and the criticality with respect to phase transitions changing the value of Planck constant and other more familiar criticalities. How closely these criticalities relate remains an open question.

A good guess is that algebraic entanglement is essential for quantum computation, which therefore might correspond to a conscious process. Hence cognition could be seen as a quantum computation like process, a more appropriate term being quantum problem solving. Living-dead dichotomy could correspond to rational-irrational or to algebraic-transcendental dichotomy: this at least when life is interpreted as intelligent life. Life would in a well defined sense correspond to islands of rationality/algebraicity in the seas of real and p-adic continua.

The view about the crucial role of rational and algebraic numbers as far as intelligent life is considered, could have been guessed on very general grounds from the analogy with the orbits of a dynamical system. Rational numbers allow a predictable periodic decimal/pinary expansion and are analogous to one-dimensional periodic orbits. Algebraic numbers are related to rationals by a finite number of algebraic operations and are intermediate between periodic and chaotic orbits allowing an interpretation as an element in an algebraic extension of any p-adic number field. The projections of the orbit to various coordinate directions of the algebraic extension represent now periodic orbits. The decimal/pinary expansions of transcendentals are un-predictable being analogous to chaotic orbits. The special role of rational and algebraic numbers was realized already by Pythagoras, and the fact that the ratios for the frequencies of the musical scale are rationals supports the special nature of rational and algebraic numbers. The special nature of the Golden Mean, which involves  $\sqrt{5}$ , conforms the view that algebraic numbers rather than only rationals are essential for life.

### 2.2.7 p-Adic physics as physics of cognition

The vision about p-adic physics as physics of cognition has gradually established itself as one of the key idea of TGD inspired theory of consciousness. There are several motivations for this idea.

The strongest motivation is the vision about living matter as something residing in the intersection of real and p-adic worlds. One of the earliest motivations was p-adic non-determinism identified tentatively as a space-time correlate for the non-determinism of imagination. p-Adic non-determinism follows from the fact that functions with vanishing derivatives are piecewise constant functions in the p-adic context. More precisely, p-adic pseudo constants depend on the pinary cutoff of their arguments and replace integration constants in p-adic differential equations. In the

case of field equations this means roughly that the initial data are replaced with initial data given for a discrete set of time values chosen in such a way that unique solution of field equations results. Solution can be fixed also in a discrete subset of rational points of the embedding space. Presumably the uniqueness requirement implies some unique binary cutoff. Thus the space-time surfaces representing solutions of p-adic field equations are analogous to space-time surfaces consisting of pieces of solutions of the real field equations. p-Adic reality is much like the dream reality consisting of rational fragments glued together in illogical manner or pieces of child's drawing of body containing body parts in more or less chaotic order.

The obvious looking interpretation for the solutions of the p-adic field equations is as a geometric correlate of imagination. Plans, intentions, expectations, dreams, and cognition in general are expected to have p-adic space-time sheets as their geometric correlates. This in the sense that p-adic space-time sheets somehow initiate the real neural processes providing symbolic counterparts for the cognitive representations provided by p-adic space-time sheets and p-adic fermions. A deep principle seems to be involved: incompleteness is characteristic feature of p-adic physics but the flexibility made possible by this incompleteness is absolutely essential for imagination and cognitive consciousness in general.

Although p-adic space-time sheets as such are not conscious, p-adic physics would provide beautiful mathematical realization for the intuitions of Descartes. The formidable challenge is to develop experimental tests for p-adic physics. The basic problem is that we can perceive p-adic reality only as "thoughts" unlike the "real" reality which represents itself to us as sensory experiences. Thus it would seem that we should be able generalize the physics of sensory experiences to physics of cognitive experiences.

## 2.3 What Is The Correspondence Between P-Adic And Real Numbers?

There must be some kind of correspondence between reals and p-adic numbers. This correspondence can depend on context. In p-adic mass calculations one must map p-adic mass squared values to real numbers in a continuous manner and canonical identification  $x = \sum x_n p^n \rightarrow Id(x) = \sum x_n p^{-n}$  is a natural first guess. Also p-adic probabilities could be mapped to their real counterparts by a suitable normalization. The minimalistic interpretation is that real and p-adic mass calculations must give same results- physics must be consistent with the existence of cognitive representations of it. In this case p-adic thermodynamics would constrain the temperature and scale parameters of real thermodynamics.

The possible existence and the nature of the correspondence at the level of embedding space and space-time surfaces is much more questionable and it is far from clear whether it is needed as a naïve map of real space-time points to p-adic space-time points by - say - canonical identification: the problem would be that symmetries are not respected if one demands continuity. One would like to various symmetries in real and p-adic variants and the correspondence should respect symmetries.

One can wonder whether p-adic valued S-matrices have any physical meaning and whether they could be obtained as algebraic continuation from a number theoretically universal S-matrix whose matrix elements are algebraic numbers allowing an interpretation as real or p-adic numbers in suitable algebraic extension: this would pose extremely strong constraints on S-matrix. If one wants to introduce p-adic physics at space-time level one must be able to relate p-adic and real space-time regions to each other. The identification along common rational points of real and various p-adic variants of the embedding space produces however problems with symmetries.

In the following these questions are discussed as I did them before the recent steps of progress summarized in the last subsection. I hope that the reader can forgive certain naïvete of the discussion: pioneering work is in question.

### 2.3.1 Generalization of the number concept

The recent view about the unification of real and p-adic physics is based on the generalization of number concept obtained by fusing together real and p-adic number fields along common rationals (see **Fig. ??** in the Appendix.

#### 1. Rational numbers as numbers common to all number fields

The unification of real physics of material work and p-adic physics of cognition leads to the

generalization of the notion of number field. Reals and various p-adic number fields are glued along common algebraic numbers defining an extension of p-adic numbers to form a fractal book like structure. Allowing all possible finite-dimensional algebraic and perhaps even transcendental extensions of rationals inducing those of p-adic numbers adds additional pages to this “Big Book”.

This suggests a generalization of the notion of manifold as real manifold and its p-adic variants glued together along common points. This generalization might make sense under very high symmetries and that it is safest to lean strongly on the physical picture provided by quantum TGD. This construction is discussed in [K38] and one must make clear that it is plagued difficulties with symmetries.

1. The most natural guess is that the coordinates of common points are rational or in some algebraic extension of rational numbers. General coordinate invariance and preservation of symmetries require preferred coordinates existing when the manifold has maximal number of isometries. This approach might make sense in the case of linear spaces- in particular Minkowski space  $M^4$ . The natural coordinates are in this case linear Minkowski coordinates. The choice of coordinates is however not completely unique and has interpretation as a geometric correlate for the choice of quantization axes for a given CD. Different choices are not equivalent.
2. As will be found, the need to have a well-defined integration based on Fourier analysis (or its generalization to harmonic analysis [A5] in symmetric spaces) poses very strong constraints and allows p-adicization only if the space has maximal symmetries. Fourier analysis requires the introduction of an algebraic extension of p-adic numbers containing sufficiently many roots of unity.
  - (a) This approach is especially natural in the case of compact symmetric spaces such as  $CP_2$  [A3].
  - (b) Also symmetric spaces such the 3-D proper time  $a = \text{constant}$  hyperboloid of  $M^4$ - call it  $H(a)$  -allowing Lorentz group as isometries allows a p-adic variant utilizing the hyperbolic counterparts for the roots of unity.  $M^4 \times H(a = 2^n a_0)$  appears as a part of the moduli space of CDs.
  - (c) For light-cone boundaries associated with CDs  $SO(3)$  invariant radial coordinate  $r_M$  defining the radius of sphere  $S^2$  defines the hyperbolic coordinate and angle coordinates of  $S^2$  would correspond to phase angles and  $M^4_{\pm}$  projections for the common points of real and p-adic variants of partonic 2-surfaces would be this kind of points. Same applies to  $CP_2$  projections.

In the “intersection of real and p-adic worlds” real and p-adic partonic 2-surfaces would obey same algebraic equations and would be obtained by an algebraic continuation from the corresponding equations making sense in the discrete variant of  $M^4_{\pm} \times CP_2$ . This connection with discrete sub-manifold geometries means very powerful constraints on the partonic 2-surfaces in the intersection.

3. The common algebraic points of real and p-adic variant of the manifold form a discrete space but one could identify the p-adic counterpart of the real discretization intervals  $(0, 2\pi/N)$  for angle like variables as p-adic numbers of norm smaller than 1 using canonical identification or some variant of it. Same applies to the the hyperbolic counterpart of this interval. The non-uniqueness of this map could be interpreted in terms of a finite measurement resolution. In particular, the condition that WCW allows Kähler geometry requires a decomposition to a union of symmetric spaces so that there are good hopes that p-adic counterpart is analogous to that assigned to  $CP_2$ .

This approach works for probabilities but has serious problems with symmetries. The only manner to circumvent the problems is based on strong form of holography and abstraction of the real-p-adic correspondence so that it is not anymore local but maps entire surfaces to each other. One must have also now discretization and co-dimension two rule holds true. For instance, space-time surfaces are replaced with a collection of 2-D objects and partonic 2-surfaces by a discrete set of points. This rule is equivalent with strong form of holography.

The correspondence would be at the level of parameters defining WCW coordinates and intersection of reality and p-adicities would consist of discrete set of 2-surfaces. As already explained, strong form of holography suggests that real and p-adic space-time sheets are obtained by continuation of the 2-surfaces to preferred extremals by assuming that the classical Noether charges associated with super-symplectic algebra vanish for the 3-surfaces at the ends of space-time surface. By conformal invariance the parameters would be naturally general coordinate invariant conformal moduli for the 2-surfaces involved, and belong to the algebraic extension of rationals in the intersection. Their continuation to various number fields would give real and p-adic space-time sheets. Also scattering amplitudes could be constructed using the data assigned with 2-surfaces in the intersection and continued algebraically to various number fields. This picture conforms also with the recipe for constructing scattering amplitudes in twistor approach [L1].

### *2. How large p-adic space-time sheets can be?*

Space-time region having finite size in the real sense can have arbitrarily large size in p-adic sense and vice versa. This raises a rather thought provoking questions. Could the p-adic space-time sheets have cosmological or even infinite size with respect to the real metric but have be p-adically finite? How large space-time surface is responsible for the p-adic representation of my body? Could the large or even infinite size of the cognitive space-time sheets explain why creatures of a finite physical size can invent the notion of infinity and construct cosmological theories? Could it be that binary cutoff  $O(p^n)$  defining the resolution of a p-adic cognitive representation would define the size of the space-time region needed to realize the cognitive representation?

These questions make sense if the real-padic correspondence is local - that is defined by the intersection real and p-adic space-time surfaces. In the more abstract approach it does not make sense.

In fact, the mere requirement that the neighborhood of a point of the p-adic space-time sheet contains points, which are p-adically infinitesimally near to it can mean that points infinitely distant from this point in the real sense are involved. A good example is provided by an integer valued point  $x = n < p$  and the point  $y = x + p^m$ ,  $m > 0$ : the p-adic distance of these points is  $p^{-m}$  whereas at the limit  $m \rightarrow \infty$  the real distance goes as  $p^m$  and becomes infinite for infinitesimally near points. The points  $n + y$ ,  $y = \sum_{k>0} x_k p^k$ ,  $0 < n < p$ , form a p-adically continuous set around  $x = n$ . In the real topology this point set is discrete set with a minimum distance  $\Delta x = p$  between neighboring points whereas in the p-adic topology every point has arbitrary nearby points. There are also rationals, which are arbitrarily near to each other both p-adically and in the real sense. Consider points  $x = m/n$ ,  $m$  and  $n$  not divisible by  $p$ , and  $y = (m/n) \times (1 + p^k r)/(1 + p^k s)$ ,  $s = r + 1$  such that neither  $r$  or  $s$  is divisible by  $p$  and  $k \gg 1$  and  $r \gg p$ . The p-adic and real distances are  $|x - y|_p = p^{-k}$  and  $|x - y| \simeq (m/n)/(r + 1)$  respectively. By choosing  $k$  and  $r$  large enough the points can be made arbitrarily close to each other both in the real and p-adic senses.

The idea about astrophysical size of the p-adic cognitive space-time sheets providing representation of body and brain is consistent with TGD inspired theory of consciousness, which forces to take very seriously the idea that even human consciousness involves astrophysical length scales. It must be however emphasized that this kind of concretization seems to be un-necessary if the correspondence is at the level of WCW.

### *3. Generalization of complex analysis*

One general idea which results as an outcome of the generalized notion of number is the idea of a universal function continuable from a function mapping rationals to rationals or to a finite extension of rationals to a function in any number field. This algebraic continuation is analogous to the analytical continuation of a real analytic function to the complex plane. Rational functions for which polynomials have rational coefficients are obviously functions satisfying this constraint. Algebraic functions for which polynomials have rational coefficients satisfy this requirement if appropriate finite-dimensional algebraic extensions of p-adic numbers are allowed.

For instance, one can ask whether residue calculus might be generalized so that the value of an integral along the real axis could be calculated by continuing it instead of the complex plane to any number field via its values in the subset of rational numbers forming the back of the book like structure (in very metaphorical sense) having number fields as its pages. If the poles of the continued function in the finitely extended number field allow interpretation as real numbers it might be possible to generalize the residue formula. One can also imagine of extending residue

calculus to any algebraic extension. An interesting situation arises when the poles correspond to extended p-adic rationals common to different pages of the “Big Book”. Could this mean that the integral could be calculated at any page having the pole common. In particular, could a p-adic residue integral be calculated in the ordinary complex plane by utilizing the fact that in this case numerical approach makes sense. Contrary to the first expectations the algebraically continued residue calculus does not seem to have obvious applications in quantum TGD.

### 2.3.2 Canonical identification

Canonical There exists a natural continuous map  $Id : R_p \rightarrow R_+$  from p-adic numbers to non-negative real numbers given by the “pinary” expansion of the real number for  $x \in R$  and  $y \in R_p$  this correspondence reads

$$\begin{aligned} y &= \sum_{k>N} y_k p^k \rightarrow x = \sum_{k<N} y_k p^{-k} , \\ y_k &\in \{0, 1, \dots, p-1\} . \end{aligned} \quad (2.2)$$

This map is continuous as one easily finds out. There is however a little difficulty associated with the definition of the inverse map since the pinary expansion like also desimal expansion is not unique ( $1 = 0.999\dots$ ) for the real numbers  $x$ , which allow pinary expansion with finite number of pinary digits

$$\begin{aligned} x &= \sum_{k=N_0}^N x_k p^{-k} , \\ x &= \sum_{k=N_0}^{N-1} x_k p^{-k} + (x_N - 1)p^{-N} + (p-1)p^{-N-1} \sum_{k=0,\dots} p^{-k} . \end{aligned} \quad (2.3)$$

The p-adic images associated with these expansions are different

$$\begin{aligned} y_1 &= \sum_{k=N_0}^N x_k p^k , \\ y_2 &= \sum_{k=N_0}^{N-1} x_k p^k + (x_N - 1)p^N + (p-1)p^{N+1} \sum_{k=0,\dots} p^k \\ &= y_1 + (x_N - 1)p^N - p^{N+1} , \end{aligned} \quad (2.4)$$

so that the inverse map is either two-valued for p-adic numbers having expansion with finite number of pinary digits or single valued and discontinuous and non-surjective if one makes pinary expansion unique by choosing the one with finite number of pinary digits. The finite number of pinary digits expansion is a natural choice since in the numerical work one always must use a pinary cutoff on the real axis.

#### 1. Canonical identification is a continuous map of non-negative reals to p-adics

The topology induced by the inverse of the canonical identification map in the set of positive real numbers differs from the ordinary topology. The difference is easily understood by interpreting the p-adic norm as a norm in the set of the real numbers. The norm is constant in each interval  $[p^k, p^{k+1})$  (see **Fig. ??**) and is equal to the usual real norm at the points  $x = p^k$ : the usual linear norm is replaced with a piecewise constant norm. This means that p-adic topology is coarser than the usual real topology and the higher the value of  $p$  is, the coarser the resulting topology is above a given length scale. This hierarchical ordering of the p-adic topologies will be a central feature as far as the proposed applications of the p-adic numbers are considered.

Ordinary continuity implies p-adic continuity since the norm induced from the p-adic topology is rougher than the ordinary norm. This allows two alternative interpretations. Either p-adic image of a physical systems provides a good representation of the system above some binary cutoff or the physical system can be genuinely p-adic below certain length scale  $L_p$  and become in good approximation real, when a length scale resolution  $L_p$  is used in its description. The first interpretation is correct if canonical identification is interpreted as a cognitive map. p-Adic continuity implies ordinary continuity from right as is clear already from the properties of the p-adic norm (the graph of the norm is indeed continuous from right, see **Fig. ??** of Appendix). This feature is one clear signature of the p-adic topology.

If one considers seriously the application of canonical identification to basic quantum TGD one cannot avoid the question about the p-adic counterparts of the negative real numbers. There is no satisfactory manner to circumvent the fact that canonical images of p-adic numbers are naturally non-negative. This is not a problem if canonical identification applies only to the coordinate interval  $(0, 2\pi/N)$  or its hyperbolic variant defining the finite measurement resolution. That p-adicization program works only for highly symmetric spaces is not a problem from the point of view of TGD.

### 2. Canonical identification relates p-adic and real statistical physics

p-Adic mass calculations based on p-adic thermodynamics were the first and rather successful application of the p-adic physics (see the four chapters in [K18] . The essential element of the approach was the replacement of the Boltzmann weight  $e^{-E/T}$  with its p-adic generalization  $p^{L_0/T_p}$ , where  $L_0$  is the Virasoro generator corresponding to scaling and representing essentially mass squared operator instead of energy.  $T_p$  is inverse integer valued p-adic temperature. The predicted mass squared averages were mapped to real numbers by canonical identification.

One could also construct a real variant of this approach by considering instead of the ordinary Boltzmann weights the weights  $p^{-L_0/T_p}$ . The quantization of temperature to  $T_p = \log(p)/n$  would be a completely ad hoc assumption. In the case of real thermodynamics all particles are predicted to be light whereas in case of p-adic thermodynamics particle is light only if the ratio for the degeneracy of the lowest massive state to the degeneracy of the ground state is integer. Immense number of particles disappear from the spectrum of light particles by this criterion. For light particles the predictions are same as of p-adic thermodynamics in the lowest non-trivial order but in the next order deviations are possible.

Also p-adic probabilities and the p-adic entropy can be mapped to real numbers by canonical identification. The general idea is that a faithful enough cognitive representation of the real physics can by the number theoretical constraints involved make predictions, which would be extremely difficult to deduce from real physics.

### 3. Variant of canonical identification commuting with division of integers

The basic problems of canonical identification is that it does not respect unitarity. For this reason it is not well suited for relating p-adic and real scattering amplitudes. The problem of the correspondence via direct rationals or roots of unity is that it does not respect continuity. The restriction of S-matrix to a discrete intersection of real and p-adic worlds is one manner to solve this difficulty.

One can also consider alternative approach to achieve a compromise between algebra and topology achieved by using a modification of canonical identification  $I_{R_p \rightarrow R}$  defined as  $I_1(r/s) = I(r)/I(s)$ . If the conditions  $r \ll p$  and  $s \ll p$  hold true, the map respects algebraic operations and also unitarity and various symmetries. It seems that this option must be used to relate p-adic transition amplitudes to real ones and vice versa [K16] . In particular, real and p-adic coupling constants are related by this map. Also some problems related to p-adic mass calculations find a nice resolution when  $I_1$  is used.

This variant of canonical identification is not equivalent with the original one using the infinite expansion of  $q$  in powers of  $p$  since canonical identification does not commute with product and division. The variant is however unique in the recent context when  $r$  and  $s$  in  $q = r/s$  have no common factors. For integers  $n < p$  it reduces to direct correspondence.

Generalized numbers would be regarded in this picture as a generalized manifold obtained by gluing different number fields together along rationals. Instead of a direct identification of real and p-adic rationals, the p-adic rationals in  $R_p$  are mapped to real rationals (or vice versa) using a variant of the canonical identification  $I_{R \rightarrow R_p}$  in which the expansion of rational number

$q = r/s = \sum r_n p^n / \sum s_n p^n$  is replaced with the rational number  $q_1 = r_1/s_1 = \sum r_n p^{-n} / \sum s_n p^{-n}$  interpreted as a p-adic number:

$$q = \frac{r}{s} = \frac{\sum_n r_n p^n}{\sum_m s_m p^m} \rightarrow q_1 = \frac{\sum_n r_n p^{-n}}{\sum_m s_m p^{-m}} . \quad (2.5)$$

$R_{p_1}$  and  $R_{p_2}$  are glued together along common rationals by an the composite map  $I_{R \rightarrow R_{p_2}} I_{R_{p_1} \rightarrow R}$ .

This variant of canonical identification seems to be an excellent candidate for mapping the predictions of p-adic mass calculations to real numbers and also for relating p-adic and real scattering amplitudes to each other [K16] . The deviations of predictions from those for standard form of canonical identification are however small.

The cautious conclusion of this section is that symmetric space approach involving both the identification along common rationals of roots of unity in large and canonical identification below the measurement resolution provide the safest approach to the p-adicization of quantum TGD. The impossibility to well-order the points below measurement resolution explains why effective p-adic topology works in real context. The discussion of integration and Fourier analysis will provide further support for the conclusion.

## 2.4 P-Adic Variants Of The Basic Mathematical Structures Relevant To Physics

The basic existential questions worrying a person planning to become a p-adic quantum physicist are rather obvious. How to define p-adic probabilities, p-adic thermodynamics, and p-adic unitarity and perhaps even p-adic Hilbert space? Is it possible to define the p-adic variant of the manifold concept? As already noticed for symmetric spaces p-adic variants might exist but what about space-time surfaces: could it be enough to consider only the p-adic variants of the partonic 2-surfaces in the manner already discussed? Can one somehow circumvent the difficulties related to the definition of the p-adic variant of definite integral? Perhaps by using Fourier analysis? How can one circumvent the fact that the basic variational principle involves integral over space-time surface which is p-adically notoriously difficult to define? Is all this just a waste of time or could it be that the enormous constraints from p-adicization could provide information about real physics not achievable otherwise (as in the case of p-adic mass calculations)?

### 2.4.1 p-Adic probabilities

p-Adic super conformal representations necessitate p-adic QM based on the p-adic unitarity and p-adic probability concepts. The concept of a p-adic probability indeed makes sense as shown by [A21] . p-Adic probabilities can be defined as relative frequencies  $N_i/N$  in a long series consisting of total number  $N$  of observations and  $N_i$  outcomes of type  $i$ . Probability conservation corresponds to

$$\sum_i N_i = N , \quad (2.6)$$

and the only difference as compared to the usual probability is that the frequencies are interpreted as p-adic numbers.

The interpretation as p-adic numbers means that the relative frequencies converge to probabilities in a p-adic rather than real sense in the limit of a large number  $N$  of observations. If one requires that probabilities are limiting values of the frequency ratios in p-adic sense one must pose restrictions on the possible numbers of the observations  $N$  if  $N$  is larger than  $p$ . For  $N$  smaller than  $p$ , the situation is similar to the real case. This means that for  $p = M_{127} \simeq 10^{38}$ , appropriate for the particle physics experiments, p-adic probability differs in no observable manner from the ordinary probability.

If the number of observations is larger than  $p$ , the situation changes. If  $N_1$  and  $N_2$  are two numbers of observations they are near to each other in the p-adic sense if they differ by a large power of  $p$ . A possible interpretation of this restriction is that the observer at the  $p$ :th level of



the condensate cannot choose the number of the observations freely. The restrictions to this freedom come from the requirement that the sensible statistical questions in a p-adically conformally invariant world must respect p-adic conformal invariance [A14].

The most important application of the p-adic probability is the description of the particle massivation based on p-adic thermodynamics. Instead of energy, Virasoro generator  $l$  is thermalized and in the low temperature phase temperature is quantized in the sense that the counterpart of the Boltzmann weight  $\exp(H/T)$  is  $p^{L_0/T}$ , where  $T = 1/n$  from the requirement that Boltzmann weight exists ( $L_0$  has integer spectrum). The surprising success of the mass calculations shows that p-adic probability theory is much more than a formal possibility.

In particle physics context coupling constant evolution is replaced with a discrete p-adic coupling constant evolution and the renormalization is related to the change of the reduction of the p-adic length scale  $L_p$  in the length scale hierarchy rather than p-adic fractality for a fixed value of  $p$ . In ZEO the evolution corresponds to the hierarchy of CDs with scales coming as powers of 2 in accordance with p-adic length scale hypothesis.

1. p-Adic probabilities and p-adic fractals

p-Adic probabilities are natural in the statistical description of the fractal structures, which can contain same structural detail with all possible sizes.

1. The concept of a structural detail in a fractal seems to be reasonably well defined concept. The structural detail is clearly fixed by its topology and p-adic conformal invariants associated with it. Clearly, a finite resolution defined by some power of  $p$  of the p-adic cutoff scale must be present in the definition. For example, p-adic angles are conformal invariants in the p-adic case, too. The overall size of the detail doesn't matter. Let us therefore assume that it is possible to make a list, possibly infinite, of the structural details appearing in the p-adic fractal.
2. What kind of questions related to the structural details of the p-adic fractal one can ask? The first thing one can ask is how many times  $i$ :th structural detail appears in a finite region of the fractal structure: although this number is infinite as a real number it might possess (and probably does so!) finite norm as a p-adic number and provides a useful p-adic invariant of the fractal. If a complete list about the structural details of the fractal is at use one can calculate also the total number of structural details defined as  $N = \sum_i N_i$ . This means that one can also define p-adic probability for the appearance of  $i$ :th structural detail as a relative frequency  $p_i = N_i/N$ .
3. One can consider conditional probabilities, too. It is natural to ask what is the probability for the occurrence of the structural detail subject to the condition that part of the structural detail is fixed (apart from the p-adic conformal transformations). In order to evaluate these probabilities as relative frequencies one needs to look only for those structural details containing the substructure in question.
4. The evaluation of the p-adic probabilities of occurrence can be done by evaluating the required numbers  $N_i$  and  $N$  in a given resolution. A better estimate is obtained by increasing the resolution and counting the numbers of the hitherto unobserved structural details. The increase in the resolution greatly increases the number of the observations in case of p-adic fractal and the fluctuations in the values of  $N_i$  and  $N$  increase with the resolution so that  $N_i/N$  has no well defined limit as a real number although one can define the probabilities of occurrence as a resolution dependent concept. In the p-adic sense the increase in the values of  $N_i$  and fluctuations are small and the procedure should converge rapidly so that reliable estimates should result with quite a reasonable resolution. Notice that the increase of the fluctuations in the real sense, when resolution is increased is in accordance with the criticality of the system.
5. p-Adic frequencies and probabilities define via the canonical correspondence real valued invariants of the fractal structure.

p-Adic fractality in this sense could have practical applications only for small values of  $p$ . They could be important in the macroscopic length scales if the hierarchy of Planck constants meaning

scaling up  $L_p \rightarrow \sqrt{r}L_p$ ,  $r = h_{eff}/h$ , of the p-adic length scales. In elementary particle physics  $L_p$  is of the order of the Compton length associated with the particle for  $r = 1$  and already in the first downward step  $CP_2$  length scale  $R$  is achieved whereas upward step gives astrophysical length scale in the case of electron ( $p = M_{127} = 2^{127} - 1$ ) for instance. For large enough values of Planck constant and for small p-adic primes  $p$  the situation changes.

2. Relationship between p-adic and real probabilities

There are uniqueness problems related to the mapping of p-adic probabilities to real ones. These problems find a nice resolution from the requirement that the map respects probability conservation. The implied modification of the original mapping does not change measurably the predictions for the masses of light particles.

a) *How unique the map of p-adic probabilities and mass squared values are mapped to real numbers is?*

The mapping of p-adic thermodynamical probabilities and mass squared values to real numbers is not completely unique.

1. The canonical identification  $Id : \sum x_n p^n \rightarrow \sum x_n p^{-n}$  takes care of this mapping but does not respect the sum of probabilities so that the real images  $I(p_n)$  of the probabilities must be normalized. This is a somewhat alarming feature.
2. The modification of the canonical identification mapping rationals by the formula  $I(r/s) = I(r)/I(s)$  has appeared naturally in various applications, in particular because it respects unitarity of unitary matrices with rational elements with  $r < p, s < p$ . In the case of p-adic thermodynamic the formula  $I(g(n)p^n/Z) \rightarrow I(g(n)p^n)/I(Z)$  would be very natural although  $Z$  need not be rational anymore. For  $g(n) < p$  the real counterparts of the p-adic probabilities would sum up to one automatically for this option. One cannot deny that this option is more convincing than the original one. The generalization of this formula to map p-adic mass squared to a real one is obvious.
3. Options 1) and 2) differ dramatically when the  $n = 0$  massless ground state has ground state degeneracy  $D > 1$ . For option 1) the real mass is predicted to be of order  $CP_2$  mass whereas for option 2) it would be by a factor  $1/D$  smaller than the minimum mass predicted by the option 1). Thus option 2) would predict a large number of additional exotic states. For those states which are light for option 1), the two options make identical predictions as far as the significant two lowest order terms are considered. Hence this interpretation would not change the predictions of the p-adic mass calculations in this respect. Option 2) is definitely more in accord with the real physics based intuitions and the main role of p-adic thermodynamics would be to guarantee the quantization of the temperature and fix practically uniquely the spectrum of the "Hamiltonian".

b) *Under what conditions the mapping of p-adic ensemble probabilities to real probabilities respects probability conservation?*

One can consider also a more general situation. Assume that one has an ensemble consisting of independent elementary events such that the number of events of type  $i$  is  $N_i$ . The probabilities are given by  $p_i = N_i/N$  and  $N = \sum N_i$  is the total number of elementary events. Even in the case that  $N$  is infinite as a real number it is natural to map the p-adic probabilities to their real counterparts using the rational canonical identification  $I(p_i) = I(N_i)/I(N)$ . Of course,  $N_i$  and  $N$  exist as well defined p-adic numbers under very stringent conditions only.

The question is under what conditions this map respects probability conservation. The answer becomes obvious by looking at the binary expansions of  $N_i$  and  $N$ . If the integers  $N_i$  (possibly infinite as real integers) have binary expansions having no common binary digits, the sum of probabilities is conserved in the map. Note that this condition can assign also to a finite ensemble with finite number of a unique value of  $p$ .

This means that the selection of a basis for independent events corresponds to a decomposition of the set of integers labelling binary digits to disjoint sets and brings in mind the selection of

orthonormalized basis of quantum states in quantum theory. What is physically highly non-trivial that this “orthogonalization” alone puts strong constraints on probabilities of the allowed elementary events. One can say that the probabilities define distributions of binary digits analogous to non-negative probability amplitudes in the space of integers labelling binary digits, and the probabilities of independent events must be orthogonal with respect to the inner product defined by point-wise multiplication in the space of binary digits.

p-Adic thermodynamics for which Boltzmann weights  $g(E)exp(-E/T)$  are replaced by  $g(E)p^{E/T}$  such that one has  $g(E) < p$  and  $E/T$  is integer valued, satisfies this constraint. The quantization of  $E/T$  to integer values implies quantization of both  $T$  and “energy” spectrum and forces so called super conformal invariance [A14, A15] in TGD applications, which is indeed a basic symmetry of the theory.

There are infinitely many ways to choose the elementary events and each choice corresponds to a decomposition of the infinite set of integers  $n$  labelling the powers of  $p$  to disjoint subsets. These subsets can be also infinite. One can assign to this kind of decomposition a resolution which is the poorer the larger the subsets involved are. p-Adic thermodynamics would represent the situation in which the resolution is maximal since each set contains only single binary digit. Note the analogy with the basis of completely localized wave functions in a lattice.

*c) How to map p-adic transition probabilities to real ones?*

p-Adic variants of TGD, if they exist, give rise to S-matrices and transition probabilities  $P_{ij}$ , which are p-adic numbers.

1. The p-adic probabilities defined by rows of S-matrix mapped to real numbers using canonical identification respecting the  $q = r/s$  decomposition of rational number or its appropriate generalization should define real probabilities.
2. The simplest example would be simple renormalization for the real counterparts of the p-adic probabilities  $(P_{ij})_R$  obtained by canonical identification (or more probably its appropriate modification).

$$\begin{aligned}
 P_{ij} &= \sum_{k \geq 0} P_{ij}^k p^k , \\
 P_{ij} &\rightarrow \sum_{k \geq 0} P_{ij}^k p^{-k} \equiv (P_{ij})_R , \\
 (P_{ij})_R &\rightarrow \frac{(P_{ij})_R}{\sum_j (P_{ij})_R} \equiv P_{ij}^R .
 \end{aligned}
 \tag{2.7}$$

The procedure converges rapidly in powers of  $p$  and resembles renormalization procedure of quantum field theories. The procedure automatically divides away one four-momentum delta function from the square of S-matrix element containing the square of delta function with no well defined mathematical meaning. Usually one gets rid of the delta function interpreting it as the inverse of the four-dimensional measurement volume so that transition rate instead of transition probability is obtained. Of course, also now same procedure should work either as a discrete or a continuous version.

3. Probability interpretation would suggest that the real counterparts of p-adic probabilities sum up to unity. This condition is rather strong since it would hold separately for each row and column of the S-matrix.
4. A further condition would be that the real counterparts of the p-adic probabilities for a given prime  $p$  are identical with the transition probabilities defined by the real S-matrix for real space-time sheets with effective p-adic topology characterized by  $p$ . This condition might allow to deduce all relevant phase information about real and corresponding p-adic S-matrices using as an input only the observable transition probabilities.

d) *What it means that p-adically independent events are not independent in real sense?*

A further condition would be that p-adic quantum transitions represent also in the real sense independent elementary events so that the real counterpart for a sum of the p-adic probabilities for a finite number of transitions equals to the sum of corresponding real probabilities. This condition is definitely too strong in the general case since only a single transition could correspond to a given p-adic norm of transition probability  $P_{ij}$  with  $i$  fixed. In p-adic thermodynamics it can be satisfied if the degeneracy for an energy eigenstate for a given eigen value  $L_0 = n$  is not larger than  $p$ . This condition fails for large values of  $n$  for super Virasoro representations since the degeneracy grows exponentially. This has not practical implications for the large values of  $p$  considered.

The crucial question concerns the physical difference between the real counterpart for the sum of the p-adic transition probabilities and for the sum of the real counterparts of these probabilities, which are in general different:

$$\left(\sum_j P_{ij}\right)_R \neq \sum_j (P_{ij})_R . \quad (2.8)$$

The suggestion is that p-adic sum of the transition probabilities corresponds to the experimental situation, when one does not monitor individual transitions but using some common experimental signature only looks whether the transition leads to this set of the final states or not. When one looks each transition separately or effectively performs different experiment by considering only one transition channel in each experiment one must use the sum of the real probabilities. More precisely, the choice of the experimental signatures divides the set  $U$  of the final states to a disjoint union  $U = \cup_i U_i$  and one must define the real counterparts for the transition probabilities  $P_{iU_k}$  as

$$\begin{aligned} P_{iU_k} &= \sum_{j \in U_k} P_{ij} , \\ P_{iU_k} &\rightarrow (P_{iU_k})_R , \\ (P_{iU_k})_R &\rightarrow \frac{(P_{iU_k})_R}{\sum_l (P_{iU_l})_R} \equiv P_{iU_k}^R . \end{aligned} \quad (2.9)$$

The assumption means deep a departure from the ordinary probability theory. If p-adic physics is the physics of cognitive systems, there need not be anything mysterious in the dependence of the behavior of system on how it is monitored. At least half-jokingly one might argue that the behavior of an intelligent system indeed depends strongly on whether the boss is nearby or not. The precise definition for the monitoring could be based on the decomposition of the density matrix representing the entangled subsystem into a direct sum over the subspaces associated with the degenerate eigenvalues of the density matrix. This decomposition provides a natural definition for the notions of the monitoring and resolution.

The renormalization procedure is in fact familiar from standard physics. Assume that the labels  $j$  correspond to momenta. The division of momentum space to cells of a given size so that the individual momenta inside cells are not monitored separately means that momentum resolution is finite. Therefore one must perform p-adic summation over the cells and define the real probabilities in the proposed manner. p-Adic effects resulting from the difference between p-adic and real summations could be the counterpart of the renormalization effects in QFT. It should be added that similar resolution can be defined also for the initial states by decomposing them into a union of disjoint subsets.

An alternative interpretation for the degenerate eigenvalues has emerged years after writing this. The sub-spaces corresponding to given eigenvalue of density matrix represent entangled states resulting in state function reduction interpreted as measurement of density matrix. This entanglement would be negentropic and represent a rule/concept, whose instances the superposed state pairs are. The information measure would Shannon entropy based on the replacement of the probability appearing as argument of logarithm with its p-adic norm. This entropy would be negative and therefore measure the information associated with the entanglement. This number

theoretic entropy characterizes two particle state rather than single particle state and has nothing to do with the ordinary Shannon entropy.

Maybe one could say that finite measurement resolution implies automatically conceptualization and rule building. Abstractions are indeed obtained by dropping out the details.

2. p-Adic thermodynamics

The p-adic field theory limit as such is not expected to give a realistic theory at elementary particle physics level. The point is that particles are expected to be either massless or possess mass of order  $10^{-4}$  Planck mass. The p-adic description of particle massivation described in [K18] shows that p-adic thermodynamics provides the proper formulation of the problem. What is thermalized is Virasoro generator  $L_0$  (mass squared contribution is not included to  $L_0$  so that states do not have a fixed conformal weight). Temperature is quantized purely number theoretically in low temperature limit ( $exp(H/kT) \rightarrow p^{L_0/T}$ ,  $T = 1/n$ ): in fact, the partition function does not even exist in high temperature phase. The extremely small mixing of massless states with Planck mass states implies massivation and predictions of the p-adic thermodynamics for the fermionic masses are in excellent agreement with experimental masses. Thermodynamic approach also explains the emergence of the length scale  $L_p$  for a given p-adic condensation level and one can develop arguments explaining why primes near prime powers of two are favored.

It should be noticed that rational p-adic temperatures  $1/T = k/n$  are possible, if one poses the restriction that thermal probabilities are non-vanishing only for some subalgebra of the Super Virasoro algebra isomorphic to the Super Virasoro algebra itself. The generators  $L_{kn}, G_{kn}$ , where  $k$  is a positive integer, indeed span this kind of a subalgebra by the fractality of the Super Virasoro algebra and  $p^{L_0/T}$  is integer valued with this restriction.

One might apply thermodynamics approach should also in the calculation of S-matrix. What is needed is thermodynamical expectation value for the transition amplitudes squared over incoming and outgoing states. In this expectation value 3-momenta are fixed and only mass squared varies.

3. Generalization of the notion of information

TGD inspired theory of consciousness, in particular the formulation of Negentropy Maximization Principle (NMP) in p-adic context, has forced to rethink the notion of the information concept. In TGD state preparation process is realized as a sequence of self measurements. Each self measurement means a decomposition of the sub-system involved to two unentangled parts. The decomposition is fixed highly uniquely from the requirement that the reduction of the entanglement entropy is maximal.

The additional assumption is that bound state entanglement is stable against self measurement. This assumption is somewhat ad hoc and it would be nice to get rid of it. The only manner to achieve this seems to be a generalized definition of entanglement entropy allowing to assign a negative value of entanglement entropy to the bound state entanglement, so that bound state entanglement would actually carry information, in fact conscious information (experience of understanding). This would be very natural since macro-temporal quantum coherence corresponds to a generation of bound state entanglement, and is indeed crucial for ability to have long lasting non-entropic mental images.

The generalization of the notion of number concept leads immediately to the basic problem. How to generalize the notion of entanglement entropy that it makes sense for a genuinely p-adic entanglement? What about the number-theoretically universal entanglement with entanglement probabilities, which correspond to finite extension of rational numbers? One can also ask whether the generalized notion of information could make sense at the level of the space-time as suggested by quantum-classical correspondence.

In the real context Shannon entropy is defined for an ensemble with probabilities  $p_n$  as

$$S = - \sum_n p_n \log(p_n) . \tag{2.10}$$

As far as theory of consciousness is considered, the basic problem is that Shannon entropy is always non-negative so that as such it does not define a genuine information measure. One could define information as a change of Shannon entropy and this definition is indeed attractive in the sense that quantum jump is the basic element of conscious experience and involves a change. One can however

argue that the mere ability to transfer entropy to environment (say by aggressive behavior) is not all that is involved with conscious information, and even less so with the experience of understanding or moment of eureka. One should somehow generalize the Shannon entropy without losing the fundamental additivity property.

*a) p-Adic entropies*

The key observation is that in the p-adic context the logarithm function  $\log(x)$  appearing in the Shannon entropy is not defined if the argument of logarithm has p-adic norm different from 1. Situation changes if one uses an extension of p-adic numbers containing  $\log(p)$ : the conjecture is that this extension is finite-dimensional. One might however argue that Shannon entropy should be well defined even without the extension.

p-Adic thermodynamics inspires a way to achieve this. One can replace  $\log(x)$  with the logarithm  $\log_p(|x|_p)$  of the p-adic norm of  $x$ , where  $\log_p$  denotes p-based logarithm. This logarithm is integer valued ( $\log_p(p^n) = n$ ), and is interpreted as a p-adic integer. The resulting p-adic entropy

$$\begin{aligned} S_p &= \sum_n p_n k(p_n) , \\ k(p_n) &= -\log_p(|p_n|) . \end{aligned} \tag{2.11}$$

is additive: that is the entropy for two non-interacting systems is the sum of the entropies of composites. Note that this definition differs from Shannon's entropy by the factor  $\log(p)$ . This entropy vanishes identically in the case that the p-adic norms of the probabilities are equal to one. This means that it is possible to have non-entropic entanglement for this entropy.

One can consider a modification of  $S_p$  using p-adic logarithm if the extension of the p-adic numbers contains  $\log(p)$ . In this case the entropy is formally identical with the Shannon entropy:

$$S_p = -\sum_n p_n \log(p_n) = -\sum_n p_n [-k(p_n)\log(p) + p^{k_n} \log(p_n/p^{k_n})] . \tag{2.12}$$

It seems that this entropy cannot vanish.

One must map the p-adic value entropy to a real number and here canonical identification can be used:

$$\begin{aligned} S_{p,R} &= (S_p)_R \times \log(p) , \\ (\sum_n x_n p^n)_R &= \sum_n x_n p^{-n} . \end{aligned} \tag{2.13}$$

The real counterpart of the p-adic entropy is non-negative.

*b) Number theoretic entropies and metabolic energy*

In the case that the probabilities are rational or belong to a finite-dimensional extension of rationals, it is possible to regard them as real numbers or p-adic numbers in some extension of p-adic numbers for any  $p$ . The visions that rationals and their finite extensions correspond to islands of order in the seas of chaos of real and p-adic transcendentals suggests that states having entanglement coefficients in finite-dimensional extensions of rational numbers are somehow very special. This is indeed the case. The p-adic entropy  $S_p = -\sum_n p_n \log_p(|p_n|)\log(p)$  can be interpreted in this case as an ordinary rational number in an extension containing  $\log(p)$ .

What makes this entropy so interesting is that it can have also negative values in which case the interpretation as an information measure is natural. In the real context one can fix the value of the value of the prime  $p$  by requiring that  $S_p$  is maximally negative, so that the information content of the ensemble could be defined as

$$I \equiv \text{Max}\{-S_p, p \text{ prime}\} . \tag{2.14}$$

This information measure is positive when the entanglement probabilities belong to a finite-dimensional extension of rational numbers. Thus kind of entanglement is stable against NMP [K15], and has a natural interpretation as a negentropic entanglement.

There is no need to interpret negentropic entanglement as bound state entanglement as was the original proposal. This together with the vision about life as something in the intersection of the real and p-adic worlds inspires the idea about a connection between information and metabolism in living matter. Metabolic energy could be carried by negentropic entanglement and the feed of metabolic energy would be also feed of negentropy. In particular, the poorly understood high energy phosphate bond could be identified as a bond involving negentropic entanglement [K1]. The prediction would be that the negentropic states of real systems form a number theoretical hierarchy according to the prime p-adic dimension of algebraic extension characterizing the entanglement.

Number theoretical state function reduction and state preparation could be seen as information generating processes in the intersection of real and p-adic worlds.

### 2.4.2 How to define integration and p-adic Fourier analysis, integral calculus, and p-adic counterparts of geometric objects?

p-Adic differential calculus exists and obeys essentially the same rules as ordinary differential calculus. The only difference from real context is the existence of p-adic pseudo constants: any function which depends on finite number of pinary digits has vanishing p-adic derivative. This implies non-determinism of p-adic differential equations. One can define p-adic integral functions using the fact that indefinite integral is the inverse of differentiation. The basis problem with the definite integrals is that p-adic numbers are not well-ordered so that the crucial ordering of the points of real axis in definite integral is not unique. Also p-adic Fourier analysis is problematic since direct counterparts of  $\exp(ix)$  and trigonometric functions are not periodic. Also  $\exp(-x)$  fails to converge exponentially since it has p-adic norm equal to 1. Note also that these functions exist only when the p-adic norm of  $x$  is smaller than 1.

The following considerations support the view that the p-adic variant of a geometric objects, integration and p-adic Fourier analysis exists but only when one considers highly symmetric geometric objects such as symmetric spaces. This is welcome news from the point of view of physics. At the level of space-time surfaces this is problematic. The field equations associated with Kähler action and Kähler-Dirac equation make sense. Kähler action defined as integral over p-adic space-time surface fails to exist. If however the Kähler function identified as Kähler for a preferred extremal of Kähler action is rational or algebraic function of preferred complex coordinates of WCW with rational coefficients, its p-adic continuation is expected to exist.

#### 1. Circle with rotational symmetries and its hyperbolic counterparts

Consider first circle with emphasis on symmetries and Fourier analysis.

1. In this case angle coordinate  $\phi$  is the natural coordinate. It however does not make sense as such p-adically and one must consider either trigonometric functions or the phase  $\exp(i\phi)$  instead. If one wants to do Fourier analysis on circle one must introduce roots  $U_{n,N} = \exp(in2\pi/N)$  of unity. This means discretization of the circle. Introducing all roots  $U_{n,p} = \exp(i2\pi n/p)$ , such that  $p$  divides  $N$ , one can represent all  $U_{k,n}$  up to  $n = N$ . Integration is naturally replaced with sum by using discrete Fourier analysis on circle. Note that the roots of unity can be expressed as products of powers of roots of unity  $\exp(in2\pi/p^k)$ , where  $p^k$  divides  $N$ .
2. There is a number theoretical delicacy involved. By Fermat's theorem  $a^{p-1} \text{ mod } p = 1$  for  $a = 1, \dots, p-1$  for a given p-adic prime so that for any integer  $M$  divisible by a factor of  $p-1$  the  $M$ :th roots of unity exist as ordinary p-adic numbers. The problem disappears if these values of  $M$  are excluded from the discretization for a given value of the p-adic prime. The manner to achieve this is to assume that  $N$  contains no divisors of  $p-1$  and is consistent with the notion of finite measurement resolution. For instance,  $N = p^n$  is an especially natural choice guaranteeing this.
3. The p-adic integral defined as a Fourier sum does not reduce to a mere discretization of the real integral. In the real case the Fourier coefficients must approach to zero as the wave vector

$k = n2\pi/N$  increases. In the p-adic case the condition consistent with the notion of finite measurement resolution for angles is that the p-adic valued Fourier coefficients approach to zero as  $n$  increases. This guarantees the p-adic convergence of the discrete approximation of the integral for large values of  $N$  as  $n$  increases. The map of p-adic Fourier coefficients to real ones by canonical identification could be used to relate p-adic and real variants of the function to each other.

This finding would suggest that p-adic geometries -in particular the p-adic counterpart of  $CP_2$ , are discrete. Variables which have the character of a radial coordinate are in natural manner p-adically continuous whereas phase angles are naturally discrete and described in terms of algebraic extensions. The conclusion is disappointing since one can quite well argue that the discrete structures can be regarded as real. Is there any manner to escape this conclusion?

1. Exponential function  $exp(ix)$  exists p-adically for  $|x|_p \leq 1/p$  but is not periodic. It provides representation of p-adic variant of circle as group  $U(1)$ . One obtains actually a hierarchy of groups  $U(1)_{p,n}$  corresponding to  $|x|_p \leq 1/p^n$ . One could consider a generalization of phases as products  $Exp_p(N, n2\pi/N + x) = exp(in2\pi n/N)exp(ix)$  of roots of unity and exponent functions with an imaginary exponent. This would assign to each root of unity p-adic continuum interpreted as the analog of the interval between two subsequent roots of unity at circle. The hierarchies of measurement resolutions coming as  $2\pi/p^n$  would be naturally accompanied by increasingly smaller p-adic groups  $U(1)_{p,n}$ .
2. p-Adic integration would involve summation plus possibly also an integration over each p-adic variant of discretization interval. The summation over the roots of unity implies that the integral of  $\int exp(ix)dx$  would appear for  $n = 0$ . Whatever the value of this integral is, it is compensated by a normalization factor guaranteeing orthonormality.
3. If one interprets the p-adic coordinate as p-adic integer without the identification of points differing by a multiple of  $n$  as different points the question whether one should require p-adic continuity arises. Continuity is obtained if  $U_n(x + mp^m) = U_n(x)$  for large values of  $m$ . This is obtained if one has  $n = p^k$ . In the spherical geometry this condition is not needed and would mean quantization of angular momentum as  $L = p^k$ , which does not look natural. If representations of translation group are considered the condition is natural and conforms with the spirit of the p-adic length scale hypothesis.

The hyperbolic counterpart of circle corresponds to the orbit of point under Lorentz group in two 2-D Minkowski space. Plane waves are replaced with exponentially decaying functions of the coordinate  $\eta$  replacing phase angle. Ordinary exponent function  $exp(x)$  has unit p-adic norm when it exists so that it is not a suitable choice. The powers  $p^n$  existing for p-adic integers however approach to zero for large values of  $x = n$ . This forces discretization of  $\eta$  or rather the hyperbolic phase as powers of  $p^x$ ,  $x = n$ . Also now one could introduce products of  $Exp_p(n \log(p) + z) = p^n exp(x)$  to achieve a p-adic continuum. Also now the integral over the discretization interval is compensated by orthonormalization and can be forgotten. The integral of exponential function would reduce to a sum  $\int Exp_p dx = \sum_k p^k = 1/(1-p)$ . One can also introduce finite-dimensional but non-algebraic extensions of p-adic numbers allowing  $e$  and its roots  $e^{1/n}$  since  $e^p$  exists p-adically.

2. Plane with translational and rotational symmetries

Consider first the situation by taking translational symmetries as a starting point. In this case Cartesian coordinates are natural and Fourier analysis based on plane waves is what one wants to define. As in the previous case, this can be done using roots of unity and one can also introduce p-adic continuum by using the p-adic variant of the exponent function. This would effectively reduce the plane to a box. As already noticed, in this case the quantization of wave vectors as multiples of  $1/p^k$  is required by continuity.

One can take also rotational symmetries as a starting point. In this case cylindrical coordinates  $(\rho, \phi)$  are natural.

1. Radial coordinate can have arbitrary values. If one wants to keep the connection  $\rho = \sqrt{x^2 + y^2}$  with the Cartesian picture square root allowing extension is natural. Also the



values of radial coordinate proportional to odd power of  $p$  are problematic since one should introduce  $\sqrt{p}$ : is this extension internally consistent? Does this mean that the points  $\rho \propto p^{2n+1}$  are excluded so that the plane decomposes to annuli?

2. As already found, angular momentum eigen states can be described in terms of roots of unity and one could obtain continuum by allowing also phases defined by p-adic exponent functions.
3. In radial direction one should define the p-adic variants for the integrals of Bessel functions and they indeed might make sense by algebraic continuation if one consistently defines all functions as Fourier expansions. Delta-function renormalization causes technical problems for a continuum of radial wave vectors. One could avoid the problem by using exponentially decaying variants of Bessel function in the regions far from origin, and here the already proposed description of the hyperbolic counterparts of plane waves is suggestive.
4. One could try to understand the situation also using Cartesian coordinates. In the case of sphere this is achieved by introducing two coordinate patches with Cartesian coordinates. Pythagorean phases are rational phases (orthogonal triangles for which all sides are integer valued) and form a dense set on circle. Complex rationals (orthogonal triangles with integer valued short sides) define a more general dense subset of circle. In both cases it is difficult to imagine a discretized version of integration over angles since discretization with constant angle increment is not possible.

### 3. The case of sphere and more general symmetric space

In the case of sphere spherical coordinates are favored by symmetry considerations. For spherical coordinates  $\sin(\theta)$  is analogous to the radial coordinate of plane. Legendre polynomials expressible as polynomials of  $\sin(\theta)$  and  $\cos(\theta)$  are expressible in terms of phases and the integration measure  $\sin^2(\theta)d\theta d\phi$  reduces the integral of  $S^2$  to summation. As before one can introduce also p-adic continuum. Algebraic cutoffs in both angular momentum  $l$  and  $m$  appear naturally. Similar cutoffs appear in the representations of quantum groups and there are good reasons to expect that these phenomena are correlated.

Exponent of Kähler function appears in the integration over WCW. From the expression of Kähler gauge potential given by  $A_\alpha = J_\alpha^\theta \partial_\theta K$  one obtains using  $A_\alpha = \cos(\theta)\delta_{\alpha,\phi}$  and  $J_{\theta\phi} = \sin(\theta)$  the expression  $\exp(K) = \sin(\theta)$ . Hence the exponent of Kähler function is expressible in terms of spherical harmonics.

The completion of the discretized sphere to a p-adic continuum- and in fact any symmetric space- could be performed purely group theoretically.

1. Exponential map maps the elements of the Lie-algebra to elements of Lie-group. This recipe generalizes to arbitrary symmetric space  $G/H$  by using the Cartan decomposition  $g = t + h$ ,  $[h, h] \subset h, [h, t] \subset t, [t, t] \subset h$ . The exponentiation of  $t$  maps  $t$  to  $G/H$  in this case. The exponential map has a p-adic generalization obtained by considering Lie algebra with coefficients with p-adic norm smaller than one so that the p-adic exponent function exists. As a matter fact, one obtains a hierarchy of Lie-algebras corresponding to the upper bounds of the p-adic norm coming as  $p^{-k}$  and this hierarchy naturally corresponds to the hierarchy of angle resolutions coming as  $2\pi/p^k$ . By introducing finite-dimensional transcendental extensions containing roots of  $e$  one obtains also a hierarchy of p-adic Lie-algebras associated with transcendental extensions.
2. In particular, one can exponentiate the complement of the  $SO(2)$  sub-algebra of  $SO(3)$  Lie-algebra in p-adic sense to obtain a p-adic completion of the discrete sphere. Each point of the discretized sphere would correspond to a p-adic continuous variant of sphere as a symmetric space. Similar construction applies in the case of  $CP_2$ . Quite generally, a kind of fractal or holographic symmetric space is obtained from a discrete variant of the symmetric space by replacing its points with the p-adic symmetric space.
3. In the N-fold discretization of the coordinates of M-dimensional space  $t$  one  $(N - 1)^M$  discretization volumes which is the number of points with non-vanishing  $t$ -coordinates. It would

be nice if one could map the p-adic discretization volumes with non-vanishing  $t$ -coordinates to their positive valued real counterparts by applying canonical identification. By group invariance it is enough to show that this works for a discretization volume assignable to the origin. Since the p-adic numbers with norm smaller than one are mapped to the real unit interval, the p-adic Lie algebra is mapped to the unit cell of the discretization lattice of the real variant of  $t$ . Hence by a proper normalization this mapping is possible.

The above considerations suggests that the hierarchies of measurement resolutions coming as  $\Delta\phi = 2\pi/p^n$  are in a preferred role. One must be however cautious in order to avoid too strong assumptions. The above considerations suggest that the hierarchies of measurement resolutions coming as  $\Delta\phi = 2\pi/p^n$  are in a preferred role. One must be however cautious in order to avoid too strong assumptions. The following arguments however support this identification.

1. The vision about p-adicization characterizes finite measurement resolution for angle measurement in the most general case as  $\Delta\phi = 2\pi M/N$ , where  $M$  and  $N$  are positive integers having no common factors. The powers of the phases  $\exp(i2\pi M/N)$  define identical Fourier basis irrespective of the value of  $M$  unless one allows only the powers  $\exp(i2\pi kM/N)$  for which  $kM < N$  holds true: in the latter case the measurement resolutions with different values of  $M$  correspond to different numbers of Fourier components. Otherwise the measurement resolution is just  $\Delta\phi = 2\pi/p^n$ . If one regards  $N$  as an ordinary integer, one must have  $N = p^n$  by the p-adic continuity requirement.
2. One can also interpret  $N$  as a p-adic integer and assume that state function reduction selects one particular prime (no superposition of quantum states with different p-adic topologies). For  $N = p^n M$ , where  $M$  is not divisible by  $p$ , one can express  $1/M$  as a p-adic integer  $1/M = \sum_{k \geq 0} M_k p^k$ , which is infinite as a real integer but effectively reduces to a finite integer  $K(p) = \sum_{k=0}^{N-1} M_k p^k$ . As a root of unity the entire phase  $\exp(i2\pi M/N)$  is equivalent with  $\exp(i2\pi R/p^n)$ ,  $R = K(p)M \pmod{p^n}$ . The phase would non-trivial only for p-adic primes appearing as factors in  $N$ . The corresponding measurement resolution would be  $\Delta\phi = R2\pi/N$ . One could assign to a given measurement resolution all the p-adic primes appearing as factors in  $N$  so that the notion of multi-p p-adicity would make sense. One can also consider the identification of the measurement resolution as  $\Delta\phi = |N/M|_p = 2\pi/p^k$ . This interpretation is supported by the approach based on infinite primes [K26].

4. What about integrals over partonic 2-surfaces and space-time surface?

One can of course ask whether also the integrals over partonic 2-surfaces and space-time surface could be p-adicized by using the proposed method of discretization. Consider first the p-adic counterparts of the integrals over the partonic 2-surface  $X^2$ .

1. WCW Hamiltonians and Kähler form are expressible using flux Hamiltonians defined in terms of  $X^2$  integrals of  $JH_A$ , where  $H_A$  is  $\delta CD \times CP_2$  Hamiltonian, which is a rational function of the preferred coordinates defined by the exponentials of the coordinates of the sub-space  $t$  in the appropriate Cartan algebra decomposition. The flux factor  $J = \epsilon^{\alpha\beta} J_{\alpha\beta} \sqrt{g_2}$  is scalar and does not actually depend on the induced metric.
2. The notion of finite measurement resolution would suggest that the discretization of  $X^2$  is somehow induced by the discretization of  $\delta CD \times CP_2$ . The coordinates of  $X^2$  could be taken to be the coordinates of the projection of  $X^2$  to the sphere  $S^2$  associated with  $\delta M_{\pm}^4$  or to the homologically non-trivial geodesic sphere of  $CP_2$  so that the discretization of the integral would reduce to that for  $S^2$  and to a sum over points of  $S^2$ .
3. To obtain an algebraic number as an outcome of the summation, one must pose additional conditions guaranteeing that both  $H_A$  and  $J$  are algebraic numbers at the points of discretization (recall that roots of unity are involved). Assume for definiteness that  $S^2$  is  $r_M = \text{constant}$  sphere. If the remaining preferred coordinates are functions of the preferred  $S^2$  coordinates mapping phases to phases at discretization points, one obtains the desired outcome. These conditions are rather strong and mean that the various angles defining  $CP_2$  coordinates -at least the two cyclic angle coordinates- are integer multiples of those assignable to  $S^2$  at the

points of discretization. This would be achieved if the preferred complex coordinates of  $CP_2$  are powers of the preferred complex coordinate of  $S^2$  at these points. One could say that  $X^2$  is algebraically continued from a rational surface in the discretized variant of  $\delta CD \times CP_2$ . Furthermore, if the measurement resolutions come as  $2\pi/p^n$  as p-adic continuity actually requires and if they correspond to the p-adic group  $G_{p,n}$  for which group parameters satisfy  $|t|_p \leq p^{-n}$ , one can precisely characterize how a p-adic prime characterizes the real partonic 2-surface. This would be a fulfilment of one of the oldest dreams related to the p-adic vision.

A even more ambitious dream would be that even the integral of the Kähler action for preferred extremals could be defined using a similar procedure. The conjectured slicing of Minkowskian space-time sheets by string world sheets and partonic 2-surfaces encourages these hopes.

1. One could introduce local coordinates of  $H$  at both ends of CD by introducing a continuous slicing of  $M^4 \times CP_2$  by the translates of  $\delta M_{\pm}^4 \times CP_2$  in the direction of the time-like vector connecting the tips of CD. As space-time coordinates one could select four of the eight coordinates defining this slicing. For instance, for the regions of the space-time sheet representable as maps  $M^4 \rightarrow CP_2$  one could use the preferred  $M^4$  time coordinate, the radial coordinate of  $\delta M_{\pm}^4$ , and the angle coordinates of  $r_M = \text{constant}$  sphere.
2. Kähler action density should have algebraic values and this would require the strengthening of the proposed conditions for  $X^2$  to apply to the entire slicing meaning that the discretized space-time surface is a rational surface in the discretized  $CD \times CP_2$ . If this condition applies to the entire space-time surface it would effectively mean the discretization of the classical physics to the level of finite geometries. This seems quite strong implication but is consistent with the preferred extremal property implying the generalized Bohr rules.

5. Tentative conclusions

These findings suggest following conclusions.

1. Exponent functions play a key role in the proposed p-adicization. This is not an accident since exponent functions play a fundamental role in group theory and p-adic variants of real geometries exist only under symmetries- possibly maximal possible symmetries- since otherwise the notion of Fourier analysis making possible integration does not exist. The inner product defined in terms of integration reduce for functions representable in Fourier basis to sums and can be carried out by using orthogonality conditions. Convolution involving integration reduces to a product for Fourier components. In the case of embedding space and WCW these conditions are satisfied but for space-time surfaces this is not possible.
2. There are several ways to choose the Cartan algebra already in the case of sphere. In the case of plane one can consider either translations or rotations and this leads to different p-adic variants of plane. Also the realization of the hierarchy of Planck constants leads to the conclusion that the extended embedding space and therefore also WCW contains sectors corresponding to different choices of quantization axes meaning that quantum measurement has a direct geometric correlate. One can imagine also other discretizations and choices of preferred coordinates and the interpretation is that they correspond to different cognitive representations and to different p-adic physics. This means a refinement of General Coordinate Invariance taking into account cognition.
3. The above described 2-D examples represent symplectic geometries for which one has natural decomposition of coordinates to canonical pairs of cyclic coordinate (phase angle) and corresponding canonical conjugate coordinate. p-Adicization depends on whether the conjugate corresponds to an angle or non-compact coordinate. In both cases it is however possible to define integration. For instance, in the case of  $CP_2$  one would have two canonically conjugate pairs and one can define the p-adic counterparts of  $CP_2$  partial waves by generalizing the procedure applied to spherical harmonics. Products of functions expressible using partial waves can be decomposed by tensor product decomposition to spherical harmonics and can be integrated. In particular inner products can be defined as integrals. The Hamiltonians generating isometries are rational functions of phases: this inspires the hope that also WCW

Hamiltonians also rational functions of preferred WCW coordinates and thus allow p-adic variants.

4. Discretization by introducing algebraic extensions seems unavoidable in the p-adicization of geometrical objects but one can have p-adic continuum as the analog of the discretization interval and in the function basis expressible in terms of phase factors and p-adic counterparts of exponent functions. As already described, the exponential map for Lie group provide an elegant manner to realize this. This would give a precise meaning for the p-adic counterparts of the embedding space and WCW if the latter is a symmetric space allowing coordinatization in terms of phase angles and conjugate coordinates. The intersection of p-adic and real worlds in a given measurement resolution would be unique and correspond to the points defining the discretization.

### 2.4.3 p-Adic embedding space

The construction of both quantum TGD and p-adic QFT limit requires p-adicization of the embedding space geometry. Also the fact that p-adic Poincare invariance throws considerable light to the p-adic length scale hypothesis suggests that p-adic geometry is really needed. The construction of the p-adic version of the embedding space geometry and spinor structure relies on the symmetry arguments and to the generalization of the analytic formulas of the real case almost. The essential element is the notion of finite measurement resolution leading to discretization in large and to p-adicization below the resolution scale. This approach leads to a highly nontrivial generalization of the symmetry concept and p-adic Poincare invariance throws light to the p-adic length scale hypothesis. An important delicacy is related to the identification of the fundamental p-adic length scale, which corresponds to the unit element of the p-adic number field and is mapped to the unit element of the real number field in the canonical identification mapping p-adic mass squared to its real counterpart.

#### 1. *p-Adic Riemannian geometry depends on cognitive representation*

p-Adic Riemann geometry is a direct formal generalization of the ordinary Riemann geometry. In the minimal purely algebraic generalization one does not try to define concepts like arch length and volume involving definite integrals but simply defines the p-adic geometry via the metric identified as a quadratic form in the tangent space of the p-adic manifold. Canonical identification would make it possible to define p-adic variant of Riemann integral formally allowing to calculate arc lengths and similar quantities but looks like a trick. The realization that the p-adic variant of harmonic analysis makes it possible to define definite integrals in the case of symmetric space became possible only after a detailed vision about what quantum TGD is [K37] had emerged.

Symmetry considerations dictate the p-adic counterpart of the Riemann geometry for  $M_+^4 \times CP_2$  to a high degree but not uniquely. This non-uniqueness might relate to the distinction between different cognitive representations. For instance, in the case of Euclidian plane one can introduce linear or cylindrical coordinates and the manifest symmetries dictating the preferred coordinates correspond to translational and rotational symmetries in these two cases and give rise to different p-adic variants of the plane. Both linear and cylindrical coordinates are fixed only modulo the action of group consisting of translations and rotations and the degeneracy of choices can be interpreted in terms of a choice of quantization axes of angular momentum and momenta.

The most natural looking manner to define the p-adic counterpart of  $M^4$  is by using a p-adic completion for a subset of rational points in coordinates which are preferred on physical basis. In case of  $M^4$  linear Minkowski coordinates are an obvious choice but also the counterparts of Robertson-Walker coordinates for  $M_+^4$  defined as  $[t, (z, x, y)] = a \times [\cosh(\eta), \sinh(\eta)(\cos(\theta), \sin(\theta)\cos(\phi), \sin(\theta)\sin(\phi))]$  expressible in terms of phases and their hyperbolic counterparts and transforming nicely under the Cartan algebra of Lorentz group are possible. p-Adic variant is obtained by introducing finite measurement resolution for angle and replacing angle range by finite number of roots of unity. Same applies to hyperbolic angles.

Rational  $CP_2$  could be defined as a coset space  $SU(3, Q)/U(2, Q)$  associated with complex rational unitary  $3 \times 3$ -matrices.  $CP_2$  could be defined as coset space of complex rational matrices by choosing one point in each coset  $SU(3, Q)/U(2, Q)$  as a complex rational  $3 \times 3$ -matrix representable in terms of Pythagorean phases [A11] and performing a completion for the elements of this matrix

by multiplying the elements with the p-adic exponentials  $\exp(iu)$ ,  $|u|_p < 1$  such that one obtains p-adically unitary matrix.

This option is not very natural as far as integration is considered.  $CP_2$  however allows the analog of spherical coordinates for  $S^2$  expressible in terms of angle variables alone and this suggests the introduction of the variant of  $CP_2$  for which the coordinate values correspond to roots of unity. Completion would be performed in the same manner as for rational  $CP_2$ . This non-uniqueness need not be a drawback but could reflect the fact that the p-adic cognitive representation of real geometry are geometrically non-equivalent. This means a refinement of the principle of General Coordinate Invariance taking into account the fact that the cognitive representation of the real world affects the world with cognition included in a delicate manner.

2. The identification of the fundamental p-adic length scale

The fundamental p-adic length scale corresponds to the p-adic unit  $e = 1$  and is mapped to the unit of the real numbers in the canonical identification. The correct physical identification of the fundamental p-adic length scale is of crucial importance since the predictions of the theory for p-adic masses depend on the choice of this scale.

In TGD the “radius”  $R$  of  $CP_2$  is the fundamental length scale ( $2\pi R$  is by definition the length of the  $CP_2$  geodesics). In accordance with the idea that p-adic QFT limit makes sense only above length scales larger than the radius of  $CP_2$   $R$  is of same order of magnitude as the p-adic length scale defined as  $l = \pi/m_0$ , where  $m_0$  is the fundamental mass scale and related to the “cosmological constant”  $\Lambda$  ( $R_{ij} = \Lambda s_{ij}$ ) of  $CP_2$  by

$$m_0^2 = 2\Lambda . \tag{2.15}$$

The relationship between  $R$  and  $l$  is uniquely fixed:

$$R^2 = \frac{3}{m_0^3} = \frac{3}{2\Lambda} = \frac{3l^2}{\pi^2} . \tag{2.16}$$

Consider now the identification of the fundamental length scale.

1. One must use  $R^2$  or its integer multiple, rather than  $l^2$ , as the fundamental p-adic length scale squared in order to avoid the appearance of the p-adically ill defined  $\pi$ :s in various formulas of  $CP_2$  geometry.
2. The identification for the fundamental length scale as  $1/m_0$  leads to difficulties.
  - (a) The p-adic length for the  $CP_2$  geodesic is proportional to  $\sqrt{3}/m_0$ . For the physically most interesting p-adic primes satisfying  $p \bmod 4 = 3$  so that  $\sqrt{-1}$  does not exist as an ordinary p-adic number,  $\sqrt{3} = i\sqrt{-3}$  belongs to the complex extension of the p-adic numbers. Hence one has troubles in getting real length for the  $CP_2$  geodesic.
  - (b) If  $m_0^2$  is the fundamental mass squared scale then general quark states have mass squared, which is integer multiple of  $1/3$  rather than integer valued as in string models.
3. These arguments suggest that the correct choice for the fundamental length scale is as  $1/R$  so that  $M^2 = 3/R^2$  appearing in the mass squared formulas is p-adically real and all values of the mass squared are integer multiples of  $1/R^2$ . This does not affect the real counterparts of the thermal expectation values of the mass squared in the lowest p-adic order but the effects, which are due to the modulo arithmetics, are seen in the higher order contributions to the mass squared. As a consequence, one must identify the p-adic length scale  $l$  as

$$l \equiv \pi R ,$$

rather than  $l = \pi/m_0$ . This is indeed a very natural identification. What is especially nice is that this identification also leads to a solution of some longstanding problems related to the p-adic mass calculations. It would be highly desirable to have the same p-adic temperature

$T_p = 1$  for both the bosons and fermions rather than  $T_p = 1/2$  for bosons and  $T_p = 1$  for fermions. For instance, black hole elementary particle analogy as well as the need to get rid of light boson exotics suggests this strongly. It indeed turns out possible to achieve this with the proposed identification of the fundamental mass squared scale.

3. *p-Adic counterpart of  $M_+^4$*

The construction of the p-adic counterpart of  $M_+^4$  seems a relatively straightforward task and should reduce to the construction of the p-adic counterpart of the real axis with the standard metric. As already noticed, linear Minkowski coordinates are physically and mathematically preferred coordinates and it is natural to construct the metric in these coordinates.

There are some quite interesting delicacies related to the p-adic version of the Poincare invariance. Consider first translations. In order to have imaginary unit needed in the construction of the ordinary representations of the Poincare group one must have  $p \bmod 4 = 3$  to guarantee that  $\sqrt{-1}$  does not exist as an ordinary p-adic number. It however seems that the construction of the representations is at least formally possible by replacing imaginary unit with the square root of some other p-adic number not existing as a p-adic number.

It seems that only the discrete group of translations allows representations consisting of orthogonal plane waves. p-Adic plane waves can be defined in the lattice consisting of the multiples of  $x_0 = m/n$  consisting of points with p-adic norm not larger than  $|x_0|_p$  and the points  $p^n x_0$  define fractally scaled-down versions of this set. In canonical identification these sets corresponds to volumes scaled by factors  $p^{-n}$ .

A physically interesting question is whether the Lorentz group should contain only the elements obtained by exponentiating the Lie-algebra generators of the Lorentz group or whether also large Lorentz transformations, containing as a subgroup the group of the rational Lorentz transformations, should be allowed. If the group contains only small Lorentz transformations, the quantization volume of  $M_+^4$  (say the points with coordinates  $m^k$  having p-adic norm not larger than one) is also invariant under Lorentz transformations. This means that the quantization of the theory in the p-adic cube  $|m^k| < p^n$  is a Poincare invariant procedure unlike in the real case.

The appearance of the square root of  $p$ , rather than the naïvely expected  $p$ , in the expression of the p-adic length scale can be understood if the p-adic version of  $M^4$  metric contains  $p$  as a scaling factor:

$$\begin{aligned} ds^2 &= pR^2 m_{kl} dm^k dm^l \ , \\ R &\leftrightarrow 1 \ , \end{aligned} \tag{2.17}$$

where  $m_{kl}$  is the standard  $M^4$  metric  $(1, -1, -1, -1)$ . The p-adic distance function is obtained by integrating the line element using p-adic integral calculus and this gives for the distance along the k:th coordinate axis the expression

$$s = R\sqrt{p}m^k \ . \tag{2.18}$$

The map from p-adic  $M^4$  to real  $M^4$  is canonical identification plus a scaling determined from the requirement that the real counterpart of an infinitesimal p-adic geodesic segment is same as the length of the corresponding real geodesic segment:

$$m^k \rightarrow \pi(m^k)_R \ . \tag{2.19}$$

The p-adic distance along the k:th coordinate axis from the origin to the point  $m^k = (p-1)(1+p+p^2+\dots) = -1$  on the boundary of the set of the p-adic numbers with norm not larger than one, corresponds to the fundamental p-adic length scale  $L_p = \sqrt{p}l = \sqrt{p}\pi R$ :

$$\sqrt{p}((p-1)(1+p+\dots))R \rightarrow \pi R \frac{(p-1)(1+p^{-1}+p^{-2}+\dots)}{\sqrt{p}} = L_p \ . \tag{2.20}$$

What is remarkable is that the shortest distance in the range  $m^k = 1, ..m - 1$  is actually  $L/\sqrt{p}$  rather than  $l$  so that p-adic numbers in range span the entire  $R_+$  at the limit  $p \rightarrow \infty$ . Hence p-adic topology approaches real topology in the limit  $p \rightarrow \infty$  in the sense that the length of the discretization step approaches to zero.

4. The two variants of  $CP_2$

As noticed,  $CP_2$  allows two variants based on rational discretization and on the discretization based on roots of unity. The root of unity option corresponds to the phases associated with  $1/(1+r^2) = \tan^2(u/2) = (1-\cos(u))/(1+\cos(u))$  and implies that integrals of spherical harmonics can be reduced to summations when angular resolution  $\Delta u = 2\pi/N$  is introduced. In the p-adic context, one can replace distances with trigonometric functions of distances along zig zag curves connecting the points of the discretization. Physically this notion of distance is quite reasonable since distances are often measured using interferometer.

In the case of rtional variant of  $CP_2$  one can proceed by defining the p-adic counterparts of  $SU(3)$  and  $U(2)$  and using the identification  $CP_2 = SU(3)/U(2)$ . The p-adic counterpart of  $SU(3)$  consists of all  $3 \times 3$  unitary matrices satisfying p-adic unitarity conditions (rows/columns are mutually orthogonal unit vectors) or its suitable subgroup: the minimal subgroup corresponds to the exponentials of the Lie-algebra generators. If one allows algebraic extensions of the p-adic numbers, one obtains several extensions of the group. The extension allowing the square root of a p-adically real number is the most interesting one in this respect since the general solution of the unitarity conditions involves square roots.

The subgroup of  $SU(3)$  obtained by exponentiating the Lie-algebra generators of  $SU(3)$  normalized so that their non-vanishing elements have unit p-adic norm, is of the form

$$SU(3)_0 = \{x = \exp(\sum_k it_k X_k) ; |t_k|_p < 1\} = \{x = 1 + iy ; |y|_p < 1\} . \tag{2.21}$$

The diagonal elements of the matrices in this group are of the form  $1 + O(p)$ . In order  $O(p)$  these matrices reduce to unit matrices.

Rational  $SU(3)$  matrices do not in general allow a representation as an exponential. In the real case all  $SU(3)$  matrices can be obtained from diagonalized matrices of the form

$$h = \text{diag}\{\exp(i\phi_1), \exp(i\phi_2), \exp(\exp(-i(\phi_1 + \phi_2)))\} . \tag{2.22}$$

The exponentials are well defined provided that one has  $|\phi_i|_p < 1$  and in this case the diagonal elements are of form  $1 + O(p)$ . For  $p \bmod 4 = 3$  one can however consider much more general diagonal matrices

$$h = \text{diag}\{z_1, z_2, z_3\} ,$$

for which the diagonal elements are rational complex numbers

$$z_i = \frac{(m_i + in_i)}{\sqrt{m_i^2 + n_i^2}} \tag{2.23}$$

satisfying  $z_1 z_2 z_3 = 1$  such that the components of  $z_i$  are integers in the range  $(0, p - 1)$  and the square roots appearing in the denominators exist as ordinary p-adic numbers. These matrices indeed form a group as is easy to see. By acting with  $SU(3)_0$  to each element of this group and by applying all possible automorphisms  $h \rightarrow ghg^{-1}$  using rational  $SU(3)$  matrices one obtains entire  $SU(3)$  as a union of an infinite number of disjoint components.

The simplest (unfortunately not physical) possibility is that the “physical”  $SU(3)$  corresponds to the connected component of  $SU(3)$  represented by the matrices, which are unit matrices in order  $O(p)$ . In this case the construction of  $CP_2$  is relatively straightforward and the real formalism

should generalize as such. In particular, for  $p \bmod 4 = 3$  it is possible to introduce complex coordinates  $\xi_1, \xi_2$  using the complexification for the Lie-algebra complement of  $su(2) \times u(1)$ . The real counterparts of these coordinates vary in the range  $[0, 1)$  and the end points correspond to the values of  $t_i$  equal to  $t_i = 0$  and  $t_i = -p$ . The p-adic sphere  $S^2$  appearing in the definition of the p-adic light cone is obtained as a geodesic sub-manifold of  $CP_2$  ( $\xi_1 = \xi_2$  is one possibility). From the requirement that real  $CP_2$  can be mapped to its p-adic counterpart it is clear that one must allow all connected components of  $CP_2$  obtained by applying discrete unitary matrices having no exponential representation to the basic connected component. In practice this corresponds to the allowance of all possible values of the p-adic norm for the components of the complex coordinates  $\xi_i$  of  $CP_2$ .

The simplest approach to the definition of the  $CP_2$  metric is to replace the expression of the Kähler function in the real context with its p-adic counterpart. In standard complex coordinates for which the action of  $U(2)$  subgroup is linear, the expression of the Kähler function reads as

$$\begin{aligned} K &= \log(1 + r^2) , \\ r^2 &= \sum_i \bar{\xi}_i \xi_i . \end{aligned} \tag{2.24}$$

p-Adic logarithm exists provided  $r^2$  is of order  $O(p)$ . This is the case when  $\xi_i$  is of order  $O(p)$ . The definition of the Kähler function in a more general case, when all possible values of the p-adic norm are allowed for  $r$ , is based on the introduction of a p-adic pseudo constant  $C$  to the argument of the Kähler function

$$K = \log\left(\frac{1 + r^2}{C}\right) . \tag{2.25}$$

$C$  guarantees that the argument is of the form  $\frac{1+r^2}{C} = 1 + O(p)$  allowing a well-defined p-adic logarithm. This modification of the Kähler function leaves the definition of Kähler metric, Kähler form and spinor connection invariant.

A more elegant manner to avoid the difficulty is to use the exponent  $\Omega = \exp(K) = 1 + r^2$  of the Kähler function instead of Kähler function, which indeed well defined for all coordinate values. In terms of  $\Omega$  one can express the Kähler metric as

$$g_{k\bar{l}} = \frac{\partial_k \partial_{\bar{l}} \Omega}{\Omega} - \frac{\partial_k \Omega \partial_{\bar{l}} \Omega}{\Omega^2} . \tag{2.26}$$

The p-adic metric can be defined as

$$s_{i\bar{j}} = R^2 \partial_i \partial_{\bar{j}} K = R^2 \frac{(\delta_{i\bar{j}} r^2 - \bar{\xi}_i \xi_j)}{(1 + r^2)^2} . \tag{2.27}$$

The expression for the Kähler form is the same as in the real case and the components of the Kähler form in the complex coordinates are numerically equal to those of the metric apart from the factor of  $i$ . The components in arbitrary coordinates can be deduced from these by the standard transformation formulas.

## 2.5 What Could Be The Origin Of Preferred P-Adic Primes And P-Adic Length Scale Hypothesis?

p-Adic mass calculations [K18] allow to conclude that elementary particles correspond to one or possible several preferred primes assigning p-adic effective topology to the real space-time sheets in discretization in some length scale range. TGD inspired theory of consciousness leads to the identification of p-adic physics as physics of cognition. Quite recent progress (2015) leads to the



proposal that quantum TGD is adelic: all p-adic number fields are involved and each gives one particular view about physics.

Adelic approach [K12, K22] plus the view about evolution as emergence of increasingly complex extensions of rationals leads to a possible answer to the question of the title. The algebraic extensions of rationals are characterized by preferred rational primes, namely those which are ramified when expressed in terms of the primes of the extensions. These primes would be natural candidates for preferred p-adic primes. An argument relying on what I call weak form of NMP in turn allows to understand why primes near powers of 2 are preferred: as a matter of fact, also primes near powers of other primes are predicted to be favoured.

### 2.5.1 Earlier attempts

How the preferred primes emerge in TGD framework? I have made several attempts to answer this question. As a matter of fact, the question has been slightly different: what determines the p-adic prime assigned to elementary particle by p-adic mass calculations [K13]. The recent view assigns to particle entire adèle but some p-adic number fields in it are different.

1. Classical non-determinism at space-time level for real space-time sheets could in some length scale range involving rational discretization for space-time surface itself or for parameters characterizing it as a preferred extremal correspond to the non-determinism of p-adic differential equations due to the presence of pseudo constants which have vanishing p-adic derivative. Pseudo-constants are functions depend on finite number of binary digits of its arguments.
2. The quantum criticality of TGD [?] is suggested to be realized in terms of infinite hierarchies of super-symplectic symmetry breakings in the sense that only a sub-algebra with conformal weights which are  $n$ -ples of those for the entire algebra act as conformal gauge symmetries [K24]. This might be true for all conformal algebras involved. One has fractal hierarchy since the sub-algebras in question are isomorphic: only the scale of conformal gauge symmetry increases in the phase transition increasing  $n$ . The hierarchies correspond to sequences of integers  $n(i)$  such that  $n(i)$  divides  $n(i+1)$ . These hierarchies would very naturally correspond to hierarchies of inclusions of hyper-finite factors and  $m(i) = n(i+1)/n(i)$  could correspond to the integer  $n$  characterizing the index of inclusion, which has value  $n \geq 3$ . Possible problem is that  $m(i) = 2$  would not correspond to Jones inclusion. Why the scaling by power of two would be different? The natural question is whether the primes dividing  $n(i)$  or  $m(i)$  could define the preferred primes.
3. Negentropic entanglement corresponds to entanglement for which density matrix is projector [K15]. For  $n$ -dimensional projector any prime  $p$  dividing  $n$  gives rise to negentropic entanglement in the sense that the number theoretic entanglement entropy defined by Shannon formula by replacing  $p_i$  in  $\log(p_i) = \log(1/n)$  by its p-adic norm  $N_p(1/n)$  is negative if  $p$  divides  $n$  and maximal for the prime for which the dividing power of prime is largest power-of-prime factor of  $n$ . The identification of p-adic primes as factors of  $n$  is highly attractive idea. The obvious question is whether  $n$  corresponds to the integer characterizing a level in the hierarchy of conformal symmetry breakings.
4. The adelic picture about TGD led to the question whether the notion of unitarity could be generalized. S-matrix would be unitary in adelic sense in the sense that  $P_m = (SS^\dagger)_{mm} = 1$  would generalize to adelic context so that one would have product of real norm and p-adic norms of  $P_m$ . In the intersection of the realities and p-adicities  $P_m$  for reals would be rational and if real and p-adic  $P_m$  correspond to the same rational, the condition would be satisfied. The condition that  $P_m \leq 1$  seems however natural and forces separate unitarity in each sector so that this option seems too tricky.

These are the basic ideas that I have discussed hitherto.

### 2.5.2 Could preferred primes characterize algebraic extensions of rationals?

The intuitive feeling is that the notion of preferred prime is something extremely deep and the deepest thing I know is number theory. Does one end up with preferred primes in number theory?

This question brought to my mind the notion of ramification of primes (see <http://tinyurl.com/hdd1jlf>) (more precisely, of prime ideals of number field in its extension), which happens only for special primes in a given extension of number field, say rationals. Could this be the mechanism assigning preferred prime(s) to a given elementary system, such as elementary particle? I have not considered their role earlier also their hierarchy is highly relevant in the number theoretical vision about TGD.

1. Stating it very roughly (I hope that mathematicians tolerate this language): As one goes from number field  $K$ , say rationals  $Q$ , to its algebraic extension  $L$ , the original prime ideals in the so called integral closure (see <http://tinyurl.com/js6fpvr>) over integers of  $K$  decompose to products of prime ideals of  $L$  (prime is a more rigorous manner to express primeness).

Integral closure for integers of number field  $K$  is defined as the set of elements of  $K$ , which are roots of some monic polynomial with coefficients, which are integers of  $K$  and having the form  $x^n + a_{n-1}x^{n-1} + \dots + a_0$ . The integral closures of both  $K$  and  $L$  are considered. For instance, integral closure of algebraic extension of  $K$  over  $K$  is the extension itself. The integral closure of complex numbers over ordinary integers is the set of algebraic numbers.

2. There are two further basic notions related to ramification and characterizing it. Relative discriminant is the ideal divided by all ramified ideals in  $K$  and relative different is the ideal of  $L$  divided by all ramified  $P_i$ :s. Note that the general ideal is analog of integer and these ideas represent the analogous of product of preferred primes  $P$  of  $K$  and primes  $P_i$  of  $L$  dividing them.
3. A physical analogy is provided by decomposition of hadrons to valence quarks. Elementary particles becomes composite of more elementary particles in the extension. The decomposition to these more elementary primes is of form  $P = \prod P_i^{e(i)}$ , where  $e_i$  is the ramification index - the physical analog would be the number of elementary particles of type  $i$  in the state (see <http://tinyurl.com/h9528p1>). Could the ramified rational primes could define the physically preferred primes for a given elementary system?

In TGD framework the extensions of rationals (see <http://tinyurl.com/h9528p1>) and p-adic number fields (see <http://tinyurl.com/zq22tvb>) are unavoidable and interpreted as an evolutionary hierarchy physically and cosmological evolution would have gradually proceeded to more and more complex extensions. One can say that string world sheets and partonic 2-surfaces with parameters of defining functions in increasingly complex extensions of prime emerge during evolution. Therefore ramifications and the preferred primes defined by them are unavoidable. For p-adic number fields the number of extensions is much smaller for instance for  $p > 2$  there are only 3 quadratic extensions.

1. In p-adic context a proper definition of counterparts of angle variables as phases allowing definition of the analogs of trigonometric functions requires the introduction of algebraic extension giving rise to some roots of unity. Their number depends on the angular resolution. These roots allow to define the counterparts of ordinary trigonometric functions - the naïve generalization based on Taylors series is not periodic - and also allows to defined the counterpart of definite integral in these degrees of freedom as discrete Fourier analysis. For the simplest algebraic extensions defined by  $x^n - 1$  for which Galois group is abelian are unramified so that something else is needed. One has decomposition  $P = \prod P_i^{e(i)}$ ,  $e(i) = 1$ , analogous to  $n$ -fermion state so that simplest cyclic extension does not give rise to a ramification and there are no preferred primes.
2. What kind of polynomials could define preferred algebraic extensions of rationals? Irreducible polynomials are certainly an attractive candidate since any polynomial reduces to a product of them. One can say that they define the elementary particles of number theory. Irreducible polynomials have integer coefficients having the property that they do not decompose to products of polynomials with rational coefficients. IT would be wrong to say that only these algebraic extensions can appear but there is a temptation to say that one can reduce the study of extensions to their study. One can even consider the possibility that string world sheets associated with products of irreducible polynomials are unstable against decay to those characterize irreducible polynomials.

3. What can one say about irreducible polynomials? Eisenstein criterion (see <http://tinyurl.com/47kxjz>) states following. If  $Q(x) = \sum_{k=0, \dots, n} a_k x^k$  is  $n$ :th order polynomial with integer coefficients and with the property that there exists at least one prime dividing all coefficients  $a_i$  except  $a_n$  and that  $p^2$  does not divide  $a_0$ , then  $Q$  is irreducible. Thus one can assign one or more preferred primes to the algebraic extension defined by an irreducible polynomial  $Q$  of this kind - in fact any polynomial allowing ramification. There are also other kinds of irreducible polynomials since Eisenstein's condition is only sufficient but not necessary.
4. Furthermore, in the algebraic extension defined by  $Q$ , the prime ideals  $P$  having the above mentioned characteristic property decompose to an  $n$  :th power of single prime ideal  $P_i$ :  $P = P_i^n$ . The primes are maximally/completely ramified. The physical analog  $P = P_0^n$  is Bose-Einstein condensate of  $n$  bosons. There is a strong temptation to identify the preferred primes of irreducible polynomials as preferred p-adic primes.

A good illustration is provided by equations  $x^2 + 1 = 0$  allowing roots  $x_{\pm} = \pm i$  and equation  $x^2 + 2px + p = 0$  allowing roots  $x_{\pm} = -p \pm \sqrt{p}p - 1$ . In the first case the ideals associated with  $\pm i$  are different. In the second case these ideals are one and the same since  $x_+ = -x_- + p$ : hence one indeed has ramification. Note that the first example represents also an example of irreducible polynomial, which does not satisfy Eisenstein criterion. In more general case the  $n$  conditions on defined by symmetric functions of roots imply that the ideals are one and same when Eisenstein conditions are satisfied.

5. What does this mean in p-adic context? The identity of the ideals can be stated by saying  $P = P_0^n$  for the ideals defined by the primes satisfying the Eisenstein condition. Very loosely one can say that the algebraic extension defined by the root involves  $n$ :th root of p-adic prime  $p$ . This does not work! Extension would have a number whose  $n$ :th power is zero modulo  $p$ . On the other hand, the p-adic numbers of the extension modulo  $p$  should be finite field but this would not be field anymore since there would exist a number whose  $n$ :th power vanishes. The algebraic extension simply does not exist for preferred primes. The physical meaning of this will be considered later.
6. What is so nice that one could readily construct polynomials giving rise to given preferred primes. The complex roots of these polynomials could correspond to the points of partonic 2-surfaces carrying fermions and defining the ends of boundaries of string world sheet. It must be however emphasized that the form of the polynomial depends on the choices of the complex coordinate. For instance, the shift  $x \rightarrow x + 1$  transforms  $(x^n - 1)/(x - 1)$  to a polynomial satisfying the Eisenstein criterion. One should be able to fix allowed coordinate changes in such a way that the extension remains irreducible for all allowed coordinate changes.

Already the integral shift of the complex coordinate affects the situation. It would seem that only the action of the allowed coordinate changes must reduce to the action of Galois group permuting the roots of polynomials. A natural assumption is that the complex coordinate corresponds to a complex coordinate transforming linearly under subgroup of isometries of the embedding space.

In the general situation one has  $P = \prod P_i^{e(i)}$ ,  $e(i) \geq 1$  so that also now there are preferred primes so that the appearance of preferred primes is completely general phenomenon.

### 2.5.3 A connection with Langlands program?

In Langlands program (see <http://tinyurl.com/ycej7s43>) [A28, A27] the great vision is that the  $n$ -dimensional representations of Galois groups  $G$  characterizing algebraic extensions of rationals or more general number fields define  $n$ -dimensional adelic representations of adelic Lie groups, in particular the adelic linear group  $Gl(n, A)$ . This would mean that it is possible to reduce these representations to a number theory for adèles. This would be highly relevant in the vision about TGD as a generalized number theory. I have speculated with this possibility earlier [K12] but the mathematics is so horribly abstract that it takes decade before one can have even hope of building a rough vision.

One can wonder whether the irreducible polynomials could define the preferred extensions  $K$  of rationals such that the maximal abelian extensions of the fields  $K$  would in turn define the adèles

utilized in Langlands program. At least one might hope that everything reduces to the maximally ramified extensions.

At the level of TGD string world sheets with parameters in an extension defined by an irreducible polynomial would define an adèle containing various p-adic number fields defined by the primes of the extension. This would define a hierarchy in which the prime ideals of previous level would decompose to those of the higher level. Each irreducible extension of rationals would correspond to some physically preferred p-adic primes.

It should be possible to tell what the preferred character means in terms of the adelic representations. What happens for these representations of Galois group in this case? This is known.

1. For Galois extensions ramification indices are constant:  $e(i) = e$  and Galois group acts transitively on ideals  $P_i$  dividing  $P$ . One obtains an  $n$ -dimensional representation of Galois group. Same applies to the subgroup of Galois group  $G/I$  where  $I$  is subgroup of  $G$  leaving  $P_i$  invariant. This group is called inertia group. For the maximally ramified case  $G$  maps the ideal  $P_0$  in  $P = P_0^n$  to itself so that  $G = I$  and the action of Galois group is trivial taking  $P_0$  to itself, and one obtains singlet representations.
2. The trivial action of Galois group looks like a technical problem for Langlands program and also for TGD unless the singletness of  $P_i$  under  $G$  has some physical interpretation. One possibility is that Galois group acts as like a gauge group and here the hierarchy of sub-algebras of super-symplectic algebra labelled by integers  $n$  is highly suggestive. This raises obvious questions. Could the integer  $n$  characterizing the sub-algebra of super-symplectic algebra acting as conformal gauge transformations, define the integer defined by the product of ramified primes?  $P_0^n$  brings in mind the  $n$  conformal equivalence classes which remain invariant under the conformal transformations acting as gauge transformations. . Recalling that relative discriminant is an of  $K$  ideal divisible by ramified prime ideals of  $K$ , this means that  $n$  would correspond to the relative discriminant for  $K = Q$ . Are the preferred primes those which are “physical” in the sense that one can assign to the states satisfying conformal gauge conditions?

If the Galois group corresponds to gauge symmetries for these primes, it is physically natural that the p-adic algebraic extension does not exist and that p-adic variant of the Galois group is absent. Nothing is lost from cognition since there is nothing to cognize!

#### 2.5.4 What could be the origin of p-adic length scale hypothesis?

The argument would explain the existence of preferred p-adic primes. It does not yet explain p-adic length scale hypothesis [K21, K13] stating that p-adic primes near powers of 2 are favored. A possible generalization of this hypothesis is that primes near powers of prime are favored. There indeed exists evidence for the realization of 3-adic time scale hierarchies in living matter [?] (see <http://tinyurl.com/jbh9m27>) and in music both 2-adicity and 3-adicity could be present, this is discussed in TGD inspired theory of music harmony and genetic code [K23].

The weak form of NMP might come in rescue here.

1. Entanglement negentropy for a negentropic entanglement [K15] characterized by  $n$ -dimensional projection operator is the  $\log(N_p(n))$  for some  $p$  whose power divides  $n$ . The maximum negentropy is obtained if the power of  $p$  is the largest power of prime divisor of  $p$ , and this can be taken as definition of number theoretic entanglement negentropy. If the largest divisor is  $p^k$ , one has  $N = k \times \log(p)$ . The entanglement negentropy per entangled state is  $N/n = k \log(p)/n$  and is maximal for  $n = p^k$ . Hence powers of prime are favoured which means that p-adic length scale hierarchies with scales coming as powers of  $p$  are negentropically favored and should be generated by NMP. Note that  $n = p^k$  would define a hierarchy of  $h_{eff}/h = p^k$ . During the first years of  $h_{eff}$  hypothesis I believe that the preferred values obey  $h_{eff} = r^k$ ,  $r$  integer not far from  $r = 2^{11}$ . It seems that this belief was not totally wrong.
2. If one accepts this argument, the remaining challenge is to explain why primes near powers of two (or more generally  $p$ ) are favoured.  $n = 2^k$  gives large entanglement negentropy for the final state. Why primes  $p = n_2 = 2^k - r$  would be favored? The reason could be following.

$n = 2^k$  corresponds to  $p = 2$ , which corresponds to the lowest level in p-adic evolution since it is the simplest p-adic topology and farthest from the real topology and therefore gives the poorest cognitive representation of real preferred extremal as p-adic preferred external (Note that  $p = 1$  makes formally sense but for it the topology is discrete).

3. Weak form of NMP [K15, K34] suggests a more convincing explanation. The density matrix of the state to be reduced is a direct sum over contributions proportional to projection operators. Suppose that the projection operator with largest dimension has dimension  $n$ . Strong form of NMP would say that final state is characterized by  $n$ -dimensional projection operator. Weak form of NMP allows free will so that all dimensions  $n - k$ ,  $k = 0, 1, \dots, n - 1$  for final state projection operator are possible. 1-dimensional case corresponds to vanishing entanglement negentropy and ordinary state function reduction isolating the measured system from external world.
4. The negentropy of the final state per state depends on the value of  $k$ . It is maximal if  $n - k$  is power of prime. For  $n = 2^k = M_k + 1$ , where  $M_k$  is Mersenne prime  $n - 1$  gives the maximum negentropy and also maximal p-adic prime available so that this reduction is favoured by NMP. Mersenne primes would be indeed special. Also the primes  $n = 2^k - r$  near  $2^k$  produce large entanglement negentropy and would be favored by NMP.
5. This argument suggests a generalization of p-adic length scale hypothesis so that  $p = 2$  can be replaced by any prime.

This argument together with the hypothesis that preferred prime is ramified would correlate the character of the irreducible extension and character of super-conformal symmetry breaking. The integer  $n$  characterizing super-symplectic conformal sub-algebra acting as gauge algebra would depend on the irreducible algebraic extension of rational involved so that the hierarchy of quantum criticalities would have number theoretical characterization. Ramified primes could appear as divisors of  $n$  and  $n$  would be essentially a characteristic of ramification known as discriminant. An interesting question is whether only the ramified primes allow the continuation of string world sheet and partonic 2-surface to a 4-D space-time surface. If this is the case, the assumptions behind p-adic mass calculations would have full first principle justification.

### 2.5.5 A connection with infinite primes?

Infinite primes are one of the mathematical outcomes of TGD [K26]. There are two kinds of infinite primes. There are the analogs of free many particle states consisting of fermions and bosons labelled by primes of the previous level in the hierarchy. They correspond to states of a supersymmetric arithmetic quantum field theory or actually a hierarchy of them obtained by a repeated second quantization of this theory. A connection between infinite primes representing bound states and irreducible polynomials is highly suggestive.

1. The infinite prime representing free many-particle state decomposes to a sum of infinite part and finite part having no common finite prime divisors so that prime is obtained. The infinite part is obtained from "fermionic vacuum"  $X = \prod_k p_k$  by dividing away some fermionic primes  $p_i$  and adding their product so that one has  $X \rightarrow X/m + m$ , where  $m$  is square free integer. Also  $m = 1$  is allowed and is analogous to fermionic vacuum interpreted as Dirac sea without holes.  $X$  is infinite prime and pure many-fermion state physically. One can add bosons by multiplying  $X$  with any integers having no common denominators with  $m$  and its prime decomposition defines the bosonic contents of the state. One can also multiply  $m$  by any integers whose prime factors are prime factors of  $m$ .
2. There are also infinite primes, which are analogs of bound states and at the lowest level of the hierarchy they correspond to irreducible polynomials  $P(x)$  with integer coefficients. At the second levels the bound states would naturally correspond to irreducible polynomials  $P_n(x)$  with coefficients  $Q_k(y)$ , which are infinite integers at the previous level of the hierarchy.
3. What is remarkable that bound state infinite primes at given level of hierarchy would define maximally ramified algebraic extensions at previous level. One indeed has infinite hierarchy

of infinite primes since the infinite primes at given level are infinite primes in the sense that they are not divisible by the primes of the previous level. The formal construction works as such. Infinite primes correspond to polynomials of single variable at the first level, polynomials of two variables at second level, and so on. Could the Langlands program could be generalized from the extensions of rationals to polynomials of complex argument and that one would obtain infinite hierarchy?

4. Infinite integers in turn could correspond to products of irreducible polynomials defining more general extensions. This raises the conjecture that infinite primes for an extension  $K$  of rationals could code for the algebraic extensions of  $K$  quite generally. If infinite primes correspond to real quantum states they would thus correspond the extensions of rationals to which the parameters appearing in the functions defining partonic 2-surfaces and string world sheets.

This would support the view that partonic 2-surfaces associated with algebraic extensions defined by infinite integers and thus not irreducible are unstable against decay to partonic 2-surfaces which corresponds to extensions assignable to infinite primes. Infinite composite integer defining intermediate unstable state would decay to its composites. Basic particle physics phenomenology would have number theoretic analog and even more.

5. According to Wikipedia, Eisenstein's criterion (<http://tinyurl.com/47kxjz>) allows generalization and what comes in mind is that it applies in exactly the same form also at the higher levels of the hierarchy. Primes would be only replaced with prime polynomials and there would be at least one prime polynomial  $Q(y)$  dividing the coefficients of  $P_n(x)$  except the highest one such that its square would not divide  $P_0$ . Infinite primes would give rise to an infinite hierarchy of functions of many complex variables. At first level zeros of function would give discrete points at partonic 2-surface. At second level one would obtain 2-D surface: partonic 2-surfaces or string world sheet. At the next level one would obtain 4-D surfaces. What about higher levels? Does one obtain higher dimensional objects or something else. The union of  $n$  2-surfaces can be interpreted also as  $2n$ -dimensional surface and one could think that the hierarchy describes a hierarchy of unions of correlated partonic 2-surfaces. The correlation would be due to the preferred extremal property of Kähler action.

One can ask whether this hierarchy could allow to generalize number theoretical Langlands to the case of function fields using the notion of prime function assignable to infinite prime. What this hierarchy of polynomials of arbitrary many complex arguments means physically is unclear. Do these polynomials describe many-particle states consisting of partonic 2-surface such that there is a correlation between them as sub-manifolds of the same space-time sheet representing a preferred extremals of Kähler action?

This would suggest strongly the generalization of the notion of p-adicity so that it applies to infinite primes.

1. This looks sensible and maybe even practical! Infinite primes can be mapped to prime polynomials so that the generalized p-adic numbers would be power series in prime polynomial - Taylor expansion in the coordinate variable defined by the infinite prime. Note that infinite primes (irreducible polynomials) would give rise to a hierarchy of preferred coordinate variables. In terms of infinite primes this expansion would require that coefficients are smaller than the infinite prime  $P$  used. Are the coefficients lower level primes? Or also infinite integers at the same level smaller than the infinite prime in question? This criterion makes sense since one can calculate the ratios of infinite primes as real numbers.
2. I would guess that the definition of infinite-P p-adicity is not a problem since mathematicians have generalized the number theoretical notions to such a level of abstraction much above of a layman like me. The basic question is how to define p-adic norm for the infinite primes (infinite only in real sense, p-adically they have unit norm for all lower level primes) so that it is finite.
3. There exists an extremely general definition of generalized p-adic number fields (see <http://tinyurl.com/y5zreeg>). One considers Dedekind domain  $D$ , which is a generalization

of integers for ordinary number field having the property that ideals factorize uniquely to prime ideals. Now  $D$  would contain infinite integers. One introduces the field  $E$  of fractions consisting of infinite rationals.

Consider element  $e$  of  $E$  and a general fractional ideal  $eD$  as counterpart of ordinary rational and decompose it to a ratio of products of powers of ideals defined by prime ideals, now those defined by infinite primes. The general expression for the p-adic norm of  $x$  is  $x^{-ord(P)}$ , where  $n$  defines the total number of ideals  $P$  appearing in the factorization of a fractional ideal in  $E$ : this number can be also negative for rationals. When the residue field is finite (finite field  $G(p,1)$  for p-adic numbers), one can take  $c$  to the number of its elements ( $c = p$  for p-adic numbers).

Now it seems that this number is not finite since the number of ordinary primes smaller than  $P$  is infinite! But this is not a problem since the topology for completion does not depend on the value of  $c$ . The simple infinite primes at the first level (free many-particle states) can be mapped to ordinary rationals and q-adic norm suggests itself: could it be that infinite-P p-adicity corresponds to q-adicity discussed by Khrennikov [A19]. Note however that q-adic numbers are not a field.

Finally a loosely related question. Could the transition from infinite primes of  $K$  to those of  $L$  takes place just by replacing the finite primes appearing in infinite prime with the decompositions? The resulting entity is infinite prime if the finite and infinite part contain no common prime divisors in  $L$ . This is not the case generally if one can have primes  $P_1$  and  $P_2$  of  $K$  having common divisors as primes of  $L$ : in this case one can include  $P_1$  to the infinite part of infinite prime and  $P_2$  to finite part.

### 3 TGD And Classical Number Fields

This section is devoted to the vision about TGD as a generalized number theory. The basic theme is the role of classical number fields [A10, A4, A13] in quantum TGD. A central notion is  $M^8 - H$  duality which might be also called number theoretic compactification. This duality allows to identify embedding space equivalently either as  $M^8$  or  $M^4 \times CP_2$  and explains the symmetries of standard model number theoretically. These number theoretical symmetries induce also the symmetries dictating the geometry of the “world of classical worlds” (WCW) as a union of symmetric spaces [A16]. This infinite-dimensional Kähler geometry is expected to be highly unique from the mere requirement of its existence requiring infinite-dimensional symmetries provided by the generalized conformal symmetries of the light-cone boundary  $\delta M^4_+ \times S$  and of light-like 3-surfaces and the answer to the question what makes 8-D embedding space and  $S = CP_2$  so unique would be the reduction of these symmetries to number theory.

ZEO has become the corner stone of also number theoretical vision. In ZEO either light-like or space-like 3-surfaces can be identified as the fundamental dynamical objects, and the extension of general coordinate invariance leads to effective 2-dimensionality (strong form of holography) in the sense that the data associated with partonic 2-surfaces and the distribution of 4-D tangent spaces at them located at the light-like boundaries of causal diamonds (CDs) defined as intersections of future and past directed light-cones code for quantum physics and the geometry of WCW. Also the hierarchy of Planck constants [K10] plays a role but not so important one.

The basic number theoretical structures are complex numbers, quaternions [A13] and octonions [A10], and their complexifications obtained by introducing additional commuting imaginary unit  $\sqrt{-1}$ . Hyper-octonionic (-quaternionic,-complex) sub-spaces for which octonionic imaginary units are multiplied by commuting  $\sqrt{-1}$  have naturally Minkowskian signature of metric. The question is whether and how the hyper-structures could allow to understand quantum TGD in terms of classical number fields. The answer which looks the most convincing one relies on the existence of octonionic representation of 8-D gamma matrix algebra.

1. The first guess is that associativity condition for the sub-algebras of the local Clifford algebra defined in this manner could select 4-D surfaces as associative (hyper-quaternionic) sub-spaces of this algebra and define WCW purely number theoretically. The associative sub-spaces in question would be spanned by the modified gamma matrices defined by the Kähler-Dirac

action fixed by the variational principle (Kähler action) selecting space-time surfaces as preferred extremals [K37].

2. This condition is quite not enough: one must strengthen it with the condition that a preferred commutative and thus hyper-complex sub-algebra is contained in the tangent space of the space-time surface. This condition actually generalizes somewhat since one can introduce a family of so called Hamilton-Jacobi coordinates for  $M^4$  allowing an integrable distribution of decompositions of tangent space to the space of non-physical and physical polarizations [K3]. The physical interpretation is as a number theoretic realization of gauge invariance selecting a preferred local commutative plane of non-physical polarizations.
3. Even this is not yet the whole story: one can define also the notions of co-associativity and co-commutativity applying in the regions of space-time surface with Euclidian signature of the induced metric. The basic unproven conjecture is that the decomposition of space-time surfaces to associative and co-associative regions containing preferred commutative *resp.* co-commutative 2-plane in the 4-D tangent plane is equivalent with the preferred extremal property of Kähler action and the hypothesis that space-time surface allows a slicing by string world sheets and by partonic 2-surfaces [K37].

### 3.1 Notations

Some notational conventions are in order before continuing. The fields of quaternions *resp.* octonions having dimension 4 *resp.* 8 and will be denoted by  $Q$  and  $O$ . Their complexified variants will be denoted by  $Q_C$  and  $O_C$ . The sub-spaces of hyper-quaternions  $HQ$  and hyper-octonions  $HO$  are obtained by multiplying the quaternionic and octonionic imaginary units by  $\sqrt{-1}$ . These sub-spaces are very intimately related with the corresponding algebras, and can be seen as Euclidian and Minkowkian variants of the same basic structure. Also the Abelianized versions of the hyper-quaternionic and -octonionic sub-spaces can be considered: these algebras have a representation in the space of spinors of embedding space  $H = M^4 \times CP_2$ .

### 3.2 Quaternion And Octonion Structures And Their Hyper Counterparts

In this introductory section the notions of quaternion and octonion structures and their hyper counterparts are introduced with strong emphasis on the physical interpretation. Literature contains several variants of these structures (Hyper-Kähler structure [A6] and quaternion Kähler structure possessed also by  $CP_2$  [A34]). The notion introduced here is inspired by the physical motivations coming from TGD. As usual the first proposal based on the notions of (hyper-)quaternion and (hyper-)octonion analyticity was not the correct one. Much later a local variant of the notion based on tangent space emerged.

#### 3.2.1 Octonions and quaternions

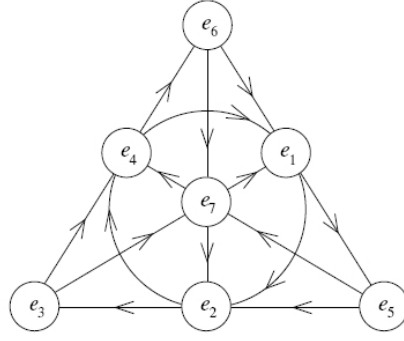
In the following only the basic definitions relating to octonions and quaternions are given (see **Fig. 1**). There is an excellent article by John Baez [A10] describing octonions and their relations to the rest of mathematics and physics.

Octonions can be expressed as real linear combinations  $\sum_k x^k I_k$  of the octonionic real unit  $I_0 = 1$  (counterpart of the unit matrix) and imaginary units  $I_a$ ,  $a = 1, \dots, 7$  satisfying

$$\begin{aligned} I_0^2 &= I_0 \equiv 1 \quad , \\ I_a^2 &= -I_0 = -1 \quad , \\ I_0 I_a &= I_a \quad . \end{aligned} \tag{3.1}$$

Octonions are closed with respect to the ordinary sum of the 8-dimensional vector space and with respect to the octonionic multiplication, which is neither commutative ( $ab \neq ba$  in general) nor associative ( $a(bc) \neq (ab)c$  in general).





**Figure 1:** Octonionic triangle: the six lines and one circle containing three vertices define the seven associative triplets for which the multiplication rules of the ordinary quaternion imaginary units hold true. The arrow defines the orientation for each associative triplet. Note that the product for the units of each associative triplets equals to real unit apart from sign factor.

A concise manner to summarize octonionic multiplication is by using octonionic triangle. Each line (6 altogether) containing 3 octonionic imaginary units forms an associative triple which together with  $I_0 = 1$  generate a division algebra of quaternions. Also the circle spanned by the 3 imaginary units at the middle of the sides of the triangle is associative triple. The multiplication rules for each associative triple are simple:

$$I_a I_b = \epsilon_{abc} I_c, \quad (3.2)$$

where  $\epsilon_{abc}$  is 3-dimensional permutation symbol.  $\epsilon_{abc} = 1$  for the clockwise sequence of vertices (the direction of the arrow along the circumference of the triangle and circle). As a special case this rule gives the multiplication table of quaternions. A crucial observation for what follows is that any pair of imaginary units belongs to one associative triple.

The non-vanishing structure constants  $d_{ab}^c$  of the octonionic algebra can be read directly from the octonionic triangle. For a given pair  $I_a, I_b$  one has

$$\begin{aligned} I_a I_b &= d_{ab}^c I_c, \\ d_{ab}^c &= \epsilon_{ab}^c, \\ I_a^2 &= d_{aa}^0 I_0 = -I_0, \\ I_0^2 &= d_{00}^0 I_0, \\ I_0 I_a &= d_{0a}^a I_a. \end{aligned} \quad (3.3)$$

For  $\epsilon_{abc}$   $c$  belongs to the same associative triple as  $ab$ .

Non-associativity means that is not possible to represent octonions as matrices since matrix product is associative. Quaternions can be represented and the structure constants provide the defining representation as  $I_a \rightarrow d_{abc}$ , where  $b$  and  $c$  are regarded as matrix indices of  $4 \times 4$  matrix. The algebra automorphisms of octonions form 14-dimensional group  $G_2$ , one of the so called exceptional Lie-groups. The isotropy group of imaginary octonion unit is the group  $SU(3)$ . The Euclidian inner product of the two octonions is defined as the real part of the product  $\bar{x}y$

$$\begin{aligned} (x, y) &= Re(\bar{x}y) = \sum_{k=0, \dots, 7} x_k y_k, \\ \bar{x} &= x^0 I_0 - \sum_{i=1, \dots, 7} x^i I_i, \end{aligned} \quad (3.4)$$

and is just the Euclidian norm of the 8-dimensional space.

### 3.2.2 Hyper-octonions and hyper-quaternions

The Euclidicity of the quaternion norm suggests that octonions are not a sensible concept in TGD context. One can imagine two ways to circumvent this conclusion.

1. Minkowskian metric for octonions and quaternions is obtained by identifying Minkowski inner product  $xy$  as the real counterpart of the product

$$x \cdot y \equiv \text{Re}(xy) = x^0 y^0 - \sum_k x^k y^k . \quad (3.5)$$

$SO(1,7)$  ( $SO(1,3)$  in quaternionic case) Lorentz invariance appears completely naturally as the symmetry of the real part of the octonion (quaternion) product and hence of octonions/quaternions and there is no need to perform the complexification of the octonion algebra. Furthermore, only the signature  $(1,7)$  ( $(1,3)$  in the quaternionic case) is possible and this would raise  $M_+^4 \times CP_2$  in a preferred position.

This norm does not give rise to a number theoretic norm defining a homomorphism to real numbers. Indeed, the number theoretic norm defined by the determinant of the linear mapping defined by the multiplication with quaternion or octonion, is inherently Euclidian. This is in conflict with the idea that quaternionic and octonionic primes and their infinite variants should have key role in TGD [K26] .

2. Hyper-octonions and hyper-quaternions provide a possible solution to these problems. These are obtained by multiplying imaginary units by commutative and associative  $\sqrt{-1}$ . These numbers form a sub-space of complexified octonions/quaternions and the cross product of imaginary parts leads out from this sub-space. In this case number theoretic norm induced from  $Q_C/O_C$  gives the fourth/eighth power of Minkowski length and Lorentz group acts as its symmetries. Light-like hyper-quaternions and -octonions causing the failure of the number field property have also a clear physical interpretation.

A criticism against the notion of hyper-quaternionic and octonionic primeness is that the tangent space as an algebra property is lost and the notion of primeness is inherited from  $Q_C/O_C$ . Also non-commutativity and non-associativity could cause difficulties.

ZEO leads to a possible physical interpretation of complexified octonions. The moduli space for causal diamonds corresponds to a Cartesian product of  $M^4 \times CP_2$  whose points label the position of either tip of  $CD \times CP_2$  and space  $I$  whose points label the relative position of the second tip with respect to the first one. p-Adic length scale hypothesis results if one assumes that the proper time distance between the tips comes in powers of two so that one has union of hyperboloids  $H_n \times CP_2$ ,  $H_n = \{m \in M_+^4 | a = 2^n a_0\}$ . A further quantization of hyperboloids  $H_n$  is obtained by replacing it with a lattice like structure is highly suggestive and would correspond to an orbit of a point of  $H_n$  under a subgroup of  $SL(2, Q_C)$  or  $SL(2, Z_C)$  acting as Lorentz transformations in standard manner. Also algebraic extensions of  $Q_C$  and  $Z_C$  can be considered. Also in the case of  $CP_2$  discretization is highly suggestive so that one would have an orbit of a point of  $CP_2$  under a discrete subgroup of  $SU(3, Q)$ .

The outcome could be interpreted by saying that the moduli space in question is  $H \times I$  such that  $H$  corresponds to hyper-octonions and  $I$  to a discretized version of  $\sqrt{-1}H$  and thus a subspace of complexified octonions. An open question whether the quantization has some deeper mathematical meaning.

### 3.2.3 Basic constraints

Before going to details it is useful to make clear the constraints on the concept of the hyper-octonionic structure implied by TGD view about physics.

$M^4 \times CP_2$  cannot certainly be regarded as having any global octonionic structure (for instance in the sense that it could be regarded as a coset space associated with some exceptional group). There are however clear indications for the importance of the hyper-quaternionic and -octonionic structures.

1.  $SU(3)$  is the only simple 8-dimensional Lie-group and acts as the group of isometries of  $CP_2$ : if  $SU(3)$  had some kind of octonionic structure,  $CP_2$  would become unique candidate for the space  $S$ . The decomposition  $SU(3) = h + t$  to  $U(2)$  subalgebra and its complement corresponds rather closely to the decomposition of (hyper-)octonions to (hyper-)quaternionic sub-space and its complement. The electro-weak  $U(2)$  algebra has a natural 1+3 decomposition and generators allow natural hyper-quaternionic structure. The components of the Weyl tensor of  $CP_2$  behave with respect to multiplication like quaternionic imaginary units but only one of them is covariantly constant so that hyper Kähler structure [A6] with three covariantly constant quaternionic imaginary units represented by Kähler forms is not possible. These tensors and metric tensor however define quaternionic structure [A34] .
2.  $M_+^4$  has a natural 1+3 decomposition and a unique cosmic time coordinate defined as the light cone proper time. Hyper-quaternionic structure is consistent with the Minkowskian signature of the inner product and hyper quaternion units have a natural representation in terms of covariantly constant self-dual symplectic forms [A26, A18, A17] and their contractions with sigma matrices. It is not however clear whether this representation is physically interesting.

### 3.2.4 How to define hyper-quaternionic and hyper-octonionic structures?

I have considered several proposals for how to define quaternionic and octonionic structures and their hyper-counterparts.

1. (Hyper-)octonionic manifolds would be obtained by gluing together coordinate patches using (hyper-)octonion analytic functions with real Laurent coefficients (this guarantees associativity and commutativity). This definition does not yet involve metric or any other structures (such as Kähler structure). This approach does not seem to be physically realistic.
2. Second option is based on the idea of representing quaternionic and octonionic imaginary units as antisymmetric tensors. This option makes sense for quaternionic manifolds [A12] and  $CP_2$  indeed represents an example of this kind of manifold. The problem with the octonionic structure is that antisymmetric tensors cannot define non-associative product.
3. If the manifold is endowed with metric, octonionic structure should be defined as a local tangent space structure analogous to eight-bein structure and local gauge algebra structures. This can be achieved by contracting octo-bein vectors with the standard octonionic basis to get octonion form  $I_k$ . Each vector field  $a^k$  defines naturally octonion field  $A = a^k I_k$ . The product of two vector fields can be defined by the octonionic multiplication and this leads to the introduction of a tensor field  $d_{klm}$  of these structure constants obtained as the contraction of the octo-bein vectors with the octonionic structure constants  $d_{abc}$ . Hyper-octonion structure can be defined in a completely analogous manner.

It is possible to induce octonionic structure to any 4-dimensional space-time surface by forming the projection of  $I_k$  to the space-time surface and redefining the products of  $I_k$ 's by dropping away that part of the product, which is orthogonal to the space-time surface. This means that the structure constants of the new 4-dimensional algebra are the projections of  $d_{klm}$  to the space-time surface. One can also define similar induced algebra in the 4-dimensional normal space of the space-time surface. The hypothesis would be that the induced tangential is associative or hyper-quaternionic algebra. Also co-associativity defined as associativity of the normal space algebra is possible. This property would give for the 4-dimensionality of the space-time surface quite special algebraic meaning. The problem is now that there is no direct connection with quantum TGD proper- in particular the connection with the classical dynamics defined by Kähler action is lacking.

4. 8-dimensional gamma matrices allow a representation in terms of tensor products of octonions and  $2 \times 2$  matrices. Genuine matrices are of course not in question since the product of the gamma matrices fails to be associative. An associative representation is obtained by restricting the matrices to a quaternionic plane of complex octonions. If the space-time surface is hyper-quaternionic in the sense that induced gamma matrices define a quaternionic plane of complexified octonions at each point of space-time surface the resulting local Clifford

algebra is associative and structure constants define a matrix representation for the induced gamma matrices.

A more general definition allows gamma matrices to be Kähler-Dirac gamma matrices defined by Kähler action appearing in the Kähler-Dirac action and forced both by internal consistency and super-conformal symmetry [K37]. The Kähler-Dirac gamma matrices associated with Kähler action do not in general define tangent space of the space-time surface as the induced gamma matrices do. Also co-associativity can be considered if one can identify a preferred imaginary unit such that the multiplication of the Kähler-Dirac gamma matrices with this unit gives a quaternionic basis. This condition makes sense only if the preferred extremals of the action are hyper-quaternionic surfaces in the sense defined by the action. That this is true for Kähler action at least is an unproven conjecture.

In the sequel only the fourth option will be considered.

### 3.2.5 How to end up to quantum TGD from number theory?

An interesting possibility is that quantum TGD could emerge from a condition that a local version of hyper-finite factor of type  $II_1$  represented as a local version of infinite-dimensional Clifford algebra exists. The conditions are that “center or mass” degrees of freedom characterizing the position of CD separate uniquely from the “vibrational” degrees of freedom being represented in terms of octonions and that for physical states associativity holds true. The resulting local Clifford algebra would be identifiable as the local Clifford algebra of WCW (being an analog of local gauge groups and conformal fields [A14]).

The uniqueness of  $M^8$  and  $M^4 \times CP_2$  as well as the role of hyper-quaternionic space-time surfaces as fundamental dynamical objects indeed follow from rather weak conditions if one restricts the consideration to gamma matrices and spinors instead of assuming that  $M^8$  coordinates are hyper-octonionic as was done in the first attempts.

1. The unique feature of  $M^8$  and any 8-dimensional space with Minkowski signature of metric is that it is possible to have an octonionic representation of the complexified gamma matrices [K37, K7] and of spinors. This does not require octonionic coordinates for  $M^8$ . The restriction to a quaternionic plane for both gamma matrices and spinors guarantees the associativity.
2. One can also consider a local variant of the octonionic Clifford algebra in  $M^8$ . This algebra contains associative subalgebras for which one can assign to each point of  $M^8$  a hyper-quaternionic plane. It is natural to assume that this plane is either a tangent plane of 4-D manifold defined naturally by the induced gamma matrices defining a basis of tangent space or more generally, by Kähler-Dirac gamma matrices defined by a variational principle (these gamma matrices do not define tangent space in general). Kähler action defines a unique candidate for the variational principle in question. Associativity condition would automatically select sub-algebras associated with 4-D hyper-quaternionic space-time surfaces.
3. This vision bears a very concrete connection to quantum TGD. In [K7] the octonionic formulation of the Kähler-Dirac equation is studied and shown to lead to a highly unique general solution ansatz for the equation working also for the matrix representation of the Clifford algebra. An open question is whether the resulting solution as such defined also solutions of the Kähler-Dirac equation for the matrix representation of gammas. Also a possible identification for 8-dimensional counterparts of twistors as octo-twistors follows: associativity implies that these twistors are very closely related to the ordinary twistors. In TGD framework octo-twistors provide an attractive manner to get rid of the difficulties posed by massive particles for the ordinary twistor formalism.
4. Associativity implies hyperquaternionic space-time surfaces (in a more general sense as usual) and this leads naturally to the notion of WCW and local Clifford algebra in this space. Number theoretic arguments imply  $M^8 - H$  duality. The resulting infinite-dimensional Clifford algebra would differ from von Neumann algebras in that the Clifford algebra and spinors assignable to the center of mass degrees of freedom of causal diamond CD would be expressed in terms of octonionic units although they are associative at space-time surfaces.

One can therefore say that quantum TGD follows by assuming that the tangent space of the embedding space corresponds to a classical number field with maximal dimension.

5. The slicing of the Minkowskian space-time surface inside CD by stringy world sheets and by partonic 2-surfaces inspires the question whether the Kähler-Dirac gamma matrices associated with the stringy world sheets *resp.* partonic 2-surfaces could be commutative *resp.* co-commutative. Commutativity would also be seen as the justification for why the fundamental objects are effectively 2-dimensional.

This formulation is undeniably the most convincing one found hitherto since the notion of hyper-quaternionic structure is local and has elegant formulation in terms of Kähler-Dirac gamma matrices.

### 3.3 Number Theoretic Compactification And $M^8 - H$ Duality

This section summarizes the basic vision about number theoretic compactification reducing the classical dynamics to associativity or co-associativity. Originally  $M^8 - H$  duality was introduced as a number theoretic explanation for  $H = M^4 \times CP_2$ . Much later it turned out that the completely exceptional twistorial properties of  $M^4$  and  $CP_2$  are enough to justify  $X^4 \subset H$  hypothesis. Skeptic could therefore criticize the introduction of  $M^8$  (or even its complexification) as an un-necessary mathematical complication producing only unproven conjectures and bundle of new statements to be formulated precisely.

#### 3.3.1 The basic ideas in nutshell

The vision about the physical role of the classical number fields relies on certain speculative questions and ideas.

1. Could space-time surfaces  $X^4$  be regarded as associative or co-associative (“quaternionic” is equivalent with “associative”) surfaces of  $H$  endowed with octonionic structure in the sense that tangent space of space-time surface would be associative (co-associative) sub-space of octonions at each point of  $X^4$  [K28]. This is certainly possible and an interesting conjecture is that the preferred extremals of Kähler action include associative and co-associative surfaces of  $H$ . Signature of  $M^8$  could be a problem in  $M^8$ :  $M^8$  can be regarded as linear sub-space of complexified octonions and the product of  $M^8$  points does not belong to  $M^8$ . For tangent space this is not the case since one can complexify tangent space.
2. Could the notion of compactification generalize to that of number theoretic compactification in the sense that one can map associative (co-associative) surfaces of  $M^8$  regarded as octonionic linear space to surfaces in  $M^4 \times CP_2$  [K28]? This conjecture -  $M^8 - H$  duality - would give for  $M^4 \times CP_2$  deep number theoretic meaning.  $CP_2$  would parametrize associative planes of octonion space containing fixed complex plane  $M^2 \subset M^8$  and  $CP_2$  point would thus characterize the tangent space of  $X^4 \subset M^8$ . The point of  $M^4$  would be obtained by projecting the point of  $X^4 \subset M^8$  to a point of  $M^4$  identified as tangent space of  $X^4$ . This would guarantee that the dimension of space-time surface in  $H$  would be four. The conjecture is that the preferred extremals of Kähler action include these surfaces.
3.  $M^8 - H$  duality can be generalized to a duality  $H \rightarrow H$  if the images of the associative surface in  $M^8$  is associative surface in  $H$ . One can start from associative surface of  $H$  and assume that it contains the preferred  $M^2$  tangent plane in 8-D tangent space of  $H$  or integrable distribution  $M^2(x)$  of them, and its points to  $H$  by mapping  $M^4$  projection of  $H$  point to itself and associative tangent space to  $CP_2$  point. This point need not be the original one! If the resulting surface is also associative, one can iterate the process indefinitely.
4.  $G_2$  defines the automorphism group of octonions, and one might hope that the maps of octonions to octonions such that the action of Jacobian in the tangent space of associative or co-associative surface reduces to that of  $G_2$  could produce new associative/co-associative surfaces. The action of  $G_2$  would be analogous to that of gauge group.

5. One can also ask whether the notions of commutativity and co-commutativity could have physical meaning. The well-definedness of em charge as quantum number for the modes of the induced spinor field requires their localization to 2-D surfaces (right-handed neutrino is an exception) - string world sheets and partonic 2-surfaces. This can be possible only for Kähler action and could have commutativity and co-commutativity as a number theoretic counterpart. The basic vision would be that the dynamics of Kähler action realizes number theoretical geometrical notions like associativity and commutativity and their co-notions.

One can go even further and ask whether one could somehow construct the preferred extremals of Kähler action using real-octonion analytic functions, call them generically  $f$ . For some time I believed to this idea but it seems I was wrong. The fact that octonion real-analytic functions in  $M^8$  section of  $M_c^8$  have values in the space of complexified octonions makes the complexification of octonions necessary. The simplest guess would be that quaternionic 4-surfaces correspond to the loci at which the values of function  $f$  are real quaternionic. One clearly obtains quaternionic planes as trivial solutions but it is not clear whether their inverse images in general case are quaternionic surfaces and whether non-trivial surfaces with physical properties are obtained. In complex case Riemann zeta serves as a discouraging much simpler analogy since real sub-manifolds of complex plane are just pieces of real axis. Quaternionicity would be replaced with reality and the loci of zeros of the imaginary part of function should be pieces of real axes. Zeta is real at real axis and also at the line  $Im(s) = 1/2$  but the inverse image of this line is not real line. Therefore this approach does not look promising.

### 3.3.2 Is Kähler action needed also at the level of $M^8$

One can question the feasibility of  $M^8 - H$  duality if the dynamics is purely number theoretic at the level of  $M^8$  and determined by Kähler action at the level of  $H$ . Situation becomes more democratic if Kähler action defines the dynamics in both  $M^8$  and  $H$ : this might mean that associativity could imply field equations for preferred extremals or vice versa or there might be equivalence between two. This means the introduction Kähler structure at the level of  $M^8$ , and motivates also the coupling of Kähler gauge potential to  $M^8$  spinors characterized by Kähler charge or em charge. One could call this form of duality strong form of  $M^8 - H$  duality.

The strong form  $M^8 - H$  duality boils down to the assumption that space-time surfaces can be regarded either as 4-surfaces of  $H$  or as surfaces of  $M^8$  composed of associative and co-associative regions identifiable as regions of space-time possessing Minkowskian *resp.* Euclidian signature of the induced metric.

Could they have the same induced metric and Kähler form and WCW associated with  $H$  should be essentially the same as that associated with  $M^8$ . Associativity corresponds to (hyper-)quaternionicity at the level of tangent space and co-associativity to co-(hyper-)quaternionicity - that is associativity/hyper-quaternionicity of the normal space. Both are needed to cope with known extremals. Since in Minkowskian context precise language would force to introduce clumsy terms like hyper-quaternionicity and co-hyper-quaternionicity, it is better to speak just about associativity or co-associativity.

For the octonionic spinor fields the octonionic analogs of electroweak couplings reduce to mere Kähler or electromagnetic coupling and the solutions reduce to those for spinor d'Alembertian in 4-D harmonic potential breaking  $SO(4)$  symmetry. Due to the enhanced symmetry of harmonic oscillator, one expects that partial waves are classified by  $SU(4)$  and by reduction to  $SU(3) \times U(1)$  by em charge and color quantum numbers just as for  $CP_2$  - at least formally.

Harmonic oscillator potential defined by self-dual em field splits  $M^8$  to  $M^4 \times E^4$  and implies Gaussian localization of the spinor modes near origin so that  $E^4$  effectively compactifies. The resulting physics brings strongly in mind low energy physics, where only electromagnetic interaction is visible directly, and one cannot avoid associations with low energy hadron physics. These are some of the reasons for considering  $M^8 - H$  duality as something more than a mere mathematical curiosity.

Kähler form for  $M^8$  non-trivial only in  $E^4 \subset M^8$  implies unique decomposition  $M^8 = M^4 \times E^4$  making possible to identify  $M^4$  point in  $M^8 - H$  duality uniquely. It however turns out that  $M^4$  point corresponds naturally to a projection of  $M^8$  point to the quaternionic tangent space.

### 3.3.3 Definition of complexified octonions and quaternions

The Minkowskian signatures of  $M^8$  and  $M^4$  produce technical nuisance if one tries to define octonion-real-analyticity. One might try to overcome it by Wick rotation, which is however somewhat questionable trick.  $M_c^8 = O_c$  provides another approach giving hopes. Complexified tangent space must be introduced in any case so that its detailed definition deserves to be discussed.

1. The proper formulation for tangent space is in terms of complexified octonions and quaternions involving the introduction of commuting imaginary unit  $j$ . If complexified quaternions are used for  $H$ , Minkowskian signature requires the introduction of two commuting imaginary units  $j$  and  $i$  meaning double complexification.
2. Hyper-quaternions/octonions define as subspace of complexified quaternions/octonions spanned by real unit and  $jI_k$ , where  $I_k$  are quaternionic units. These spaces are obviously not closed under multiplication. One can however define the notion of associativity for the sub-space of  $M^8$  by requiring that the products and sums of the tangent space vectors generate complexified quaternions.
3. Ordinary quaternions  $Q$  are expressible as  $q = q_0 + q^k I_k$ . Hyper-quaternions are expressible as  $q = q_0 + jq^k I_k$  and form a subspace of complexified quaternions  $Q_c = Q \oplus jQ$ . Similar formula applies to octonions and their hyper counterparts which can be regarded as subspaces of complexified octonions  $O \oplus jO$ .
4. One can consider two ways to identify the tangent space of  $H$ . Either as 8-D manifold for which tangent space is hyper-octonionic linear sub-space of complexified octonions  $O_c$  generated by sums and products of tangent vectors. Tangent space vectors of  $H$  could be also identified as hyper-quaternions  $q_H = q_0 + jq^k I_k + jiq_2$  defining a subspace of doubly complexified quaternions: note the appearance of two imaginary units. This would imply an asymmetry between  $M^8$  and  $H$ . The first option looks more elegant also because the composition of the duality maps can be iterated as maps of surfaces of  $H$  to those of  $H$ .

#### 1. Are gamma matrices needed at all?

The recent definitions of associativity and  $M^8 - H$ -duality has evolved slowly from in-accurate characterizations and there are still open questions.

1. The standard spinor structure of  $H$  can be regarded as quaternionic in the sense that gamma matrices are essentially tensor products of quaternionic gamma matrices and reduce in matrix representation for quaternions to ordinary gamma matrices. Therefore the idea that one should introduce octonionic gamma matrices in  $H$  or even  $M^8$  would mean doubling of the spinor structure: not an attractive idea.

It is however important to notice that the introduction of octonionic gamma matrices is not necessary. Simplest option is just the interpretation of tangent basis vectors are octonions: octonion basis is obtained as contractions of vielbein vectors with "flat space" octonions.

2. The earlier formulation was in terms of octonionic flat space gamma matrices replacing the ordinary gamma matrices so that the formulation reduces to that in  $M^8$  tangent space. This formulation is enough to define what associativity means although one can protest.
3. The known extremals provide a test for the associativity (co-associativity) hypothesis. I have not demonstrated that the associativity works for massless extremals (MEs) and vacuum extremals with the dimension of  $CP_2$  projection not larger than 2.
4. Could one define associativity in  $H$  also in terms of modified gamma matrices defined by Kähler action (certainly not  $M^8$ )? The basic problem is that the space spanned by the Kähler-Dirac gamma matrices can have dimension smaller than that of 4 (so that co-basis would have dimension larger than 4 if identified in terms of orthogonal complement). Second problem is that Kähler-Dirac gammas are in general not in the tangent space of space-time surface as vectors of the embedding space. Therefore the notions of associativity (co-associativity) defined in terms of tangent space (normal space) become problematic.

### 3.3.4 Basic formulation of $M^8 - H$ duality

If four-surfaces  $X^4 \subset M^8$  under some conditions define 4-surfaces in  $M^4 \times CP_2$  indirectly, the spontaneous compactification of super string models would correspond in TGD to two different ways to interpret the space-time surface. This correspondence could be called number theoretical compactification or  $M^8 - H$  duality.

### 3.3.5 Basic mathematical facts

The hard mathematical facts behind the notion of number theoretical compactification are following.

1. One manner to define  $M^4$  image of  $M^8$  point uniquely would be to assume that  $M^8$  has unique decomposition  $M^8 = M^4 \times E^4$  (it turns out that this is not the correct manner!). This would be most naturally due to Kähler structure in  $E^4$  defined by a self-dual Kähler form defining parallel constant electric and magnetic fields in Euclidian sense. Besides Kähler form there is vector field coupling to sigma matrix representing the analog of strong isospin: the corresponding octonionic sigma matrix however is imaginary unit times gamma matrix - say  $ie_1$  in  $M^4$  - defining a preferred plane  $M^2$  in  $M^4$ . Here it is essential that the gamma matrices of  $E^4$  defined in terms of octonion units commute to gamma matrices in  $M^4$ . What is involved becomes clear from the Fano triangle illustrating octonionic multiplication table. One can however do also without the introduction of this structure and use only the octonionic structure.
2. The space of hyper-complex structures of the hyper-octonion space - they correspond to the choices of plane  $M^2 \subset M^8$  - is parameterized by 6-sphere  $S^6 = G^2/SU(3)$ . The subgroup  $SU(3)$  of the full automorphism group  $G_2$  respects the a priori selected complex structure and thus leaves invariant one octonionic imaginary unit, call it  $e_1$ . Fixed complex structure therefore corresponds to a point of  $S^6$ .
3. Quaternionic sub-algebras of  $M^8$  are parametrized by  $G_2/U(2)$ . The quaternionic sub-algebras of octonions with fixed complex structure (that is complex sub-space defined by real and preferred imaginary unit and parametrized by a point of  $S^6$ ) are parameterized by  $SU(3)/U(2) = CP_2$  just as the complex planes of quaternion space are parameterized by  $CP_1 = S^2$ . Same applies to hyper-quaternionic sub-spaces of hyper-octonions.  $SU(3)$  would thus have an interpretation as the isometry group of  $CP_2$ , as the automorphism sub-group of octonions, and as color group. Thus the space of quaternionic structures can be parametrized by the 10-dimensional space  $G_2/U(2)$  decomposing as  $S^6 \times CP_2$  locally.
4. The basic result behind number theoretic compactification and  $M^8 - H$  duality is that associative sub-spaces  $M^4 \subset M^8$  containing a fixed commutative sub-space  $M^2 \subset M^8$  are parameterized by  $CP_2$ . The choices of a fixed hyper-quaternionic basis  $1, e_1, e_2, e_3$  with a fixed complex sub-space (choice of  $e_1$ ) are labeled by  $U(2) \subset SU(3)$ . The choice of  $e_2$  and  $e_3$  amounts to fixing  $e_2 \pm \sqrt{-1}e_3$ , which selects the  $U(2) = SU(2) \times U(1)$  subgroup of  $SU(3)$ .  $U(1)$  leaves 1 invariant and induced a phase multiplication of  $e_1$  and  $e_2 \pm e_3$ .  $SU(2)$  induces rotations of the spinor having  $e_2$  and  $e_3$  components. Hence all possible completions of  $1, e_1$  by adding  $e_2, e_3$  doublet are labeled by  $SU(3)/U(2) = CP_2$ .

#### 1. Formulation of $M^8 - H$ duality

Consider now the formulation of  $M^8 - H$  duality.

1. The idea of the standard formulation is that associative manifold  $X^4 \subset M^8$  has at its each point associative tangent plane. That is  $X^4$  corresponds to an integrable distribution of  $M^2(x) \subset M^8$  parametrized 4-D coordinate  $x$  that is map  $x \rightarrow S^6$  such that the 4-D tangent plane is hyper-quaternionic for each  $x$ .
2. One should be able to assign a unique point of  $M^4$  to a given point of  $X^4 \subset M^8$ .



- (a) The associative tangent space of space-time surface shifted to go through the origin of  $M^8$  defines the preferred  $M^4 \subset M^8$  uniquely, and one can project the point of  $M^8$  to this  $M^4$  to get  $M^4$  point. This identification implies that the dimension of tangent space projection to  $M^4$  is maximum, and one avoids the situations in which the image surface of  $H$  has dimension smaller than 4.
- (b) One can imagine also second option which however fails. Since the Kähler structure of  $M^8$  implies a unique decomposition  $M^8 = M^4 \times E^4$ , this surface in turn defines a surface in  $M^4 \times CP_2$  obtained by assigning to the point of 4-surface point  $(m, s) \in H = M^4 \times CP_2$ :  $m \in M^4$  is obtained as *projection*  $M^8 \rightarrow M^4$  (this is modification to the original definition) and  $s \in CP_2$  parametrizes the quaternionic tangent plane as point of  $CP_2$ . Here the local decomposition  $G_2/U(2) = S^6 \times CP_2$  is essential for achieving uniqueness.

One can however represent objection to this identification. The dimension of image in  $H$  is smaller than 4. For instance, hyperquaternionic plane  $M_1^4$  which has  $M^2$  the intersection with preferred  $M^4$  corresponds to constant  $CP_2$  point so that its  $H$  image is  $M^2$ .

### 2. Generalization to $H - H$ duality

As a matter of fact,  $M^8 - H$  duality might generalize to  $H - H$  duality allowing to integrate space-time surfaces and thus WCW to a category.

1. The map of space-time surfaces of  $M^8$  to those of  $H = M^4 \times CP_2$  need not imply that the image surfaces in  $H$  are quaternionic in  $H$ . If they are, then the construction can be iterated. It seems that one can continue this series ad infinitum and could generate new solutions of field equations! If this is the case, one could iterate duality as a sequence  $M^8 \rightarrow H \rightarrow H \dots$  by mapping the space-time surface to  $M^4 \times CP_2$  by the same recipe as in case of  $M^8$ . One would obtain basically a category of space-time surfaces with arrows defined by the duality. Same probably applies to co-associative surfaces. This certainly makes the heart of mathematician beat.
2. It is not proven that associativity/co-associativity implies preferred extremal property for Kähler action. One thing to understand is why Kähler action. An argument in favor of preferred role of Kähler action is that only Kähler action allows localization of spinor modes to 2-D surfaces essential for the well-definedness of em charge [K37]. These surfaces would be string world sheets and possibly also partonic 2-surfaces and their could correspond to commutative and co-commutative 2-surfaces in number theoretic vision and be well-defined also for  $M^8$ . If so, Kähler action would provide a physical representation for the number theoretic notions like associativity and commutativity and their co-notions.
3. If all goes as in dreams, the mere associativity or co-associativity in  $M^8$  would code for the preferred extremal property of Kähler action in  $H$  and would imply this property in  $H$ . The surfaces with this property would form category with arrow defined by the duality.
4. One could also map the associative surface in  $M^8$  to surface in 10-dimensional  $S^6 \times CP_2$ . In this case the metric of the image surface cannot have Minkowskian signature and one cannot assume that the induced metrics are identical. It is not known whether  $S^6$  allows genuine complex structure and Kähler structure which is essential for TGD formulation.

### 3. Some comments

A couple of comments are in order.

1. This definition differs from the first proposal for years ago stating that each point of  $X^4$  contains a *fixed*  $M^2 \subset M^4$  rather than  $M_2(x) \subset M^8$  and also from the proposal assuming integrable distribution of  $M^2(x) \subset M^4$ . The older proposals are not consistent with the properties of massless extremals and string like objects for which the counterpart of  $M^2$  depends on space-time point and is not restricted to  $M^4$ . The earlier definition  $M^2(x) \subset M^4$  was problematic in the co-associative case since for the Euclidian signature it is not clear what the counterpart of  $M^2(x)$  could be.

2. The new definition is consistent with the existence of Hamilton-Jacobi structure meaning slicing of space-time surface by string world sheets and partonic 2-surfaces with points of partonic 2-surfaces labeling the string world sheets [K3]. This structure has been proposed to characterize preferred extremals in Minkowskian space-time regions at least.
3. Co-associative Euclidian 4-surfaces, say  $CP_2$  type vacuum extremal do not contain integrable distribution of  $M^2(x)$ . It is normal space which contains  $M^2(x)$ . Does this have some physical meaning? Or does the surface defined by  $M^2(x)$  have Euclidian analog?

A possible identification of the analog would be as string world sheet at which  $W$  boson field is pure gauge so that the modes of the modified Dirac operator [K37] restricted to the string world sheet have well-defined em charge. This condition appears in the construction of solutions of Kähler-Dirac operator.

For octonionic spinor structure the  $W$  coupling is however absent so that the condition does not make sense in  $M^8$ . The number theoretic condition would be as commutative or co-commutative surface for which imaginary units in tangent space transform to real and imaginary unit by a multiplication with a fixed imaginary unit! One can also formulate co-associativity as a condition that tangent space becomes associative by a multiplication with a fixed imaginary unit.

There is also another justification for the distribution of Euclidian tangent planes. The idea about associativity as a fundamental dynamical principle can be strengthened to the statement that space-time surface allows slicing by hyper-complex or complex 2-surfaces, which are commutative or co-commutative inside space-time surface. The physical interpretation would be as Minkowskian or Euclidian string world sheets carrying spinor modes. This would give a connection with string model and also with the conjecture about the general structure of preferred extremals.

4. Minimalist could argue that the minimal definition requires octonionic structure and associativity *only* in  $M^8$ . There is no need to introduce the counterpart of Kähler action in  $M^8$  since the dynamics would be based on associativity or co-associativity alone. Not that the decomposition  $M^8 = M^4 \times E^4$  is not necessary if  $M^4$  projection is defined to the  $M^4$  defined by hyper-quaternionic tangent place.

### 3.3.6 Hyper-octonionic Pauli “matrices” and the definition of associativity

Octonionic Pauli matrices suggest an interesting possibility to define precisely what associativity means at the level of  $M^8$  using gamma matrices (for background see [L1] ).

1. According to the standard definition space-time surface  $X^4 \subset M^8$  is associative if the tangent space at each point of  $X^4$  in  $X^4 \subset M^8$  picture is associative. The definition can be given also in terms of octonionic gamma matrices whose definition is completely straightforward.
2. Could/should one define the analog of associativity at the level of  $H$ ? One can identify the tangent space of  $H$  as  $M^8$  and can define octonionic structure in the tangent space and this allows to define associativity locally. One can replace gamma matrices with their octonionic variants and formulate associativity in terms of them locally and this should be enough.

Skeptic however reminds  $M^4$  allows hyper-quaternionic structure and  $CP_2$  quaternionic structure so that complexified quaternionic structure would look more natural for  $H$ . The tangent space would decompose as  $M^8 = HQ + ijQ$ , where  $j$  is commuting imaginary unit and  $HQ$  is spanned by real unit and by units  $iI_k$ , where  $i$  second commuting imaginary unit and  $I_k$  denotes quaternionic imaginary units. There is no need to make anything associative.

There is however far from obvious that octonionic spinor structure can be (or need to be!) defined globally. The lift of the  $CP_2$  spinor connection to its octonionic variant has questionable features: in particular vanishing of the charged part and reduction of neutral part to photon. Therefore it is unclear whether associativity condition makes sense for  $X^4 \subset M^4 \times CP_2$ . What makes it so fascinating is that it would allow to iterate duality as a sequences  $M^8 \rightarrow H \rightarrow H \dots$ . This brings in mind the functional composition of octonion real-analytic functions suggested to produce associative or co-associative surfaces.

I have not been able to settle the situation. What seems the working option is associativity in both  $M^8$  and  $H$  and Kähler-Dirac gamma matrices defined by appropriate Kähler action and correlation between associativity and preferred extremal property.

### 3.3.7 Are Kähler and spinor structures necessary in $M^8$ ?

If one introduces  $M^8$  as dual of  $H$ , one cannot avoid the idea that hyper-quaternionic surfaces obtained as images of the preferred extremals of Kähler action in  $H$  are also extremals of  $M^8$  Kähler action with same value of Kähler action defining Kähler function. As found, this leads to the conclusion that the  $M^8 - H$  duality is Kähler isometry. Coupling of spinors to Kähler potential is the next step and this in turn leads to the introduction of spinor structure so that quantum TGD in  $H$  should have full  $M^8$  dual.

#### 1. Are also the 4-surfaces in $M^8$ preferred extremals of Kähler action?

It would be a mathematical miracle if associative and co-associative surfaces in  $M^8$  would be in 1-1 correspondence with preferred extremals of Kähler action. This motivates the question whether Kähler action make sense also in  $M^8$ . This does not exclude the possibility that associativity implies or is equivalent with the preferred extremal property.

One expects a close correspondence between preferred extremals: also now vacuum degeneracy is obtained, one obtains massless extremals, string like objects, and counterparts of  $CP_2$  type vacuum extremals. All known extremals would be associative or co-associative if modified gamma matrices define the notion (possible only in the case of  $H$ ).

The strongest form of duality would be that the space-time surfaces in  $M^8$  and  $H$  have same induced metric same induced Kähler form. The basic difference would be that the spinor connection for surfaces in  $M^8$  would be however neutral and have no left handed components and only em gauge potential. A possible interpretation is that  $M^8$  picture defines a theory in the phase in which electroweak symmetry breaking has happened and only photon belongs to the spectrum.

The question is whether one can define WCW also for  $M^8$ . Certainly it should be equivalent with WCW for  $H$ : otherwise an inflation of poorly defined notions follows. Certainly the general formulation of the WCW geometry generalizes from  $H$  to  $M^8$ . Since the matrix elements of symplectic super-Hamiltonians defining WCW gamma matrices are well defined as matrix elements involve spinor modes with Gaussian harmonic oscillator behavior, the non-compactness of  $E^4$  does not pose any technical problems.

#### 2. Spinor connection of $M^8$

There are strong physical constraints on  $M^8$  dual and they could kill the hypothesis. The basic constraint to the spinor structure of  $M^8$  is that it reproduces basic facts about electro-weak interactions. This includes neutral electro-weak couplings to quarks and leptons identified as different  $H$ -chiralities and parity breaking.

1. By the flatness of the metric of  $E^4$  its spinor connection is trivial.  $E^4$  however allows full  $S^2$  of covariantly constant Kähler forms so that one can accommodate free independent Abelian gauge fields assuming that the independent gauge fields are orthogonal to each other when interpreted as realizations of quaternionic imaginary units. This is possible but perhaps a more natural option is the introduction of just single Kähler form as in the case of  $CP_2$ .
2. One should be able to distinguish between quarks and leptons also in  $M^8$ , which suggests that one introduce spinor structure and Kähler structure in  $E^4$ . The Kähler structure of  $E^4$  is unique apart from  $SO(3)$  rotation since all three quaternionic imaginary units and the unit vectors formed from them allow a representation as an antisymmetric tensor. Hence one must select one preferred Kähler structure, that is fix a point of  $S^2$  representing the selected imaginary unit. It is natural to assume different couplings of the Kähler gauge potential to spinor chiralities representing quarks and leptons: these couplings can be assumed to be same as in case of  $H$ .
3. Electro-weak gauge potential has vectorial and axial parts. Em part is vectorial involving coupling to Kähler form and  $Z^0$  contains both axial and vector parts. The naïve replacement of sigma matrices appearing in the coupling of electroweak gauge fields takes the left handed

parts of these fields to zero so that only neutral part remains. Further, gauge fields correspond to curvature of  $CP_2$  which vanishes for  $E^4$  so that only Kähler form remains. Kähler form couples to  $3L$  and  $q$  so that the basic asymmetry between leptons and quarks remains. The resulting field could be seen as analog of photon.

4. The absence of weak parts of classical electro-weak gauge fields would conform with the standard thinking that classical weak fields are not important in long scales. A further prediction is that this distinction becomes visible only in situations, where  $H$  picture is necessary. This is the case at high energies, where the description of quarks in terms of  $SU(3)$  color is convenient whereas  $SO(4)$  QCD would require large number of  $E^4$  partial waves. At low energies large number of  $SU(3)$  color partial waves are needed and the convenient description would be in terms of  $SO(4)$  QCD. Proton spin crisis might relate to this.

### 3. Dirac equation for leptons and quarks in $M^8$

Kähler gauge potential would also couple to octonionic spinors and explain the distinction between quarks and leptons.

1. The complexified octonions representing  $H$  spinors decompose to  $1 + 1 + 3 + \bar{3}$  under  $SU(3)$  representing color automorphisms but the interpretation in terms of QCD color does not make sense. Rather, the triplet and single combine to two weak isospin doublets and quarks and leptons corresponds to “spin” states of octonion valued 2-spinor. The conservation of quark and lepton numbers follows from the absence of coupling between these states.
2. One could modify the coupling so that coupling is on electric charge by coupling it to electromagnetic charge which as a combination of unit matrix and sigma matrix is proportional to  $1 + kI_1$ , where  $I_1$  is octonionic imaginary unit in  $M^2 \subset M^4$ . The complexified octonionic units can be chosen to be eigenstates of  $Q_{em}$  so that Laplace equation reduces to ordinary scalar Laplacian with coupling to self-dual em field.
3. One expects harmonic oscillator like behavior for the modes of the Dirac operator of  $M^8$  since the gauge potential is linear in  $E^4$  coordinates. One possibility is Cartesian coordinates is  $A(A_x, A_y, A_z, A_t) = k(-y, x, t, -z)$ . The coupling would make  $E^4$  effectively a compact space.
4. The square of Dirac operator gives potential term proportional to  $r^2 = x^2 + y^2 + z^2 + t^2$  so that the spectrum of 4-D harmonic oscillator operator and  $SO(4)$  harmonics localized near origin are expected. For harmonic oscillator the symmetry enhances to  $SU(4)$ .

If one replaces Kähler coupling with em charge symmetry breaking of  $SO(4)$  to vectorial  $SO(3)$  is expected since the coupling is proportional to  $1 + ike_1$  defining electromagnetic charge. Since the basis of complexified quaternions can be chosen to be eigenstates of  $e_1$  under multiplication, octonionic spinors are eigenstates of em charge and one obtains two color singlets  $1 \pm e_1$  and color triplet and antitriplet. The color triplets cannot be however interpreted in terms of quark color.

Harmonic oscillator potential is expected to enhance  $SO(3)$  to  $SU(3)$ . This suggests the reduction of the symmetry to  $SU(3) \times U(1)$  corresponding to color symmetry and em charge so that one would have same basic quantum numbers as tof  $CP_2$  harmonics. An interesting question is how the spectrum and mass squared eigenvalues of harmonics differ from those for  $CP_2$ .

5. In the square of Dirac equation  $J^{kl}\Sigma_{kl}$  term distinguishes between different em charges ( $\Sigma_{kl}$  reduces by self duality and by special properties of octonionic sigma matrices to a term proportional to  $iI_1$  and complexified octonionic units can be chosen to be its eigenstates with eigen value  $\pm 1$ . The vacuum mass squared analogous to the vacuum energy of harmonic oscillator is also present and this contribution are expected to cancel themselves for neutrinos so that they are massless whereas charged leptons and quarks are massive. It remains to be checked that quarks and leptons can be classified to triality  $T = \pm 1$  and  $t = 0$  representations of dynamical  $SU(3)$  respectively.

#### 4. What about the analog of Kähler-Dirac equation

Only the octonionic structure in  $T(M^8)$  is needed to formulate quaternionicity of space-time surfaces: the reduction to  $O_c$ -real-analyticity would be extremely nice but not necessary ( $O_c$  denotes complexified octonions needed to cope with Minkowskian signature). Most importantly, there might be no need to introduce Kähler action (and Kähler form) in  $M^8$ . Even the octonionic representation of gamma matrices is un-necessary. Neither there is any absolute need to define octonionic Dirac equation and octonionic Kähler Dirac equation nor octonionic analog of its solutions nor the octonionic variants of embedding space harmonics.

It would be of course nice if the general formulas for solutions of the Kähler Dirac equation in  $H$  could have counterparts for octonionic spinors satisfying quaternionicity condition. One can indeed wonder whether the restriction of the modes of induced spinor field to string world sheets defined by integrable distributions of hyper-complex spaces  $M^2(x)$  could be interpreted in terms of commutativity of fermionic physics in  $M^8$ .  $M^8 - H$  correspondence could map the octonionic spinor fields at string world sheets to their quaternionic counterparts in  $H$ . The fact that only holomorphy is involved with the definition of modes could make this map possible.

#### 3.3.8 How could one solve associativity/co-associativity conditions?

The natural question is whether and how one could solve the associativity/-co-associativity conditions explicitly. One can imagine two approaches besides  $M^8 \rightarrow H \rightarrow H\dots$  iteration generating new solutions from existing ones.

##### 1. Could octonion-real analyticity be equivalent with associativity/co-associativity?

Analytic functions provide solutions to 2-D Laplace equations and one might hope that also the field equations could be solved in terms of octonion-real-analyticity at the level of  $M^8$  perhaps also at the level of  $H$ . Signature however causes problems - at least technical. Also the compactness of  $CP_2$  causes technical difficulties but they need not be insurmountable.

For  $E^8$  the tangent space would be genuinely octonionic and one can define the notion octonion-real analytic map as a generalization of real-analytic function of complex variables (the coefficients of Laurent series are real to guarantee associativity of the series). The argument is complexified octonion in  $O \oplus iO$  forming an algebra but not a field. The norm square is Minkowskian as difference of two Euclidian octonionic norms:  $N(o_1 + io_2) = N(o_1) - N(o_2)$  and vanishes at 15-D light cone boundary. Obviously, differential calculus is possible outside the light-cone boundary. Rational analytic functions have however poles at the light-cone boundary. One can wonder whether the poles at  $M^4$  light-cone boundary, which is subset of 15-D light-cone boundary could have physical significance and relevant for the role of causal diamonds in ZEO.

The candidates for associative surfaces defined by  $O_c$ -real-analytic functions (I use  $O_c$  for complexified octonions) have Minkowskian signature of metric and are 4-surfaces at which the projection of  $f(o_1 + io_2)$  to  $Im(O_1)$ ,  $iIm(O_2)$ , and  $iRe(Q_2) \oplus Im(Q_1)$  vanish so that only the projection to hyper-quaternionic Minkowskian sub-space  $M^4 = Re(Q_1) + iIm(Q_2)$  with signature  $(1, -1, -, 1-)$  is non-vanishing. The inverse image need not belong to  $M^8$  and in general it belongs to  $M_c^8$  but this is not a problem: all that is needed that the tangent space of inverse image is complexified quaternionic. If this is the case then  $M^8 - H$  duality maps the tangent space of the inverse image to  $CP_2$  point and image itself defines the point of  $M^4$  so that a point of  $H$  is obtained. Co-associative surfaces would be surfaces for which the projections of image to  $Re(O_1)$ ,  $iRe(O_2)$ , and to  $Im(O_1)$  vanish so that only the projection to  $iIm(O_2)$  with signature  $(-1, -1, -1, -1)$  is non-vanishing.

The inverse images as 4-D sub-manifolds of  $M_c^8$  (not  $M^8$ !) are excellent candidates for associative and co-associative 4-surfaces since  $M^8 - H$  duality assigns to them a 4-surface in  $M^4 \times CP_2$  if the tangent space at given point is complexified quaternionic. This is true if one believes on the analytic continuation of the intuition from complex analysis (the image of real axes under the map defined by  $O_c$ -real-analytic function is real axes in the new coordinates defined by the map: the intuition results by replacing “real” by “complexified quaternionic”). The possibility to solve field equations in this manner would be of enormous significance since besides basic arithmetic operations also the functional decomposition of  $O_c$ -real-analytic functions produces similar functions. One could speak of the algebra of space-time surfaces.

What is remarkable that the complexified octonion real analytic functions are obtained by analytic continuation from single real valued function of real argument. The real functions form naturally a hierarchy of polynomials (maybe also rational functions) and number theoretic vision suggests that there coefficients are rationals or algebraic numbers. Already for rational coefficients hierarchy of algebraic extensions of rationals results as one solves the vanishing conditions. There is a temptation to regard this hierarchy coding for space-time sheets as an analog of DNA.

Note that in the recent formulation there is no need to pose separately the condition about integrable distribution of  $M^2(x) \subset M^4$ .

### 2. Quaternionicity condition for space-time surfaces

Quaternionicity actually has a surprisingly simple formulation at the level of space-time surfaces. The following discussion applies to both  $M^8$  and  $H$  with minor modifications if one accepts that also  $H$  can allow octonionic tangent space structure, which does not require gamma matrices.

1. Quaternionicity is equivalent with associativity guaranteed by the vanishing of the associator  $A(a, b, c) = a(bc) - (ab)c$  for any triplet of imaginary tangent vectors in the tangent space of the space-time surface. The condition must hold true for purely imaginary combinations of tangent vectors.
2. If one is able to choose the coordinates in such a way that one of the tangent vectors corresponds to real unit (in the embedding map embedding space  $M^4$  coordinate depends only on the time coordinate of space-time surface), the condition reduces to the vanishing of the octonionic product of remaining three induced gamma matrices interpreted as octonionic gamma matrices. This condition looks very simple - perhaps too simple!- since it involves only first derivatives of the embedding space vectors.

One can of course whether quaternionicity conditions replace field equations or only select preferred extremals. In the latter case, one should be able to prove that quaternionicity conditions are consistent with the field equations.

3. Field equations would reduce to tri-linear equations in in the gradients of embedding space coordinates (rather than involving embedding space coordinates quadratically). Sum of analogs of  $3 \times 3$  determinants deriving from  $a \times (b \times b)$  for different octonion units is involved.
4. Written explicitly field equations give in terms of vielbein projections  $e_\alpha^A$ , vielbein vectors  $e_k^A$ , coordinate gradients  $\partial_\alpha h^k$  and octonionic structure constants  $f_{ABC}$  the following conditions stating that the projections of the octonionic associator tensor to the space-time surface vanishes:

$$\begin{aligned}
 e_\alpha^A e_\beta^B e_\gamma^C A_{ABC}^E &= 0 , \\
 A_{ABC}^E &= f_{AD}^E f_{BC}^D - f_{AB}^D f_{DC}^E , \\
 e_\alpha^A &= \partial_\alpha h^k e_k^A , \\
 \Gamma_k &= e_k^A \gamma_A .
 \end{aligned}
 \tag{3.6}$$

The very naïve idea would be that the field equations are indeed integrable in the sense that they reduce to these tri-linear equations. Tri-linearity in derivatives is highly non-trivial outcome simplifying the situation further. These equations can be formulated as the as purely algebraic equations written above plus integrability conditions

$$F_{\alpha\beta}^A = D_\alpha e_\beta^A - D_\beta e_\alpha^A = 0 . \tag{3.7}$$

One could say that vielbein projections define an analog of a trivial gauge potential. Note however that the covariant derivative is defined by spinor connection rather than this effective gauge potential which reduces to that in  $SU(2)$ . Similar formulation holds true for field equations and one should be able to see whether the field equations formulated in terms of derivatives of vielbein projections commute with the associativity conditions.

5. The quaternionicity conditions can be formulated as vanishing of generalization of Cayley's hyperdeterminant for "hypermatrix"  $a_{ijk}$  with 2-valued indices (see <http://tinyurl.com/ya7h3n9z>). Now one has 8 hyper-matrices with 3 8-valued indices associated with the vanishing  $A_{BCD}^E x^B y^C z^D = 0$  of trilinear forms defined by the associators. The conditions say something only about the octonion structure constants and since octonionic space allow quaternionic sub-spaces these conditions must be satisfied.

The inspection of the Fano triangle [A36] (see **Fig. 1**) expressing the multiplication table for octonionic imaginary units reveals that give any two imaginary octonion units  $e_1$  and  $e_2$  their product  $e_1 e_2$  (or equivalently commutator) is imaginary octonion unit (2 times octonion unit) and the three units span together with real unit quaternionic sub-algebra. There it seems that one can generate local quaternionic sub-space from two imaginary units plus real unit. This generalizes to the vielbein components of tangent vectors of space-time surface and one can build the solutions to the quaternionicity conditions from vielbein projections  $e_1, e_2$ , their product  $e_3 = k(x)e_1 e_2$  and real fourth "time-like" vielbein component which must be expressible as a combination of real unit and imaginary units:

$$e_0 = a \times 1 + b^i e_i$$

For static solutions this condition is trivial. Here summation over  $i$  is understood in the latter term. Besides these conditions one has integrability conditions and field equations for Kähler action. This formulation suggests that quaternionicity is additional - perhaps defining - property of preferred extremals.

### 3.3.9 Quaternionicity at the level of embedding space quantum numbers

From the multiplication table of octonions as illustrated by Fano triangle [A36] one finds that all edges of the triangle, the middle circle and the three the lines connecting vertices to the midpoints of opposite side define triplets of quaternionic units. This means that by taking real unit and any imaginary unit in quaternionic  $M^4$  algebra spanning  $M^2 \subset M^4$  and two imaginary units in the complement representing  $CP_2$  tangent space one obtains quaternionic algebra. This suggests an explanation for the preferred  $M^2$  contained in tangent space of space-time surface (the  $M^2$ 's could form an integrable distribution). Four-momentum restricted to  $M^2$  and  $I_3$  and  $Y$  interpreted as tangent vectors in  $CP_2$  tangent space defined quaternionic sub-algebra. This could give content for the idea that quantum numbers are quaternionic.

I have indeed proposed that the four-momentum belongs to  $M^2$ . If  $M^2(x)$  form a distribution as the proposal for the preferred extremals suggests this could reflect momentum exchanges between different points of the space-time surface such that total momentum is conserved or momentum exchange between two sheets connected by wormhole contacts.

### 3.3.10 Questions

In following some questions related to  $M^8 - H$  duality are represented.

1. Could associativity condition be formulated using modified gamma matrices?

Skeptic can criticize the minimal form of  $M^8 - H$  duality involving no Kähler action in  $M^8$  is unrealistic. Why just Kähler action? What makes it so special? The only defense that I can imagine is that Kähler action is in many respects unique choice.

An alternative approach would replace induced gamma matrices with the modified ones to get the correlation In the case of  $M^8$  this option cannot work. One cannot exclude it for  $H$ .

1. For Kähler action the Kähler-Dirac gamma matrices  $\Gamma^\alpha = \frac{\partial L_K}{\partial h_\alpha^k} \Gamma^k$ ,  $\Gamma_k = e_k^A \gamma_A$ , assign to a given point of  $X^4$  a 4-D space which need not be tangent space anymore or even its sub-space.

The reason is that canonical momentum current contains besides the gravitational contribution coming from the induced metric also the "Maxwell contribution" from the induced Kähler form not parallel to space-time surface. In the case of  $M^8$  the duality map to  $H$  is therefore lost.

2. The space spanned by the Kähler-Dirac gamma matrices need not be 4-dimensional. For vacuum extremals with at most 2-D  $CP_2$  projection Kähler-Dirac gamma matrices vanish identically. For massless extremals they span 1- D light-like subspace. For  $CP_2$  vacuum extremals the modified gamma matrices reduces to ordinary gamma matrices for  $CP_2$  and the situation reduces to the quaternionicity of  $CP_2$ . Also for string like objects the conditions are satisfied since the gamma matrices define associative sub-space as tangent space of  $M^2 \times S^2 \subset M^4 \times CP_2$ . It seems that associativity is satisfied by all known extremals. Hence Kähler-Dirac gamma matrices are flexible enough to realize associativity in  $H$ .
3. Kähler-Dirac gamma matrices in Dirac equation are required by super conformal symmetry for the extremals of action and they also guarantee that vacuum extremals defined by surfaces in  $M^4 \times Y^2$ ,  $Y^2$  a Lagrange sub-manifold of  $CP_2$ , are trivially hyper-quaternionic surfaces. The modified definition of associativity in  $H$  does not affect in any manner  $M^8 - H$  duality necessarily based on induced gamma matrices in  $M^8$  allowing purely number theoretic interpretation of standard model symmetries. One can however argue that the most natural definition of associativity is in terms of induced gamma matrices in both  $M^8$  and  $H$ .

**Remark:** A side comment not strictly related to associativity is in order. The anti-commutators of the Kähler-Dirac gamma matrices define an effective Riemann metric and one can assign to it the counterparts of Riemann connection, curvature tensor, geodesic line, volume, etc... One would have two different metrics associated with the space-time surface. Only if the action defining space-time surface is identified as the volume in the ordinary metric, these metrics are equivalent. The index raising for the effective metric could be defined also by the induced metric and it is not clear whether one can define Riemann connection also in this case. Could this effective metric have concrete physical significance and play a deeper role in quantum TGD? For instance, AdS-CFT duality leads to ask whether interactions be coded in terms of the gravitation associated with the effective metric.

Now skeptic can ask why should one demand  $M^8 - H$  correspondence if one in any case is forced to introduced Kähler also at the level of  $M^8$ ? Does  $M^8 - H$  correspondence help to construct preferred extremals or does it only bring in a long list of conjectures? I can repeat the questions of the skeptic.

### 2. Minkowskian-Euclidian $\leftrightarrow$ associative-co-associative?

The 8-dimensionality of  $M^8$  allows to consider both associativity of the tangent space and associativity of the normal space- let us call this co-associativity of tangent space- as alternative options. Both options are needed as has been already found. Since space-time surface decomposes into regions whose induced metric possesses either Minkowskian or Euclidian signature, there is a strong temptation to propose that Minkowskian regions correspond to associative and Euclidian regions to co-associative regions so that space-time itself would provide both the description and its dual.

The proposed interpretation of conjectured associative-co-associative duality relates in an interesting manner to p-adic length scale hypothesis selecting the primes  $p \simeq 2^k$ ,  $k$  positive integer as preferred p-adic length scales.  $L_p \propto \sqrt{p}$  corresponds to the p-adic length scale defining the size of the space-time sheet at which elementary particle represented as  $CP_2$  type extremal is topologically condensed and is of order Compton length.  $L_k \propto \sqrt{k}$  represents the p-adic length scale of the wormhole contacts associated with the  $CP_2$  type extremal and  $CP_2$  size is the natural length unit now. Obviously the quantitative formulation for associative-co-associative duality would be in terms  $p \rightarrow k$  duality.

### 3. Can $M^8 - H$ duality be useful?

Skeptic could of course argue that  $M^8 - H$  duality generates only an inflation of unproven conjectures. This might be the case. In the following I will however try to defend the conjecture. One can however find good motivations for  $M^8 - H$  duality: both theoretical and physical.

1. If  $M^8 - H$  duality makes sense for induced gamma matrices also in  $H$ , one obtains infinite sequence of dualities allowing to construct preferred extremals iteratively. This might relate to octonionic real-analyticity and composition of octonion-real-analytic functions.



2.  $M^8 - H$  duality could provide much simpler description of preferred extremals of Kähler action as hyper-quaternionic surfaces. Unfortunately, it is not clear whether one should introduce the counterpart of Kähler action in  $M^8$  and the coupling of  $M^8$  spinors to Kähler form. Note that the Kähler form in  $E^4$  would be self dual and have constant components: essentially parallel electric and magnetic field of same constant magnitude.
3.  $M^8 - H$  duality provides insights to low energy physics, in particular low energy hadron physics.  $M^8$  description might work when  $H$ -description fails. For instance, perturbative QCD which corresponds to  $H$ -description fails at low energies whereas  $M^8$  description might become perturbative description at this limit. Strong  $SO(4) = SU(2)_L \times SU(2)_R$  invariance is the basic symmetry of the phenomenological low energy hadron models based on conserved vector current hypothesis (CVC) and partially conserved axial current hypothesis (PCAC). Strong  $SO(4) = SU(2)_L \times SU(2)_R$  relates closely also to electro-weak gauge group  $SU(2)_L \times U(1)$  and this connection is not well understood in QCD description.  $M^8 - H$  duality could provide this connection. Strong  $SO(4)$  symmetry would emerge as a low energy dual of the color symmetry. Orbital  $SO(4)$  would correspond to strong  $SU(2)_L \times SU(2)_R$  and by flatness of  $E^4$  spin like  $SO(4)$  would correspond to electro-weak group  $SU(2)_L \times U(1)_R \subset SO(4)$ . Note that the inclusion of coupling to Kähler gauge potential is necessary to achieve respectable spinor structure in  $CP_2$ . One could say that the orbital angular momentum in  $SO(4)$  corresponds to strong isospin and spin part of angular momentum to the weak isospin.

This argument does not seem to be consistent with  $SU(3) \times U(1) \subset SU(4)$  symmetry for  $Mx$  Dirac equation. One can however argue that  $SU(4)$  symmetry combines  $SO(4)$  multiplets together. Furthermore,  $SO(4)$  represents the isometries leaving Kähler form invariant.

#### 4. $M^8 - H$ duality in low energy physics and low energy hadron physics

$M^8 - H$  can be applied to gain a view about color confinement. The basic idea would be that  $SO(4)$  and  $SU(3)$  provide provide dual descriptions of quarks using  $E^4$  and  $CP_2$  partial waves and low energy hadron physics corresponds to a situation in which  $M^8$  picture provides the perturbative approach whereas  $H$  picture works at high energies.

A possible interpretation is that the space-time surfaces vary so slowly in  $CP_2$  degrees of freedom that can approximate  $CP_2$  with a small region of its tangent space  $E^4$ . One could also say that color interactions mask completely electroweak interactions so that the spinor connection of  $CP_2$  can be neglected and one has effectively  $E^4$ . The basic prediction is that  $SO(4)$  should appear as dynamical symmetry group of low energy hadron physics and this is indeed the case.

Consider color confinement at the long length scale limit in terms of  $M^8 - H$  duality.

1. At high energy limit only lowest color triplet color partial waves for quarks dominate so that QCD description becomes appropriate whereas very higher color partial waves for quarks and gluons are expected to appear at the confinement limit. Since WCW degrees of freedom begin to dominate, color confinement limit transcends the descriptive power of QCD.
2. The success of  $SO(4)$  sigma model in the description of low lying hadrons would directly relate to the fact that this group labels also the  $E^4$  Hamiltonians in  $M^8$  picture. Strong  $SO(4)$  quantum numbers can be identified as orbital counterparts of right and left handed electro-weak isospin coinciding with strong isospin for lowest quarks. In sigma model pion and sigma boson form the components of  $E^4$  valued vector field or equivalently collection of four  $E^4$  Hamiltonians corresponding to spherical  $E^4$  coordinates. Pion corresponds to  $S^3$  valued unit vector field with charge states of pion identifiable as three Hamiltonians defined by the coordinate components. Sigma is mapped to the Hamiltonian defined by the  $E^4$  radial coordinate. Excited mesons corresponding to more complex Hamiltonians are predicted.
3. The generalization of sigma model would assign to quarks  $E^4$  partial waves belonging to the representations of  $SO(4)$ . The model would involve also 6  $SO(4)$  gluons and their  $SO(4)$  partial waves. At the low energy limit only lowest representations would be important whereas at higher energies higher partial waves would be excited and the description based on  $CP_2$  partial waves would become more appropriate.

4. The low energy quark model would rely on quarks moving  $SO(4)$  color partial waves. Left *resp.* right handed quarks could correspond to  $SU(2)_L$  *resp.*  $SU(2)_R$  triplets so that spin statistics problem would be solved in the same manner as in the standard quark model.
5. Family replication phenomenon is described in TGD framework the same manner in both cases so that quantum numbers like strangeness and charm are not fundamental. Indeed, p-adic mass calculations allowing fractally scaled up versions of various quarks allow to replace Gell-Mann mass formula with highly successful predictions for hadron masses [K19] .

To my opinion these observations are intriguing enough to motivate a concrete attempt to construct low energy hadron physics in terms of  $SO(4)$  gauge theory.

### 3.3.11 Summary

The overall conclusion is that the most convincing scenario relies on the associativity/co-associativity of space-time surfaces define by induced gamma matrices and applying both for  $M^8$  and  $H$ . The fact that the duality can be continued to an iterated sequence of duality maps  $M^8 \rightarrow H \rightarrow H\dots$  is what makes the proposal so fascinating and suggests connection with fractality.

The introduction of Kähler action and coupling of spinors to Kähler gauge potentials is highly natural. One can also consider the idea that the space-time surfaces in  $M^8$  and  $H$  have same induced metric and Kähler form: for iterated duality map this would mean that the steps in the map produce space-time surfaces which identical metric and Kähler form so that the sequence might stop.  $M^8_H$  duality might provide two descriptions of same underlying dynamics:  $M^8$  description would apply in long length scales and  $H$  description in short length scales.

## 4 Infinite Primes

The notion of prime seems to capture something very essential about what it is to be elementary building block of matter and has become a fundamental conceptual element of TGD. The notion of prime gains its generality from its reducibility to the notion of prime ideal of an algebra. Thus the notion of prime is well-defined, not only in case of quaternions and octonions, but also for their complexifications and one can speak about infinite primes in the case of hyper-quaternions and -octonions, which are especially natural physically and for which numbers having zero norm correspond physically to light-like 8-vectors.

### 4.1 Basic Ideas

#### 4.1.1 The notion of infinite prime

The original motivation for the notion of infinite prime came from the first attempts to construct TGD inspired theory of consciousness (around 1995) [K30] . Suppose very naïvely that the 4-surfaces in a given sector of the “world of classical worlds” (WCW) are labelled by a fixed p-adic prime. The natural expectation is that evolution by quantum jumps means dispersion in the space of these sectors and leads to the increase of the p-adic prime characterizing the Universe. As one moves backwards in subjective time (sequence of quantum jumps) one ends up to the situation in which the prime characterizing the universe was  $p = 2$ . Should one assume that there was the first quantum jump when everything began? If not, then it would seem that the p-adic prime characterizing the Universe must be infinite. Second problem is that the p-adic length scales are finite and if the size scale of Universe is given by p-adic length scale the Universe has finite sized: this does not make sense in TGD framework. The only way out of the problems is the assumption that the p-adic prime characterizing the entire Universe is literally infinite and that p-adic primes characterizing space-time sheets are finite.

These argument, which are by no means central for the recent view about p-adic primes, motivated the attempt to construct a theory of infinite primes and to extend quantum TGD accordingly. This turns out to be possible. The recipe for constructing infinite primes is structurally equivalent with a repeated second quantization of an arithmetic super-symmetric quantum field theory. At the lowest level one has fermionic and bosonic states labeled by finite primes and infinite primes correspond to many particle states of this theory. Also infinite primes analogous to bound

states are predicted. This hierarchy of quantizations can be continued indefinitely by taking the many particle states of the previous level as elementary particles at the next level. It must be also emphasized that the notion of infinity is relativistic. With respect to the p-adic norm infinite primes have unit norm for all finite and infinite primes so that there is nothing to become scared of!

Construction could make sense also for hyper-quaternionic and hyper-octonionic primes although non-commutativity and non-associativity pose technical challenges. One can also construct infinite number of real units as ratios of infinite integers with a precise number theoretic anatomy. The fascinating finding is that the quantum states labeled by standard model quantum numbers allow a representation as wave functions in the discrete space of these units. Space-time point becomes infinitely richly structured in the sense that one can associate to it a wave function in the space of real (or octonionic) units allowing to represent the WCW spinor fields. One can speak about algebraic holography or number theoretic Brahman=Atman identity and one can also say that the points of embedding space and space-time surface are subject to a number theoretic evolution. In philosophical mood one can of course also ask whether there exists a hierarchy of embedding spaces in which the embedding space at the lower level represents something with infinitesimal size in the sense of real topology and whether this hierarchy is accompanied also by a hierarchy of conscious entities.

This picture suggest that the Universe of quantum TGD might basically provide a physical representation of number theory allowing also infinite primes. The proposal suggests also a possible generalization of real numbers to a number system akin to hyper-reals introduced by Robinson in his non-standard calculus [A20] providing a rigorous mathematical basis for calculus. In fact, some rather natural requirements lead to a unique generalization for the concepts of integer, rational and real. Infinite integers and reals can be regarded as infinite-dimensional vector spaces with integer and real valued coefficients respectively. Same generalization could make sense for all classical number fields [A10, A4, A13] .

#### 4.1.2 Infinite primes and physics in TGD Universe

Several different views about how infinite primes, integers, and rationals might be relevant in TGD Universe have emerged.

##### 1. Infinite primes and super-symmetric quantum field theory

Consider next the physical interpretation.

1. The discovery of infinite primes suggested strongly the possibility to reduce physics to number theory. The construction of infinite primes can be regarded as a repeated second quantization of a super-symmetric arithmetic quantum field theory. This suggests that WCW spinor fields or at least the ground states of associated super-conformal representations [A15] (for super-conformal invariance see [A15] could be mapped to infinite primes in both bosonic and fermionic degrees of freedom. The process might generalize so that it applies in the case of quaternionic and octonionic primes and their hyper counterparts. This hierarchy of second quantizations means enormous generalization of physics to what might be regarded a physical counterpart for a hierarchy of abstractions about abstractions about.... The ordinary second quantized quantum physics corresponds only to the lowest level infinite primes.
2. The ordinary primes appearing as building blocks of infinite primes at the first level of the hierarchy could be identified as coding for p-adic primes assignable to fermionic and bosonic partons identified as 2-surfaces of a given space-time sheet. The hierarchy of infinite primes would correspond to hierarchy of space-time sheets defined by the topological condensate. This leads also to a precise identification of p-adic and real variants of bosonic partonic 2-surfaces as correlates of intention and action and pairs of p-adic and real fermionic partons as correlates for cognitive representations.
3. The idea that infinite primes characterize quantum states of the entire Universe, perhaps ground states of super-conformal representations, if not all states, could be taken further. It turns out that this idea makes sense when one considers discrete wave functions in the space of infinite primes and that one can indeed represent standard model quantum numbers in this manner.

4. The number theoretical supersymmetry suggests also space-time supersymmetry TGD framework. Space-time super-symmetry in its standard form is not possible in TGD Universe and this cheated me to believe that this supersymmetry is completely absent in TGD Universe. The progress in the understanding of the properties of the modified Dirac action however led to a generalization of the space-time super-symmetry as a dynamical and broken symmetry of quantum TGD [?] .

Here however emerges the idea about the number theoretic analog of color confinement. Rational (infinite) primes allow not only a decomposition to (infinite) primes of algebraic extensions of rationals but also to algebraic extensions of quaternionic and octonionic (infinite) primes. The physical analog is the decomposition of a particle to its more elementary constituents. This fits nicely with the idea about number theoretic resolution represented as a hierarchy of Galois groups defined by the extensions of rationals and realized at the level of physics in terms of Jones inclusions [K36] defined by these groups having a natural action on space-time surfaces, induced spinor fields, and on WCW spinor fields representing physical states [K7] .

### 2. Infinite primes and physics as number theory

The hierarchy of algebraic extensions of rationals implying corresponding extensions of p-adic numbers [A31, A22, A23, A19] suggests that Galois groups, which are the basic symmetry groups of number theory, should have concrete physical representations using induced spinor fields and configuration space spinor fields and also infinite primes and real units formed as infinite rationals. These groups permute zeros of polynomials and thus have a concrete physical interpretation both at the level of partonic 2-surfaces dictated by algebraic equations and at the level of braid hierarchy. The vision about the role of hyperfinite factors of  $II_1$  and of Jones inclusions as descriptions of quantum measurements with finite measurement resolution leads to concrete ideas about how these groups are realized.

$G_2$  acts as automorphisms of hyper-octonions and  $SU(3)$  as its subgroup respecting the choice of a preferred imaginary unit. The discrete subgroups of  $SU(3)$  permuting to each other hyper-octonionic primes are analogous to Galois group and turned out to play a crucial role in the understanding of the correspondence between infinite hyper-octonionic primes and quantum states predicted by quantum TGD.

### 3. The notion of finite measurement resolution as the key concept

TGD predicts several hierarchies: the hierarchy of space-time sheets, the hierarchy of infinite primes, the hierarchy of Jones inclusions identifiable in terms of finite measurement resolution [K36] , the dark matter hierarchy characterized by increasing values of  $\hbar$  [K10] , the hierarchy of extensions of a given p-adic number field. TGD inspired theory of consciousness predicts the hierarchy of selves and quantum jumps with increasing duration with respect to geometric time. These hierarchies should be closely related.

The notion of finite measurement resolution turns out to be the key concept: the p-adic norm of the rational defined by the infinite prime characterizes the angle measurement resolution for given p-adic prime  $p$ . It is essential that one has what might be called a state function reduction selecting a fixed p-adic prime which could be also infinite. This gives direct connections with cognition and with the p-adicization program relying also on angle measurement resolution. Also the value the integers characterizing the singular coverings of  $CD$  and  $CP_2$  defining as their product Planck constant characterize the measurement resolution for a given p-adic prime in  $CD$  and  $CP_2$  degrees of freedom. This conforms with the fact that elementary particles are characterized by two infinite primes. Hence finite measurement resolution ties tightly together the three threads of the number theoretic vision. Finite measurement resolution relates also closely to the inclusions of hyper-finite factors central for TGD inspired quantum measurement theory with finite measurement resolution.

### 4. Space-time correlates of infinite primes

Infinite primes code naturally for Fock states in a hierarchy of super-symmetric arithmetic quantum field theories. Quantum classical correspondence leads to ask whether infinite primes could also code for the space-time surfaces serving as symbolic representations of quantum states. This would a generalization of algebraic geometry would emerge and could reduce the dynamics of Kähler action to algebraic geometry and organize 4-surfaces to a physical hierarchy according to

their algebraic complexity. This conjecture should be consistent with two other conjectures about the dynamics of space-time surfaces (space-time surfaces as preferred extrema of Kähler action and space-time surfaces as quaternionic or co-quaternionic (as associative or co-associative) 4-surfaces of hyper-octonion space  $M^8$ ).

Quantum classical correspondence requires the map of the quantum numbers of WCW spinor fields to space-time geometry. The quantum numbers characterizing positive and negative energy parts of zero energy states couple directly to space-time geometry via the measurement interaction terms in Kähler action expressing the equality of classical conserved charges in Cartan algebra with their quantal counterparts for space-time surfaces in quantum superposition. This makes sense if classical charges parametrize zero modes. The localization in zero modes in state function reduction would be the WCW counterpart of state function collapse.

Therefore, if one wants to map infinite rationals to space-time geometry it is enough to map infinite primes to quantum numbers. This map might be achieved thanks to the detailed picture about the interpretation of the symmetries of infinite primes in terms of standard model symmetries. The notion of finite measurement resolution allows to deduce much more detailed about this correspondence. In particular, the rational defined by the infinite prime classifies the finite sub-manifold geometry defined by the discretization of the partonic 2-surface implied by the finite measurement resolution. Also a direct correlation between integers defining Planck constant and the “fermionic” part of the infinite prime emerges.

### 4.1.3 Infinite primes and cognition

The correlation of infinite primes with cognition is the first fascinating possibility and this possibility has stimulated several ideas.

1. One can define the notion of prime also for the algebraic extensions of rationals. The hierarchy of infinite primes associated with algebraic extensions of rationals leading gradually towards algebraic closure of rationals would in turn define cognitive hierarchy corresponding to algebraic extensions of p-adic numbers.
2. The introduction of infinite primes, integers, and rationals leads also to a generalization of classical number fields since an infinite algebra of real (complex, etc...) units defined by finite ratios of infinite rationals multiplied by ordinary rationals which are their inverses becomes possible. These units are not units in the p-adic sense and have a finite p-adic norm which can be differ from one. This construction generalizes also to the case of hyper- quaternions and -octonions although non-commutativity and in case of octonions also non-associativity pose technical problems. Obviously this approach differs from the standard introduction of infinitesimals in the sense that sum of infinitesimals (real zeros) is replaced by multiplication of real units meaning that the set of real and also more general units becomes infinitely degenerate.
3. Infinite primes form an infinite hierarchy so that the points of space-time and embedding space can be seen as infinitely structured and able to represent all imaginable algebraic structures. Certainly counter-intuitively, single space-time point -or more generally wave functions in the space of the units associated with the point- might be even capable of representing the quantum state of the entire physical Universe in its structure. For instance, in the real sense surfaces in the space of units correspond to the same real number 1, and single point, which is structure-less in the real sense could represent arbitrarily high-dimensional spaces as unions of real units. For real physics this structure is completely invisible and is relevant only for the physics of cognition. One can say that Universe is an algebraic hologram, and there is an obvious connection both with Brahman=Atman identity of Eastern philosophies and Leibniz’s notion of monad.
4. In ZEO hyper-octonionic units identified as ratios of the infinite integers associated with the positive and negative energy parts of the zero energy state define a representation of WCW spinor fields. The action of subgroups of  $SU(3)$  and rotation group  $SU(2)$  preserving hyper-octonionic and hyper-quaternionic primeness and identification of momentum and electro-weak charges in terms of components of hyper-octonionic primes makes this representation

unique. Hence Brahman-Atman identity has a completely concrete realization and fixes completely the quantum number spectrum including particle masses and correlations between various quantum numbers.

5. One can assign to infinite primes at  $n^{\text{th}}$  level of hierarchy rational functions of  $n$  rational arguments which form a natural hierarchical structure in that highest level corresponds to a polynomial with coefficients which are rational functions of the arguments at the lower level. One can solve one of the arguments in terms of lower ones to get a hierarchy of algebraic extensions. At the lowest level algebraic extensions of rationals emerge, at the next level algebraic extensions of space of rational functions of single variable, etc... This would suggest that infinite primes code for the correlation between quantum states and the algebraic extensions appearing in their physical description and characterizing their cognitive correlates. The hierarchy of infinite primes would also correlate with a hierarchy of logics of various orders (hierarchy of statements about statements about...).

## 4.2 Infinite Primes, Integers, And Rationals

The definition of the infinite integers and rationals is a straightforward procedure and structurally similar to a repeated second quantization of a super-symmetric quantum field theory but including also the number theoretic counterparts of bound states.

### 4.2.1 The first level of hierarchy

In the following the concept of infinite prime is developed gradually by stepwise procedure rather than giving directly the basic definitions. The hope is that the development of the concept in the same manner as it actually occurred would make it easier to understand it.

#### Step 1

One could try to define infinite primes  $P$  by starting from the basic idea in the proof of Euclid for the existence of infinite number of primes. Take the product of all finite primes and add 1 to get a new prime:

$$\begin{aligned} P &= 1 + X \quad , \\ X &= \prod_p p \quad . \end{aligned} \tag{4.1}$$

If  $P$  were divisible by finite prime then  $P - X = 1$  would be divisible by finite prime and one would encounter contradiction. One could of course worry about the possible existence of infinite primes smaller than  $P$  and possibly dividing  $P$ . The numbers  $N = P - k$ ,  $k > 1$ , are certainly not primes since  $k$  can be taken as a factor. The number  $P' = P - 2 = -1 + X$  could however be prime.  $P$  is certainly not divisible by  $P - 2$ . It seems that one cannot express  $P$  and  $P - 2$  as product of infinite integer and finite integer. Neither it seems possible to express these numbers as products of more general numbers of form  $\prod_{p \in U} p + q$ , where  $U$  is infinite subset of finite primes and  $q$  is finite integer.

## 4.3 How To Interpret The Infinite Hierarchy Of Infinite Primes?

From the foregoing it should be clear that infinite primes might play key role in quantum physics. One can even consider the possibility that physics reduces to a generalized number theory, and that infinite primes are crucial for understanding mathematically consciousness and cognition. Of course, one must leave open the question whether infinite primes really provide really the mathematics of consciousness or whether they are only a beautiful but esoteric mathematical construct. In this spirit the following subsections give only different points of view to the problem with no attempt to a coherent overall view.

### 4.3.1 Infinite primes and hierarchy of super-symmetric arithmetic quantum field theories

Infinite primes are a generalization of the notion of prime. They turn out to provide number theoretic correlates of both free, interacting and bound states of a super-symmetric arithmetic quantum field theory. It turns also possible to assign to infinite prime space-time surface as a geometric correlate although the original proposal for how to achieve this failed. Hence infinite primes serve as a bridge between classical and quantum and realize quantum classical correspondence stating that quantum states have classical counterparts, and has served as a basic heuristic guideline of TGD. More precisely, the natural hypothesis is that infinite primes code for the ground states of super-symplectic representations (for instance, ordinary particles correspond to states of this kind).

#### 1. Infinite primes and Fock states of a super-symmetric arithmetic QFT

The basic construction recipe for infinite primes is simple and generalizes to the quaternionic case.

1. Form the product of all primes and call it  $X$ :

$$X = \prod_p p .$$

2. Form the vacuum states

$$V_{\pm} = X \pm 1 .$$

3. From these vacua construct all *generating* infinite primes by the following process. Kick out from the Dirac sea some negative energy fermions: they correspond to a product  $s$  of first powers of primes:  $V \rightarrow X/s \pm s$  ( $s$  is thus square-free integer). This state represents a state with some fermions represented as holes in Dirac sea but no bosons. Add bosons by multiplying by integer  $r$ , which decomposes into parts as  $r = mn$ :  $m$  corresponding to bosons in  $X/s$  is product of powers of primes dividing  $X/s$  and  $n$  corresponds to bosons in  $s$  and is product of powers of primes dividing  $s$ . This step can be described as  $X/s \pm s \rightarrow mX/s \pm ns$ .

Generating infinite primes are thus in one-one correspondence with the Fock states of a super-symmetric arithmetic quantum field theory and can be written as

$$P_{\pm}(m, n, s) = \frac{mX}{s} \pm ns ,$$

where  $X$  is product of all primes at previous level.  $s$  is square free integer.  $m$  and  $n$  have no common factors, and neither  $m$  and  $s$  nor  $n$  and  $X/s$  have common factors.

The physical analog of the process is the creation of Fock states of a super-symmetric arithmetic quantum field theory. The factorization of  $s$  to a product of first powers of primes corresponds to many-fermion state and the decomposition of  $m$  and  $n$  to products of powers of prime correspond to bosonic Fock states since  $p^k$  corresponds to  $k$ -particle state in arithmetic quantum field theory.

#### 2. More complex infinite primes as counterparts of bound states

Generating infinite primes are not all that are possible. One can construct also polynomials of the generating primes and under certain conditions these polynomials are non-divisible by both finite primes and infinite primes already constructed. As found, the conjectured effective 2-dimensionality for hyper-octonionic primes allows the reduction of polynomial representation of hyper-octonionic primes to that for hyper-complex primes. This would be in accordance with the effective 2-dimensionality of the basic objects of quantum TGD.

The physical counterpart of  $n$ :th order irreducible polynomial is as a bound state of  $n$  particles whereas infinite integers constructed as products of infinite primes correspond to non-bound but interacting states. This process can be repeated at the higher levels by defining the vacuum state to be the product of all primes at previous levels and repeating the process. A repeated second quantization of a super-symmetric arithmetic quantum field theory is in question.

The infinite primes represented by irreducible polynomials correspond to quantum states obtained by mapping the superposition of the products of the generating infinite primes to a superposition of the corresponding Fock states. If complex rationals are the coefficient field for infinite integers, this gives rise to states in a complex Hilbert space and irreducibility corresponds to a superposition of states with varying particle number and the presence of entanglement. For instance, the superpositions of several products of type  $\prod_{i=1,\dots,n} P_i$  of  $n$  generating infinite primes are possible and in general give rise to irreducible infinite primes decomposing into a product of infinite primes in algebraic extension of rationals.

### 3. Infinite rationals viz. quantum states and space-time surfaces

The most promising answer to the question how infinite rationals correspond to space-time surfaces is discussed in detail in the next section. Here it is enough to give only the basic idea.

1. In ZEO hyper-octonionic units (in real sense) defined by ratios of infinite integers have an interpretation as representations for pairs of positive and negative energy states. Suppose that the quantum number combinations characterizing positive and negative energy quantum states are representable as superpositions of real units defined by ratios of infinite integers at each point of the space-time surface. If this is true, the quantum classical correspondence coded by the measurement interaction term of the Kähler-Dirac action maps the quantum numbers also to space-time geometry and implies a correspondence between infinite rationals and space-time surfaces.
2. The space-time surface associated with the infinite rational is in general not a union of the space-time surfaces associated with the primes composing the integers defining the rational. There the classical description of interactions emerges automatically. The description of classical states in terms of infinite integers would be analogous to the description of many particle states as finite integers in arithmetic quantum field theory. This mapping could in principle make sense both in real and p-adic sectors of WCW.

The finite primes which correspond to particles of an arithmetic quantum field theory present in Fock state, correspond to the space-time sheets of finite size serving as the building blocks of the space-time sheet characterized by infinite prime.

### 4. What is the interpretation of the higher level infinite primes?

Infinite hierarchy of infinite primes codes for a hierarchy of Fock states such that many-particle Fock states of a given level serve as elementary particles at next level. The unavoidable conclusion is that higher levels represent totally new physics not described by the standard quantization procedures. In particular, the assignment of fermion/boson property to arbitrarily large system would be in some sense exact. Topologically these higher level particles could correspond to space-time sheets containing many-particle states and behaving as higher level elementary particles.

This view suggests that the generating quantum numbers are present already at the lowest level and somehow coded by the hyper-octonionic primes taking the role of momentum quantum number they have in arithmetic quantum field theories. The task is to understand whether and how hyper-octonionic primes can code for quantum numbers predicted by quantum TGD.

The quantum numbers coding higher level states are collections of quantum numbers of lower level states. At geometric level the replacement of the coefficients of polynomials with rational functions is the equivalent of replacing single particle states with new single particle states consisting of many-particle states.

#### **4.3.2 Infinite primes, the structure of many-sheeted space-time, and the notion of finite measurement resolution**

The mapping of infinite primes to space-time surfaces codes the structure of infinite prime to the structure of space-time surface in a rather non-implicit manner, and the question arises about the concrete correspondence between the structure of infinite prime and topological structure of the space-time surface. It turns out that the notion of finite measurement resolution is the key concept: infinite prime characterizes angle measurement resolution. This gives a direct connection with the p-adicization program relying also on angle measurement resolution as well as a connection with the



hierarchy of Planck constants. Finite measurement resolution relates also closely to the inclusions of hyper-finite factors central for TGD inspired quantum measurement theory.

### 1. The first intuitions

The concrete prediction of the general vision is that the hierarchy of infinite primes should somehow correspond to the hierarchy of space-time sheets or partonic 2-surfaces if one accepts the effective 2-dimensionality. The challenge is to find space-time counterparts for infinite primes at the lowest level of the hierarchy.

One could hope that the Fock space structure of infinite prime would have a more concrete correspondence with the structure of the many-sheeted space-time. One might that the space-time sheets labeled by primes  $p$  would directly correspond to the primes appearing in the definition of infinite prime. This expectation seems to be too simplistic.

1. What seems to be a safe guess is that the simplest infinite primes at the lowest level of the hierarchy should correspond to elementary particles. If inverses of infinite primes correspond to negative energy space-time sheets, this would explain why negative energy particles are not encountered in elementary particle physics.
2. More complex infinite primes at the lowest level of the hierarchy could be interpreted in terms of structures formed by connecting these structures by join along boundaries bonds to get space-time correlates of bound states. Even simplest infinite primes must correspond to bound state structures if the condition that the corresponding polynomial has real-rational coefficients is taken seriously.

Infinite primes at the lowest level of hierarchy correspond to several finite primes rather than single finite prime. The number of finite primes is however finite.

1. A possible interpretation for multi-p property is in terms of multi-p p-adic fractality prevailing in the interior of space-time surface. The effective p-adic topology of these space-time sheets would depend on length scale. In the longest scale the topology would correspond to  $p_n$ , in some shorter length scale there would be smaller structures with  $p_{n-1} < p_n$ -adic topology, and so on... . A good metaphor would be a wave containing ripples, which in turn would contain still smaller ripples. The multi-p p-adic fractality would be assigned with the 4-D space-time sheets associated with elementary particles. The concrete realization of multi-p p-adicity would be in terms of infinite integers coming as power series  $\sum x_n N^n$  and having interpretation as p-adic numbers for any prime dividing  $N$ .
2. Effective 2-dimensionality would suggest that the individual p-adic topologies could be assigned with the 2-dimensional partonic surfaces. Thus infinite prime would characterize at the lowest level space-time sheet and corresponding partonic 2-surfaces. There are however reasons to think that even single partonic 2-surface corresponds to a multi-p p-adic topology.

### 2. Do infinite primes code for the finite measurement resolution?

The above describe heuristic picture is not yet satisfactory. In order to proceed, it is good to ask what determines the finite prime or set of them associated with a given partonic 2-surface. It is good to recall first the recent view about the p-adicization program relying crucially on the notion of finite measurement resolution.

1. The vision about p-adicization characterizes finite measurement resolution for angle measurement in the most general case as  $\Delta\phi = 2\pi M/N$ , where  $M$  and  $N$  are positive integers having no common factors. The powers of the phases  $\exp(i2\pi M/N)$  define identical Fourier basis irrespective of the value of  $M$  and measurement resolution does not depend on the value of  $M$ . Situation is different if one allows only the powers  $\exp(i2\pi kM/N)$  for which  $kM < N$  holds true: in the latter case the measurement resolutions with different values of  $M$  correspond to different numbers of Fourier components. If one regards  $N$  as an ordinary integer, one must have  $N = p^n$  by the p-adic continuity requirement.

2. One can also interpret  $N$  as a  $p$ -adic integer. For  $N = p^n M$ , where  $M$  is not divisible by  $p$ , one can express  $1/M$  as a  $p$ -adic integer  $1/M = \sum_{k \geq 0} M_k p^k$ , which is infinite as a real integer but effectively reduces to a finite integer  $K(p) = \sum_{k=0}^{N-1} M_k p^k$ . As a root of unity the entire phase  $\exp(i2\pi M/N)$  is equivalent with  $\exp(i2\pi R/p^n)$ ,  $R = K(p)M \pmod{p^n}$ . The phase would non-trivial only for  $p$ -adic primes appearing as factors in  $N$ . The corresponding measurement resolution would be  $\Delta\phi = R2\pi/N$  if modular arithmetics is used to define the measurement resolution. This works at the first level of the hierarchy but not at higher levels. The alternative manner to assign a finite measurement resolution to  $M/N$  for given  $p$  is as  $\Delta\phi = 2\pi|N/M|_p = 2\pi/p^n$ . In this case the small fermionic part of the infinite prime would fix the measurement resolution. The argument below shows that only this option works also at the higher levels of hierarchy and is therefore more plausible.
3.  $p$ -Adicization conditions in their strong form require that the notion of integration based on harmonic analysis [A5] in symmetric spaces [A16] makes sense even at the level of partonic 2-surfaces. These conditions are satisfied if the partonic 2-surfaces in a given measurement resolution can be regarded as algebraic continuations of discrete surfaces whose points belong to the discrete variant of the  $\delta M_{\pm}^4 \times CP_2$ . This condition is extremely powerful since it effectively allows to code the geometry of partonic 2-surfaces by the geometry of finite submanifold geometries for a given measurement resolution. This condition assigns the integer  $N$  to a given partonic surface and all primes appearing as factors of  $N$  define possible effective  $p$ -adic topologies assignable to the partonic 2-surface.

How infinite primes could then code for the finite measurement resolution? Can one identify the measurement resolution for  $M/N = M/(Rp^n)$  as  $\Delta\phi = ((M/R) \pmod{p^n}) \times 2\pi/p^n$  or as  $\Delta\phi = 2\pi/p^n$ ? The following argument allows only the latter option.

1. Suppose that  $p$ -adic topology makes sense also for infinite primes and that state function reduction selects power of infinite prime  $P$  from the product of lower level infinite primes defining the integer  $N$  in  $M/N$ . Suppose that the rational defined by infinite integer defines measurement resolution also at the higher levels of the hierarchy.
2. The infinite primes at the first level of hierarchy representing Fock states are in one-one correspondence with finite rationals  $M/N$  for which integers  $M$  and  $N$  can be chosen to characterize the infinite bosonic part and finite fermionic part of the infinite prime. This correspondence makes sense also at higher levels of the hierarchy but  $M$  and  $N$  are infinite integers. Also other option obtained by exchanging “bosonic” and “fermionic” but later it will be found that only the first identification makes sense.
3. The first guess is that the rational  $M/N$  characterizing the infinite prime characterizes the measurement resolution for angles and therefore partially classifies also the finite submanifold geometry assignable to the partonic 2-surface. One should define what  $M/N = ((M/R) \pmod{P^n}) \times P^{-n}$  is for infinite primes. This would require expression of  $M/R$  in modular arithmetics modulo  $P^n$ . This does not make sense.
4. For the second option the measurement resolution defined as  $\Delta\phi = 2\pi|N/M|_P = 2\pi/P^n$  makes sense. The Fourier basis obtained in this manner would be infinite but all states  $\exp(ik/P^n)$  would correspond in real sense to real unity unless one allows  $k$  to be infinite  $P$ -adic integer smaller than  $P^n$  and thus expressible as  $k = \sum_{m < n} k_m P^m$ , where  $k_m$  are infinite integers smaller than  $P$ . In real sense one obtains all roots  $\exp(iq2\pi)$  of unity with  $q < 1$  rational. For instance, for  $n = 1$  one can have  $0 < k/P < 1$  for a suitably chosen infinite prime  $k$ . Thus one would have essentially continuum theory at higher levels of the hierarchy. The purely fermionic part  $N$  of the infinite prime would code for both the number of Fourier components in discretization for each power of prime involved and the ratio characterize the angle resolution.

The proposed relation between infinite prime and finite measurement resolution implies very strong number theoretic selection rules on the reaction vertices.

1. The point is that the vertices of generalized Feynman diagrams correspond to partonic 2-surfaces at which the ends of light-like 3-surfaces describing the orbits of partonic 2-surfaces join together. Suppose that the partonic 2-surfaces appearing at both ends of the propagator lines correspond to same rational as finite sub-manifold geometries. If so, then for a given p-adic effective topology the integers assignable to all lines entering the vertex must contain this p-adic prime as a factor. Particles would correspond to integers and only the particles having common prime factors could appear in the same vertex.
2. In fact, already the work with modelling dark matter [K10] led to ask whether particle could be characterized by a collection of p-adic primes to which one can assign weak, color, em, gravitational interactions, and possibly also other interactions. It also seemed natural to assume that only the space-time sheets containing common primes in this collection can interact. This inspired the notions of relative and partial darkness. An entire hierarchy of weak and color physics such that weak bosons and gluons of given physics are characterized by a given p-adic prime  $p$  and also the fermions of this physics contain space-time sheet characterized by same p-adic prime, say  $M_{89}$  as in case of weak interactions. In this picture the decay widths of weak bosons do not pose limitations on the number of light particles if weak interactions for them are characterized by p-adic prime  $p \neq M_{89}$ . Same applies to color interactions.

The possibility of multi-p p-adicity raises the question about how to fix the p-adic prime characterizing the mass of the particle. The mass scale of the contribution of a given throat to the mass squared is given by  $p^{-n/2}$ , where  $T = 1/n$  corresponds to the p-adic temperature of throat. Hence the dominating contribution to the mass squared corresponds to the smallest prime power  $p^n$  associated with the throats of the particle. This works if the integers characterizing other particles than graviton are divisible by the gravitonic p-adic prime or a product of p-adic primes assignable to graviton. If the smallest power  $p^n$  assignable to the graviton is large enough, the mass of graviton is consistent with the empirical bounds on it. The same consideration applies in the case of photons. Recall that the number theoretically very natural condition that in ZEO the number of generalized Feynman graphs contributing to a given process is finite is satisfied if all particles have a non-vanishing but arbitrarily small p-adic thermal mass [K37].

### 3. Interpretational problem

The identification of infinite prime as a characterizer of finite measurement resolution looks nice but there is an interpretational problem.

1. The model characterizing the quantum numbers of WCW spinor fields to be discussed in the next section involves a pair of infinite primes  $P_+$  and  $P_-$  corresponding to the two vacuum primes  $X \pm 1$ . Do they correspond to two different measurement resolutions perhaps assignable to CD and  $CP_2$  degrees of freedom?
2. Different measurement resolutions in CD and  $CP_2$  degrees of freedom need not be not a problem as long as one considers only the discrete variants of symmetric spaces involved. What might be a problem is that in the general case the p-adic primes associated with CD and  $CP_2$  degrees of freedom would not be same unless the integers  $N_+$  and  $N_-$  are assumed to have have same prime factors (they indeed do if  $p^0 = 1$  is formally counted as prime power factors).
3. The idea of assigning different p-adic effective topologies to CD and  $CP_2$  does not look attractive. Both CD and  $CP_2$  and thus also partonic 2-surface could however possess simultaneously both p-adic effective topologies. This kind of option might make sense since the integers representable as infinite powers series of integer  $N$  can be regarded as p-adic integers for all prime factors of  $N$ . As a matter fact, this kind of multi-p p-adicity could make sense also for the partonic 2-surfaces characterized by a measurement resolution  $\Delta\phi = 2\pi M/N$ . One would have what might be interpreted as  $N_+N_-$ -adicity.
4. It will be found that quantum measurement means also the measurement of the p-adic prime selecting same p-adic prime from  $N_+$  and  $N_-$ . If  $N_{\pm}$  is divisible only by  $p^0 = 1$ , the corresponding angle measurement resolution is trivial. From the point of view of consciousness

state function reduction selects also the p-adic prime characterizing the cognitive representation which is very natural since quantum superpositions of different p-adic topologies are not natural physically.

### 4.3.3 How the hierarchy of Planck constants could relate to infinite primes and p-adic hierarchy?

Besides the hierarchy of space-time sheets, TGD predicts, or at least suggests, several hierarchies such as the hierarchy of infinite primes, the hierarchy of Jones inclusions identifiable in terms of finite measurement resolution [K36], the hierarchy of super-symplectic gauge symmetry breakings closely related to the dark matter hierarchy characterized by increasing values of  $h_{eff} = n \times h$  [K10], the hierarchy of extensions of given p-adic number field associated with algebraic extensions of rationals, and the hierarchy of selves and quantum jumps with increasing duration with respect to geometric time. There are good reasons to expect that these hierarchies are closely related. Number theoretical considerations give hopes about developing a more quantitative vision about the relationship between these hierarchies, in particular between the hierarchy of infinite primes, p-adic length scale hierarchy, and the hierarchy of Planck constants.

This idea can be indeed made concrete.

1. The hierarchy of algebraic extensions of rationals corresponds to the hierarchy of quantum criticalities labelled by integer  $n = h_{eff}/h$ . There is a temptation to identify  $n$  as the product of ramified primes of the algebraic extension or its power. In accordance with the number theoretic vision number theoretic criticality would correspond to quantum criticality. The idea is that ramified primes are analogous to multiple roots of polynomial and criticality indeed corresponds to this kind of situation.
2. Infinite primes at the  $n$ :th level of hierarchy representing analogs of bound states correspond to irreducible polynomials of  $n$ -variables identifiable as polynomials of  $z_n$  with coefficients, which are polynomials of  $z_1, \dots, z_{n-1}$ . At the first level of hierarchy bound states correspond to irreducible polynomials of single variable and their roots define irreducible algebraic extensions of rationals. Infinite integers in turn correspond to products of reducible polynomials defining reducible extensions. The infinite integers at the first level of hierarchy would define the hierarchy of algebraic extensions of rationals in turn defining a hierarchy of quantum criticalities. This observation might generalize to the higher levels of hierarchy of infinite primes so that infinite primes would be part of quantum TGD although in much more abstract sense as thought originally.

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